Study of the Effect of the Breakup on the Fusion Cross Section of the Systems ^{6,7}Li+⁵⁹Co

F.A. Souza, A. Szanto de Toledo, M.G. Munhoz, J. Takahashi, N. Carlin, A.A.P. Suaide, M.M. de Moura, E.M. Szanto,

Instituto de Física, Universidade de São Paulo, Caixa Postal 66318, 05315-970, São Paulo, SP, Brazil

C. Beck,

Centre National de la Recherche Scientifique Université Louis Pasteur, Strasbourg Cedex 2, France

and S.J. Sanders

University of Kansas, Lawrence, KS 66045, USA

Received on 6 October, 2003

Fusion cross section excitation functions were measured for the ${}^{6,7}\text{Li}+{}^{59}\text{Co}$ systems ($12 \le E_{lab} \le 26MeV$) aiming the investigation of the effect of the breakup process on the fusion cross section. The experimental method consisted on the detection of γ -rays from the evaporation residues. Coupled channels calculations have been performed for both systems. The comparison to the experimental results support the conclusion that there is neither supression nor enhancement at energies above the coulomb barrier. A slight enhancement at energies below the barrier is observed for the ${}^{6}\text{Li}+{}^{59}\text{Co}$ cross section over the one for ${}^{7}\text{Li}+{}^{59}\text{Co}$. In order to better quantify the effect of the breakup process, coincidence measurements for the breakup products are being performed.

The effect of collective degrees of freedom on the fusion process has been extensively investigated over the past few years [1]. A significant enhancement of the sub-barrier fusion cross section is often found as compared to the predictions of one-dimensional barrier penetration models. This enhancement is understood in terms of dynamical processes involving couplings to collective inelastic excitations of the target and/or projectile. A precise determination of the "barrier distributions" leading to the enhancement requires an understanding of the dominant channels that couple to the fusion channel [2]. However, in the case of reactions where at least one of the colliding nuclei has a sufficiently low binding energy so that breakup becomes an important process, conflicting experimental and theoretical results have been reported [3-11].

The many questions regarding the influence of breakup become more relevant with the recent availability of radioactive beams and the renewed interest in super-heavy element formation. Radioactive ion beams are likely to produce intense breakup yields. In the fusion processes, and more specifically in the fusion of weakly bound nuclei, that can be used in super-heavy element studies, two different and independent processes can be distinguished both experimentally and theoretically. One, denoted as "Complete Fusion" (CF), is associated with the capture of all of the projectile constituents by the target. The other, denoted as "Incomplete Fusion" (ICF) or Partial Fusion, occurs when part of the projectile is captured by the target and the remaining part escapes. "Total Fusion" is understood as the sum of these two processes (CF + ICF). In order to avoid misinterpretations, a clear definition of the experimental and theoretical quantities being compared is essential. This is the main difficulty in comparing data and/or calculations from different authors. Depending on the theoretical approach, different results are achieved. If fusion occurs incoherently, i.e., with possible breakup of the incoming particle [4, 5], the survival probability of this particle prior to fusion, is lower than unity, resulting in a decrease in the effective entrancechannel flux and a corresponding reduction of the fusion cross section. On the other hand, if the breakup channel is coupled coherently to the fusion channel, as expected in references [3, 6], entrance barrier fluctuations will lead to an effective lowering of this barrier, resulting in an overall cross section enhancement. Recent coupled-channel calculations based on the coupling to states in the continuum, which is split in discrete bins, (CDCC) [7], predict the coexistence of both dynamical descriptions: a coherent process at subbarrier energies leads to a cross section enhancement and, at higher energies, a significant loss of entrance-channel flux leads to an overall cross section suppression.

A similar situation is found experimentally. In the case of heavy systems, two different situations are observed. Data for fusion reactions presented in reference [8, 9, 10] indicate the occurrence of a significant cross-section enhancement. On the other hand, in reference [11] a suppression is observed. References [12] and [13] report the simultaneous occurrence of both effects in different energy regimes. In the case of light systems, a correlation between the "Fusion Inhibition Factor", $FIF = \sigma_{fusion} / \sigma_{reaction}$ and the binding energy has been observed [14, 15, 16] (σ_{fusion} represents the experimental fusion cross section at the energy of two times the fusion barrier, and $\sigma_{reaction}$ represents the reaction cross section extracted from Optical Model fits to the elastic scattering. For details see ref. [14] and [15]) (see Fig. 1). However, it has been found [15] that two factors contribute to FIF. One is the relative importance of the surface to core density distribution that contributes to fusion (which increases for very light nuclei) and another is related to the survival probability (which increases with binding energy). Fusion reactions involving medium weight nuclei were also investigated by means of the ^{6,7}Li+⁵⁹Co reactions [17]. Fusion cross sections are presented in Fig. 2 for the ⁶Li and ⁷Li induced reactions as well as the ratio $R = \sigma_{^{6}Li} / \sigma_{^{7}Li}$. (The ⁷Li induced reaction is used as reference due to its higher binding energy (2.47 MeV) when compared to 1.47 MeV for ⁶Li). Results from coupled channel calculations for these systems are presented in Fig. 2. Predictions based on CDCC [18], shown in Fig. 3, indicate that no significant suppression is expected at energies above the barrier energies but, at sub-barrier energies, a fusion cross section enhancement is predicted due to the coupling of the breakup channel. A complete accounting for the coupling to the breakup channel, as well as the correct description of the reaction dynamics, requires the explicit measurement of yields leading to breakup itself. Some experiments have already addressed this issue [19, 20, 21, 22]. In the present case of the ^{6,7}Li+⁵⁹Co reactions, we have initiated this investigation (see Fig. 4) and have identified the Li states involved in the breakup process.



Figure 2. Energy dependence of the ratio $R = \sigma_{6Li}/\sigma_{7Li}$ between the fusion cross sections for ⁶Li and ⁷Li induced reactions. Error bars reflect the large systematic errors. The solid and dashed curves correspond to the case of fits of the experimental data to Single Barrier Penetration Model. Dotted curves correspond to two uncoupled CCFULL calculations [17] with and without reorientation effects.



Figure 1. Fusion Inhibition Factor $FIF = \sigma_{fusion}/\sigma_{reaction}$ (for a set of reactions identified in reference [14]) as a function of the binding energy of the colliding nuclei.



Figure 3. Total fusion excitation functions for ${}^{6}\text{Li}+{}^{59}\text{Co}$ (full dots) and ${}^{7}\text{Li}+{}^{59}\text{Co}$ (full triangles), which are normalized with the cross sections in absence of couplings to breakup channels. For each reaction, the incident energy is normalized with the Coulomb barrier V_{B} of the bare potential [18]. The curves are drawn to guide the eyes.



Figure 4. Energy correlation plots E_{α} vs E_d , for α particles and deuterons, detected in coincidence for the ⁵⁹Co (⁶Li, α d) reaction. Projections of the E_{α} (right panel) and E_d (lower panel) are also shown.

Acknowledgments

This work was supported in part by the Conselho Nacional de Desenvolvimento Científico e Tecnológico (CNPq) and Fundação de Amparo à Pesquisa do Estado de São Paulo (FAPESP)-Brasil.

References

[1] M. Dasgupta et al., Ann. Rev. Nucl. Part. Sci. **48**, 402 (1998) and references therein.

- [2] N. Rowley et al., Phys. Lett. B254, 25 (1991).
- [3] C.H. Dasso et al., Phys. Rev. C50, R12 (1994) and Nucl. Phys. A597, 473 (1996).
- [4] L.F. Canto et al., Phys. Rev. C 52, 1 (1995).
- [5] N. Takigawa et al., Phys. Rev. C47, R2470 (19930).
- [6] K. Hagino et al., Phys. Rev. C61, 037602 (2000).
- [7] A. Diaz-Torres and I.J. Thompson, Phys. Rev. C65, 024606 (2002).
- [8] K. Yoshida et al., Phys. Rev. Lett. 389, 457 (1996).
- [9] J.J. Kolata et al., Phys. Rev. Lett. 81, 4580 (1988).
- [10] M. Trotta et al., Phys. Rev. Lett. 84, 2342 (2000).
- [11] K.E. Rehm et al., Phys. Rev. Lett. 81, 3341 (1998).
- [12] N. Dasgupta et al., Phys. Rev. Lett. 82,1395 (1999) and Phys. Rev. C66, 041602R (2002).
- [13] V. Tripathi et al., Phys. Rev. Lett. 88, 172701-1 (2001).
- [14] J. Takahashi et al., Phys. Rev. Lett. 78, 30 (1997).
- [15] A. Szanto de Toledo et al., Nucl. Phys. A679, 175 (2000) and R. Cabezas et al., Phys. Rev. C60, 067602 (1999).
- [16] R.M. Anjos et al., Phys. Lett. B 534, 295 (2002).
- [17] C. Beck et al., Phys. Rev. C 67, 54602 (2003).
- [18] A. Diaz-Torres et al., Phys. Rev. C (to be published and priv. com.).
- [19] D.J. Hinde et al., Phys. Rev. Lett. 89, 272701 (2002).
- [20] K. Rusek et al., Phys. Rev. C 67, 41604 (2003).
- [21] J.F. Liang et al., Phys. Rev. C 67, 44603 (2003).
- [22] C. Signorini et al., Phys. Rev. C 67, 44607 (2003).