

MAGNETOCALORIC AND BAROCALORIC EFFECTS IN METALS

OUTLINE

PRELIMINARIES

THEORY AND CALCULATIONS

APPLICATIONS

PROSPECTS

PRELIMINARIES

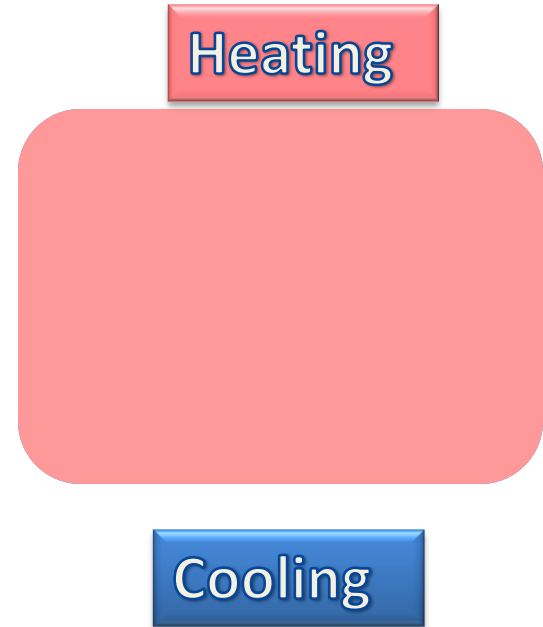
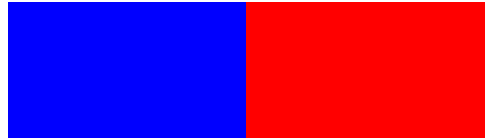
MAGNETOCALORIC EFFECT

HISTORICAL FACTS

MAGNETIC REFRIGERATION

BAROCALORIC EFFECT

MAGNETOCALORIC EFFECT



HISTORICAL FACTS

Thomson's work 1878 (Comments)

Thomson W. "Thermoelastic, thermomagnetic and pyroelectric properties of matter" *Phil. Mag. Ser 5* 5(1878)4.

Weiss and Piccard's work 1917: Experimental discovery

Weiss, P., Piccard A., "Le phénomène magnétocalorique" *J. Phys.* 7(1917)103

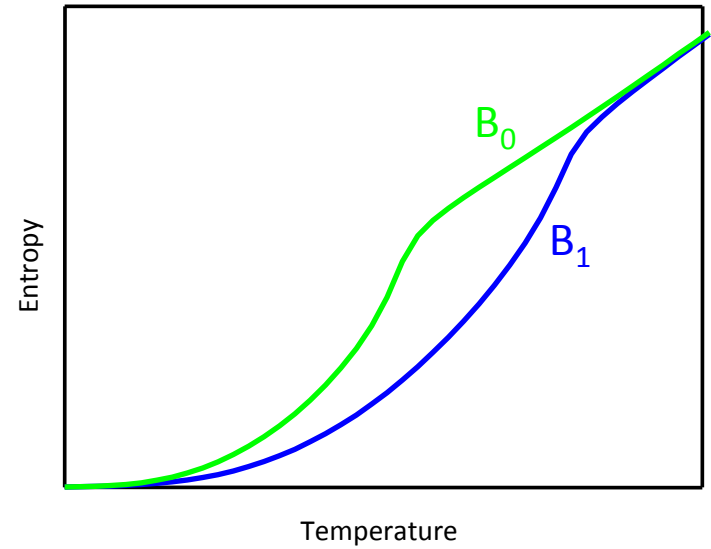
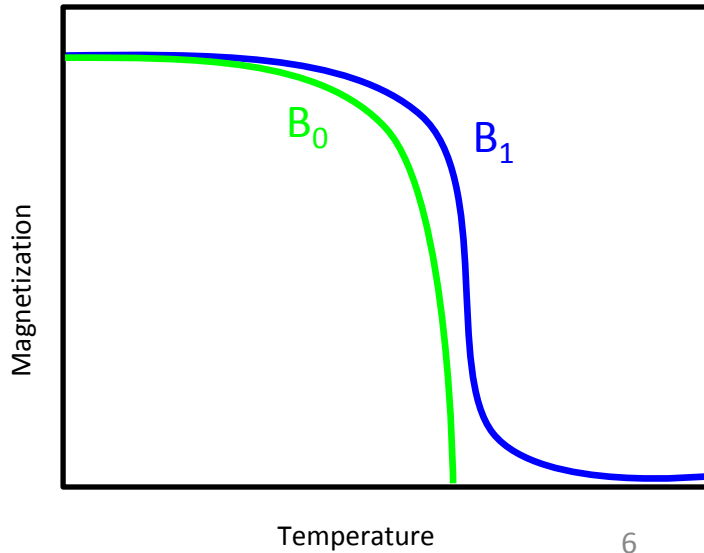
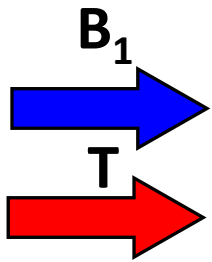
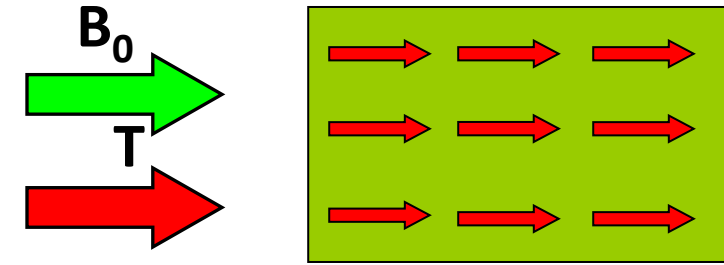
Brown's work - 1978: Room temperature magnetic refrigerator

Brown, G. V., "Heat pump near room temperature" *J. Appl. Phys.* 47(1976)3673.

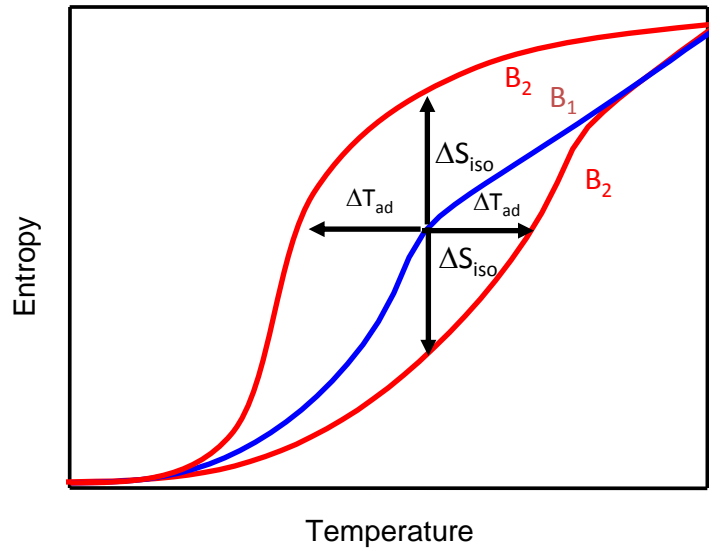
Gschneidner's work - 1997: Giant magnetocaloric effect: $\text{Gd}_5\text{Si}_2\text{Ge}_2$

V. K. Pecharsky, K. A. Gschneidner Jr, "Giant magnetocaloric effect in $\text{Gd}_5\text{Si}_2\text{Ge}_2$ " *Phys. Rev. Lett.* 78 (1997) 4494

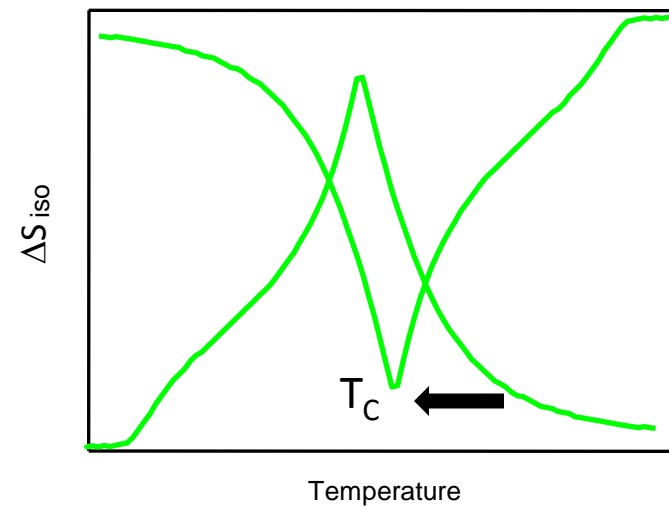
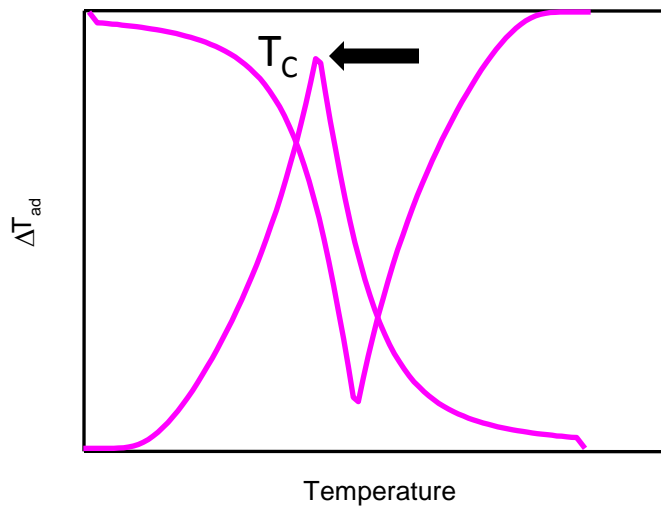
MAGNETOCALORIC EFFECT



MAGNETOCALORIC QUANTITIES



INVERSE MAGNETOCALORIC EFFECT



How to measure?

EXPERIMENTAL TECHNIQUES

Adiabatic temperature change (ΔT_{ad})

Direct measurements: Thermopar

Indirect measurements: Specific heat

Indirect measurements: Specific heat and magnetization

Isothermal entropy change (ΔS_{iso})

Indirect measurements: Specific heat

Indirect measurements: magnetization

EQUIPMENTS

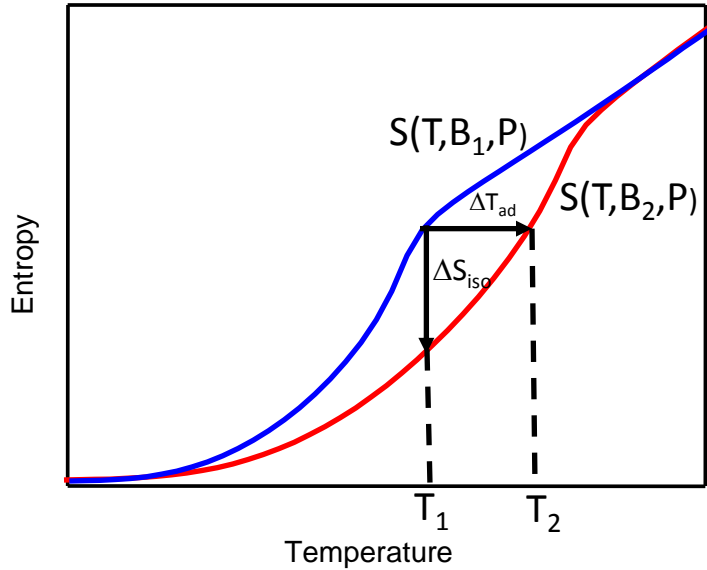
Calorimeter: ΔS_{iso} and ΔT_{ad}

Calorimeter and VSM/Squid: ΔS_{iso} and ΔT_{ad}

Squid/VSM : ΔS_{iso}

How to calculate?

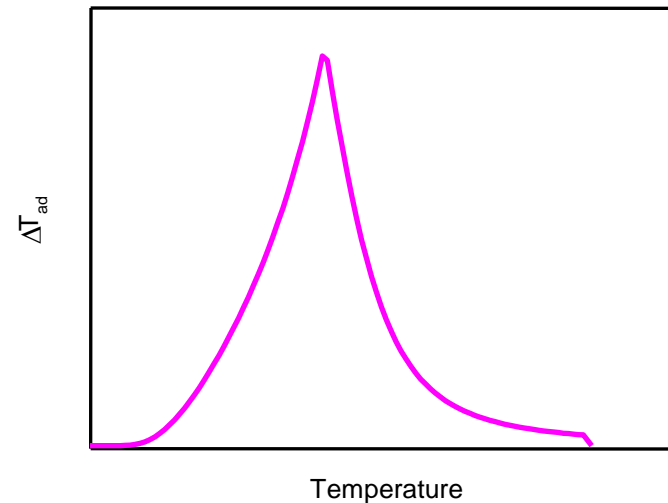
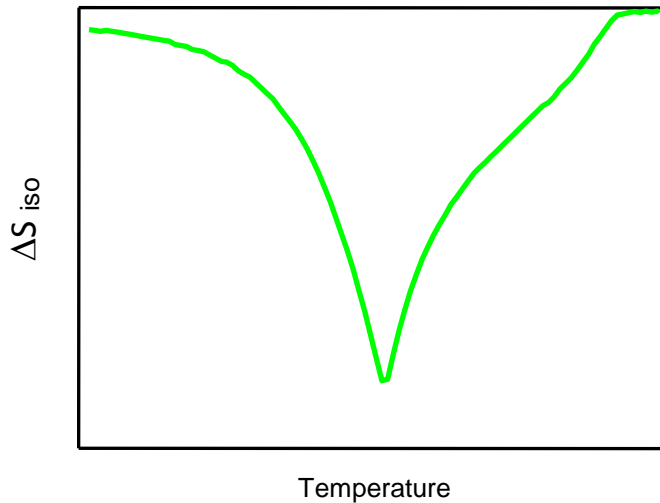
MAGNETOCALORIC EFFECT



$$\Delta S_{iso}(T, \Delta B, P) = S(T, B_2, P) - S(T, B_1, P)$$

$$\Delta T_{ad}(T, \Delta B, P) = T_2 - T_1$$

$$S(T, B_2, P) = S(T, B_1, P)$$



THERMODYNAMICS OF THE MCE (ΔS_{iso})

ENTROPY CHANGE

$$dS(T, B, P) = \left[\frac{\partial S(T, B, P)}{\partial T} + \frac{\delta S(T_C, B, P)}{\delta T} \right]_{B, P} dT + \left[\frac{\partial S(T, B, P)}{\partial B} + \frac{\delta S(T, B_C, P)}{\delta B} \right]_{T, P} dB + \left[\frac{\partial S(T, B, P)}{\partial P} + \frac{\delta S(T, B, P_C)}{\delta P} \right]_{T, B} dP$$

ISOTHERMAL / ISOBARIC PROCESS

$$\Delta S_{iso}(T, \Delta B, P) = \int_{B_1}^{B_2} \left[\frac{\partial S(T, B, P)}{\partial B} + \frac{\delta S(T, B_C, P)}{\delta B} \right]_{T, P} dB$$

SECOND ORDER TRANSITION

$$\Delta S_{iso}(T, \Delta B, P) = \int_{B_1}^{B_2} \left[\frac{\partial S(T, B, P)}{\partial B} \right]_{T, P} dB$$

THERMODYNAMICS OF THE MCE (ΔT_{ad})

ENTROPY CHANGE

$$dS(T, B, P) = \left[\frac{\partial S(T, B, P)}{\partial T} + \frac{\delta S(T_C, B, P)}{\delta T} \right]_{B, P} dT + \left[\frac{\partial S(T, B, P)}{\partial B} + \frac{\delta S(T, B_C, P)}{\delta B} \right]_{T, P} dB + \left[\frac{\partial S(T, B, P)}{\partial P} + \frac{\delta S(T, B, P_C)}{\delta P} \right]_{T, B} dP$$

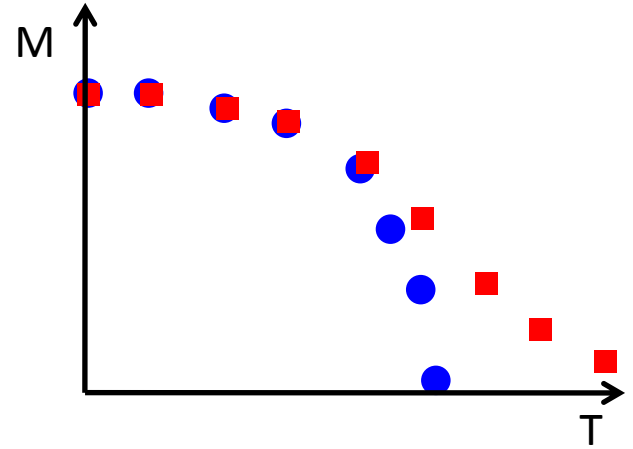
$$\Delta T_{ad}(T, \Delta B, P) = - \int_{B_1}^{B_2} \frac{1}{C(T, B, P)} \left[\frac{\partial S(T, B, P)}{\partial B} + \frac{\delta S(T, B_C, P)}{\delta B} \right]_{T, P} dB$$

$$\Delta T_{ad}(T, \Delta B, P) = - \int_{B_1}^{B_2} \frac{1}{C(T, B, P)} \left[\frac{\partial S(T, B, P)}{\partial B} \right]_{T, P} dB$$

THERMODYNAMICS: MAXWELL RELATION

MCE QUANTITIES

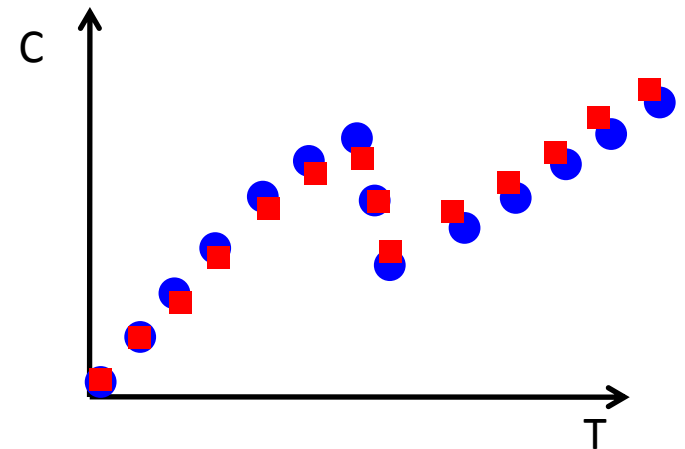
$$\Delta S_{iso}(T, \Delta B, P) = \int_{B_1}^{B_2} \left[\frac{\partial M(T, B, P)}{\partial T} \right]_{B, P} dB$$
$$\Delta T_{ad}(T, \Delta B, P) = - \int_{B_1}^{B_2} \frac{1}{C(T, B, P)} \left[\frac{\partial M(T, B, P)}{\partial T} \right]_{B, P} dB$$



MAXWELL RELATION

$$\left[\frac{\partial S(T, B, P)}{\partial B} \right]_{T, P} = \left[\frac{\partial M(T, B, P)}{\partial T} \right]_{B, P}$$

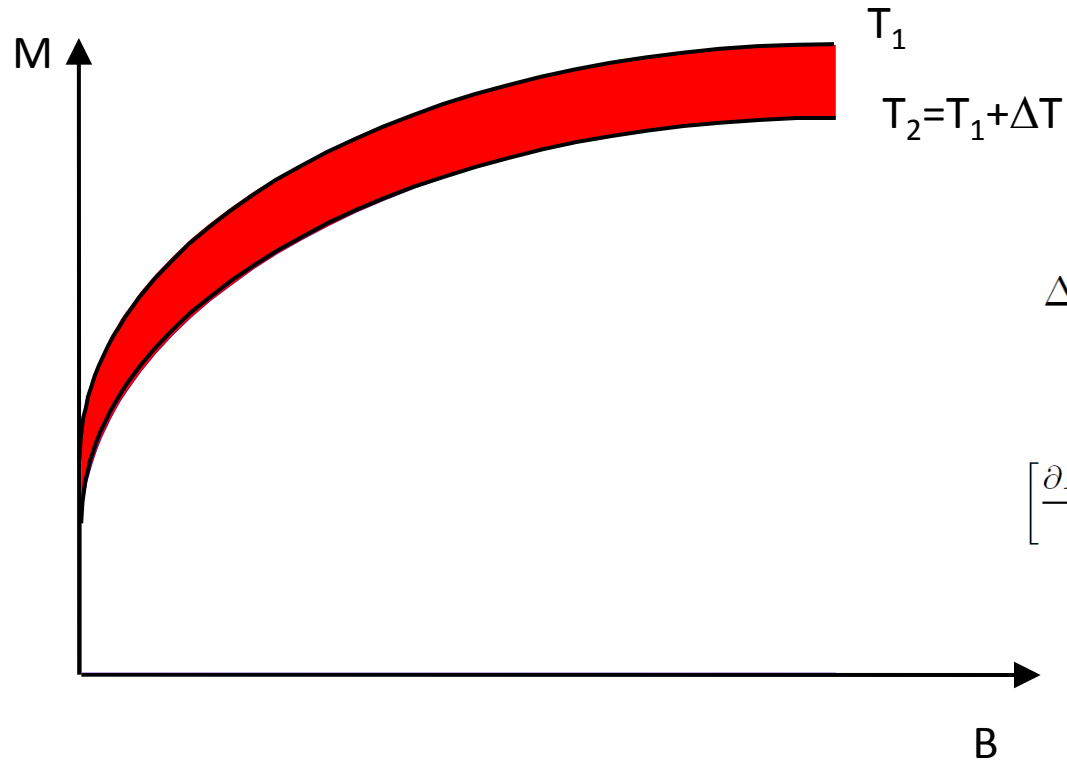
~~FIRST ORDER TRANSITION~~



SECOND ORDER TRANSITION

THERMODYNAMICS OF THE MCE

MAGNETIZATION CURVE



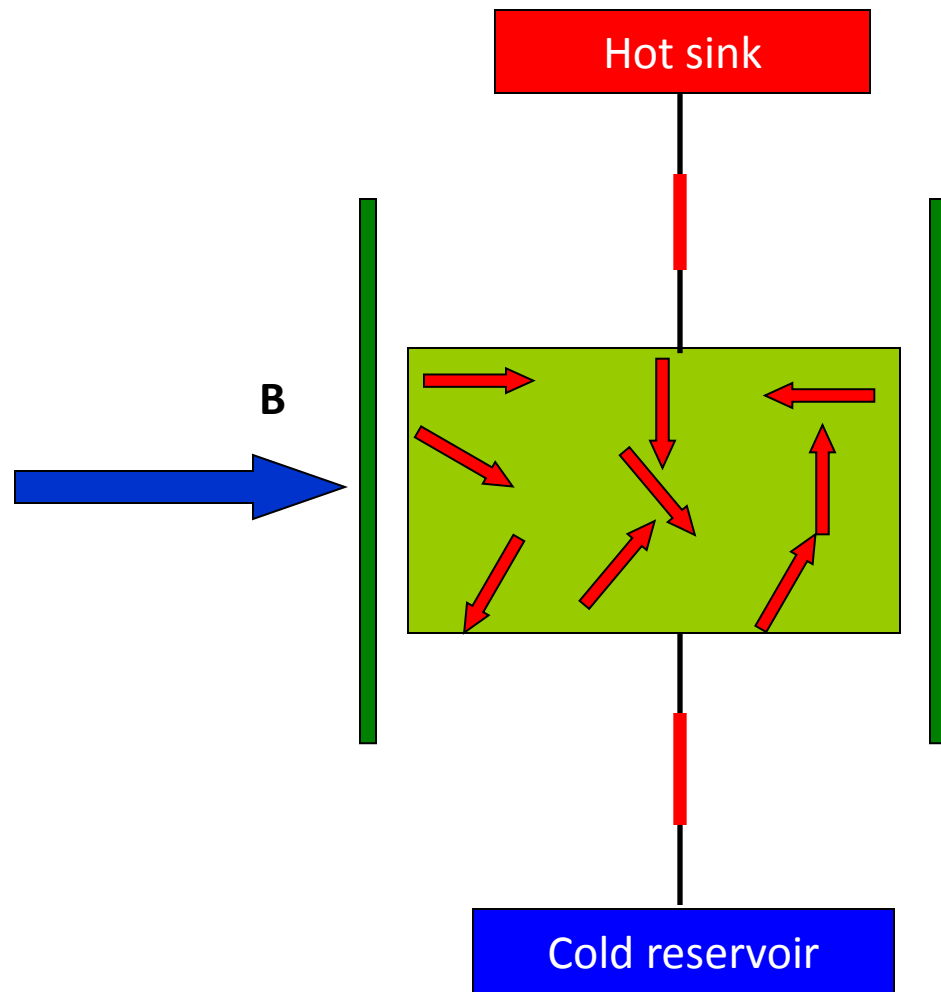
$$\Delta S_{iso}(T, \Delta B, P) = \int_{B_1}^{B_2} \left[\frac{\partial M(T, B, P)}{\partial T} \right]_{B,P} dB$$

$$\left[\frac{\partial M(T, B, P)}{\partial T} \right]_{B,P} = \lim_{\Delta T \rightarrow 0} \frac{1}{\Delta T} [M(T + \Delta T, B, P)dB - M(T, B, P)dB]$$

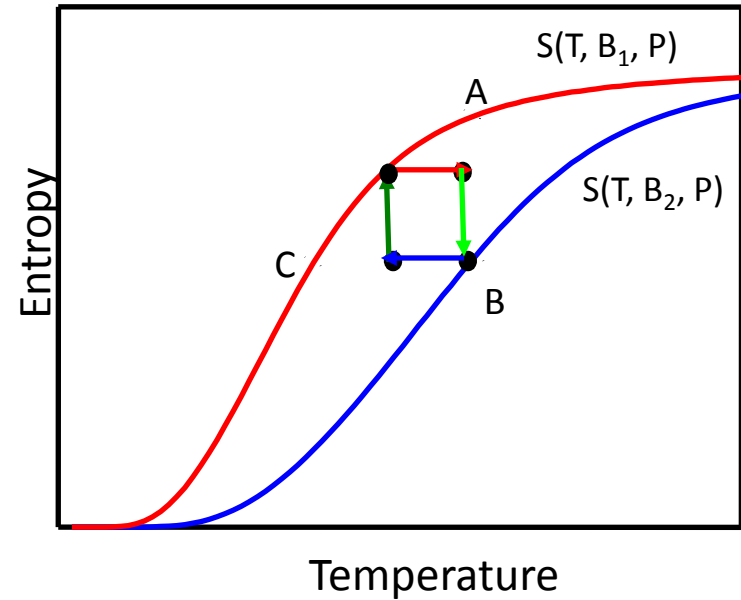
$$\Delta S_{iso}(T, \Delta B, P) = \frac{1}{\Delta T} \left[\int_{B_1}^{B_2} M(T + \Delta T, B, P)dB - \int_{B_1}^{B_2} M(T, B, P)dB \right]$$

MAGNETIC REFRIGERATOR

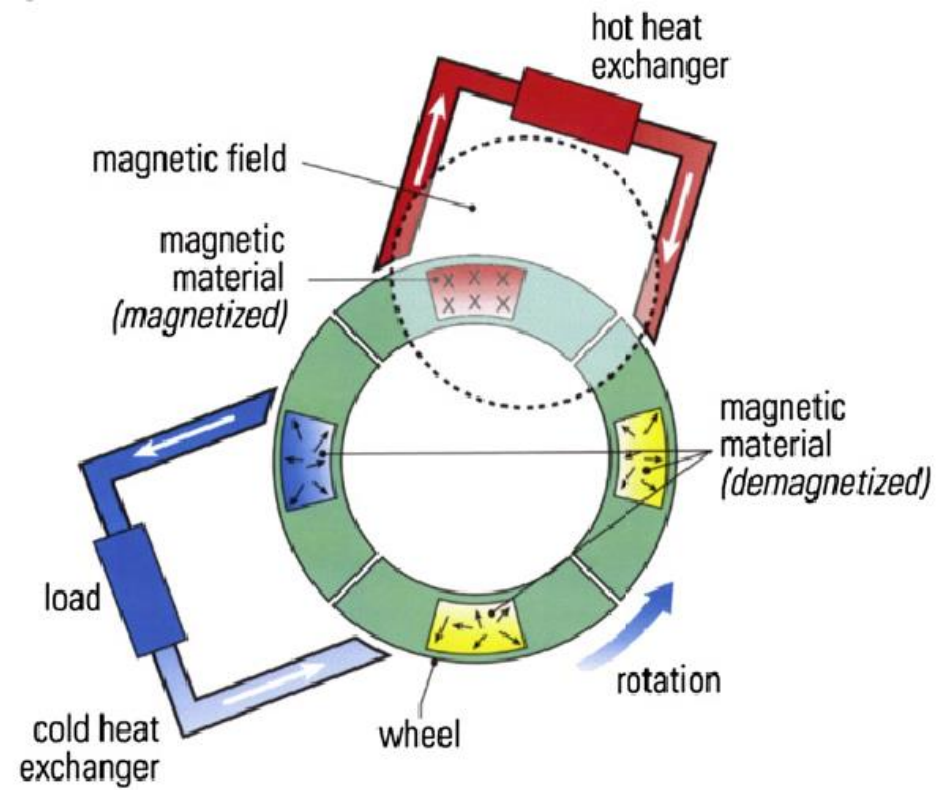
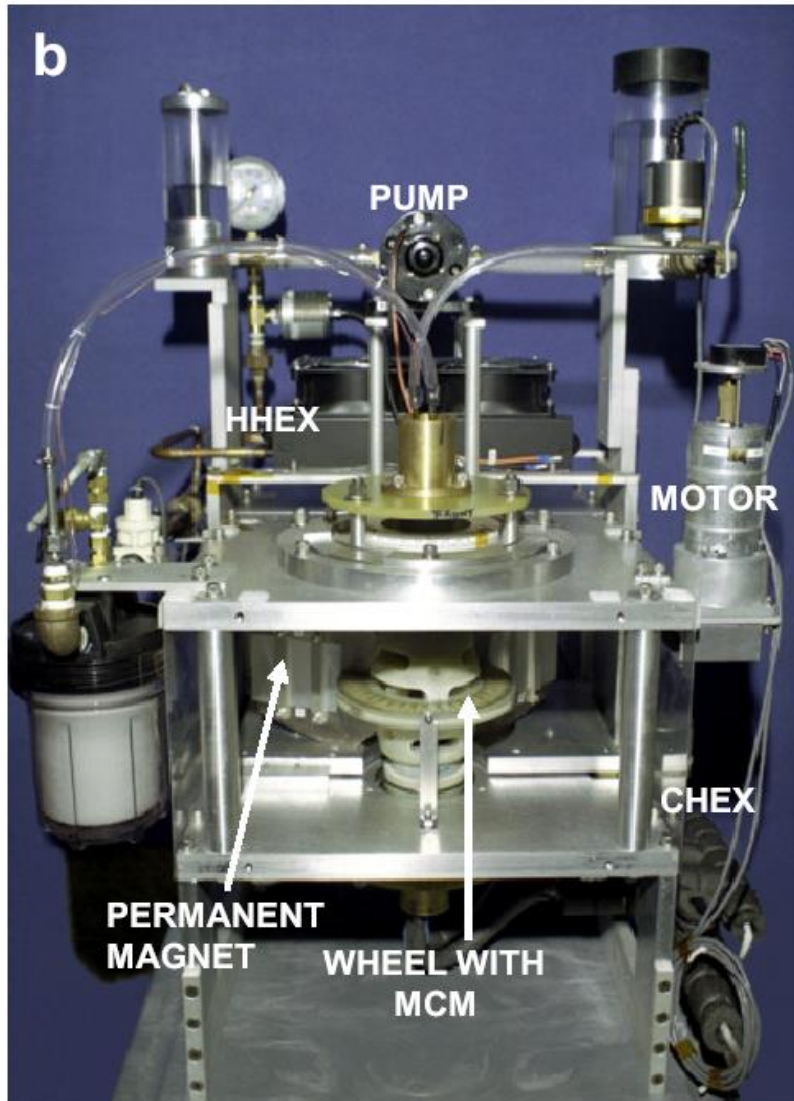
ACADEMIC PROTOTYPE



Heating (heat release) n



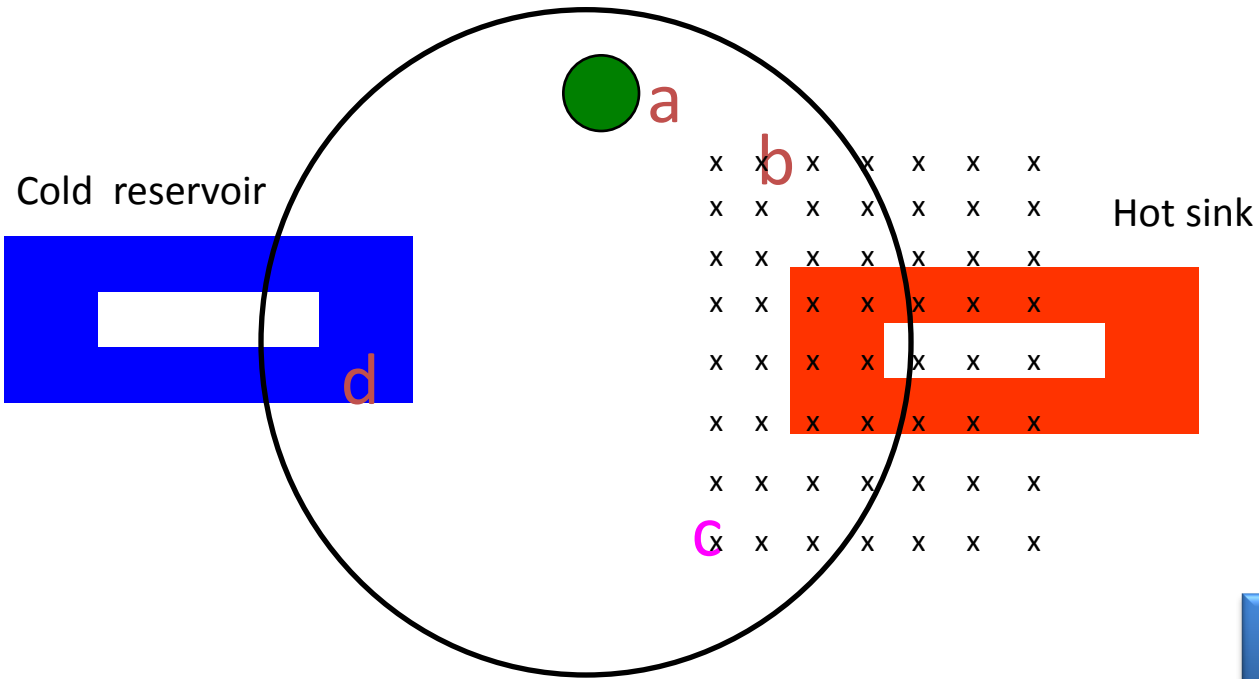
Carnot cycle



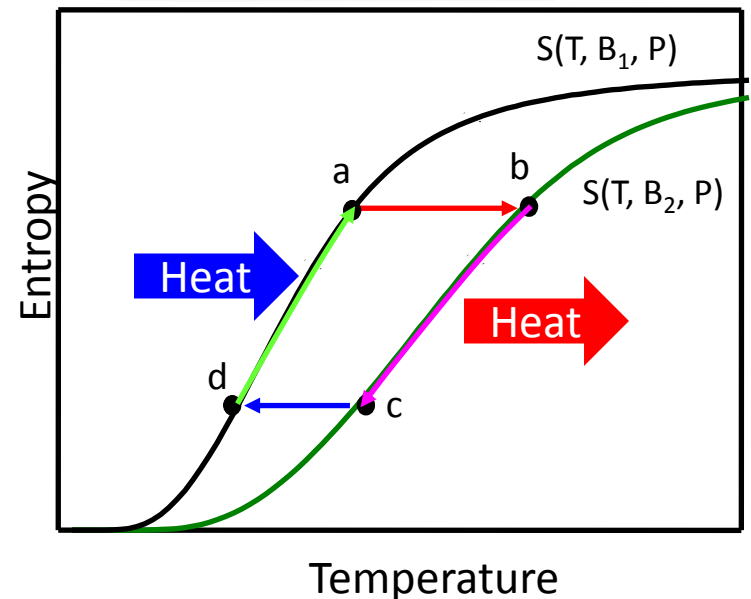
Environmental friendly

Large values of magnetic field

ACADEMIC PROTOTYPE

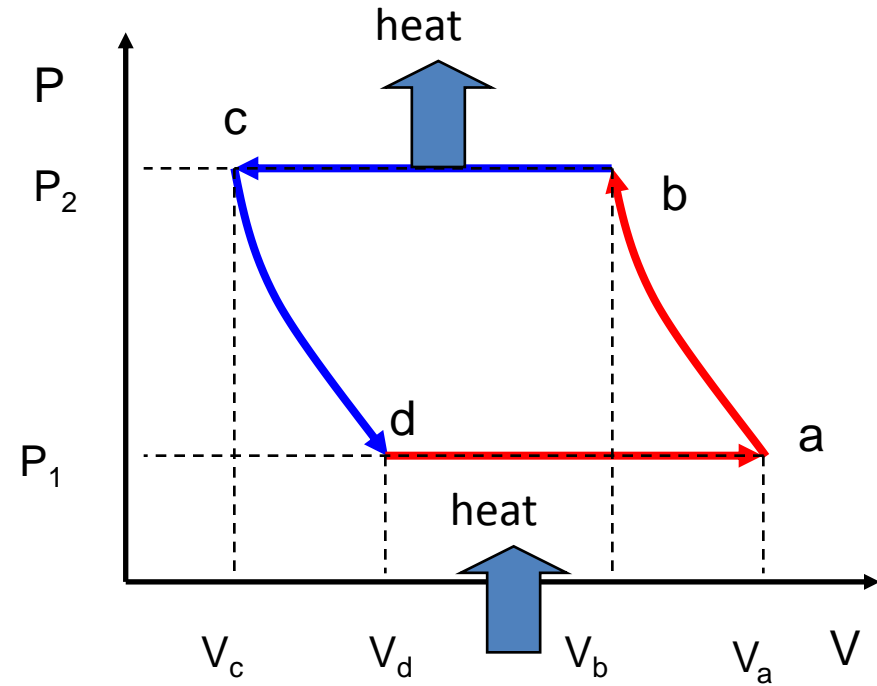


Brayton cycle

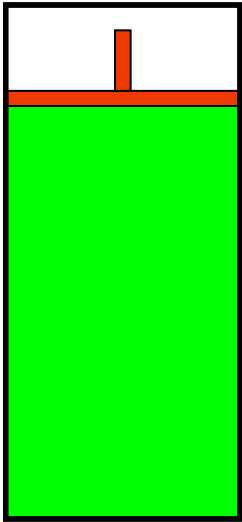


BAROCALORIC EFFECT

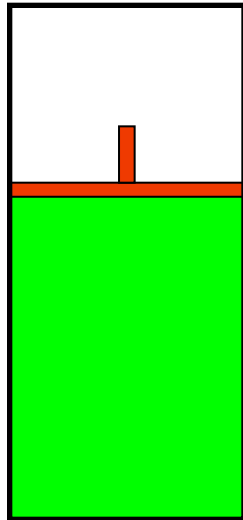
BAROCALORIC EFFECT



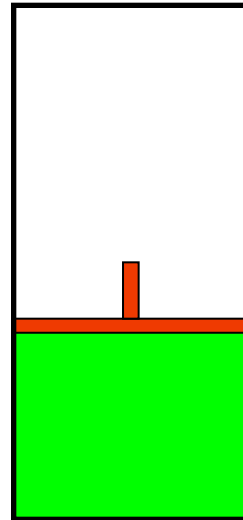
Adiabatic compression



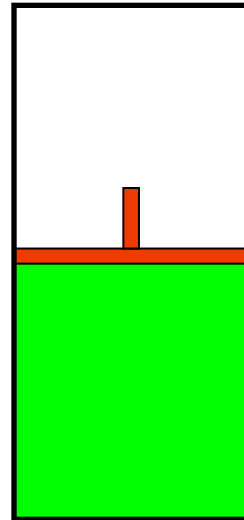
Isobaric compression



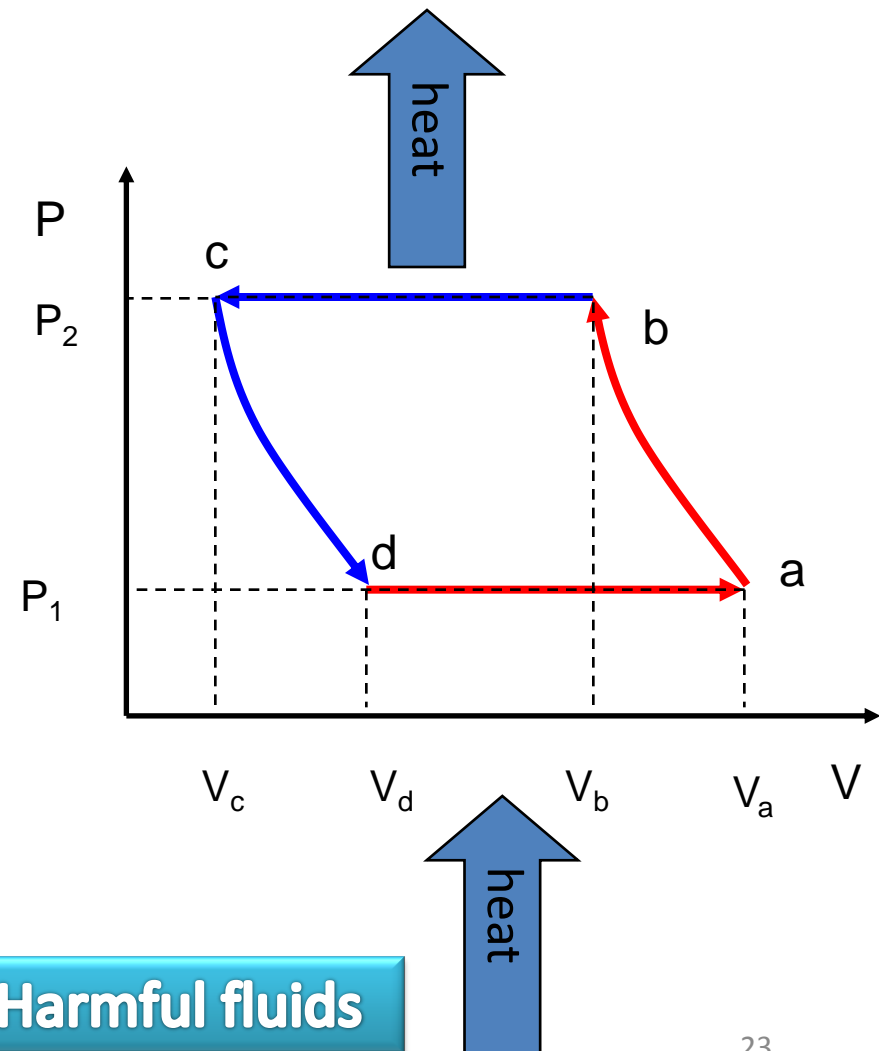
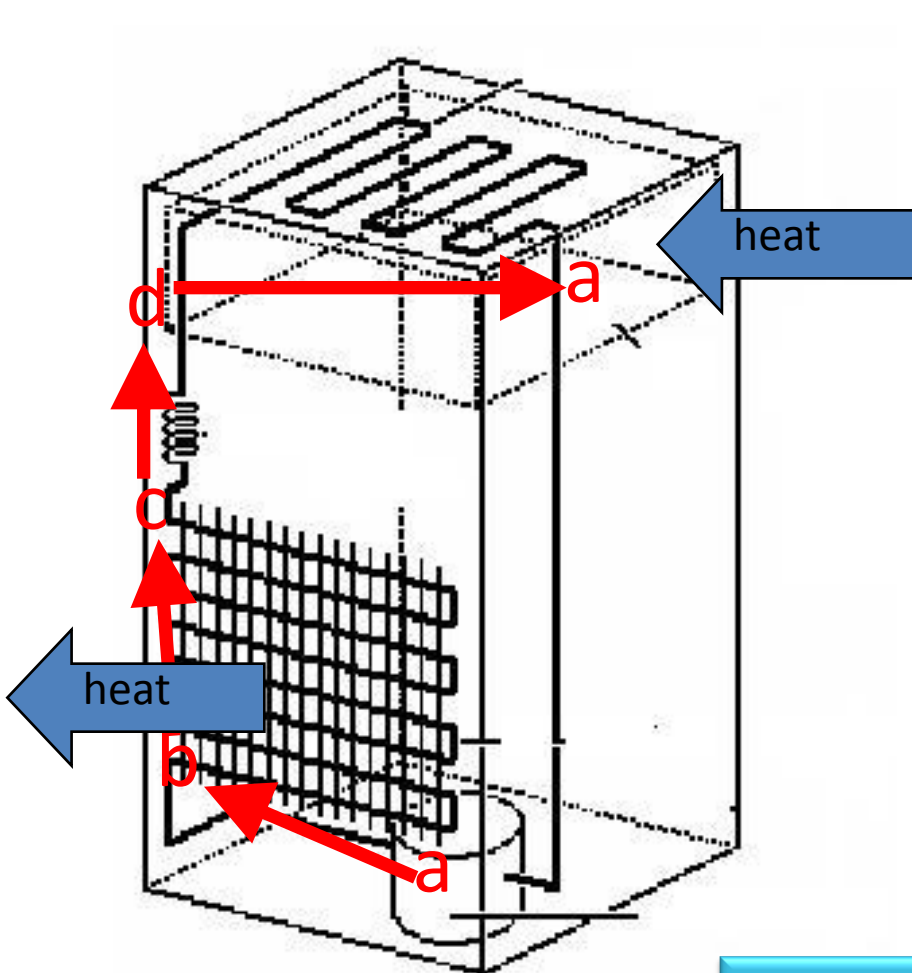
Adiabatic expansion



Isobaric expansion



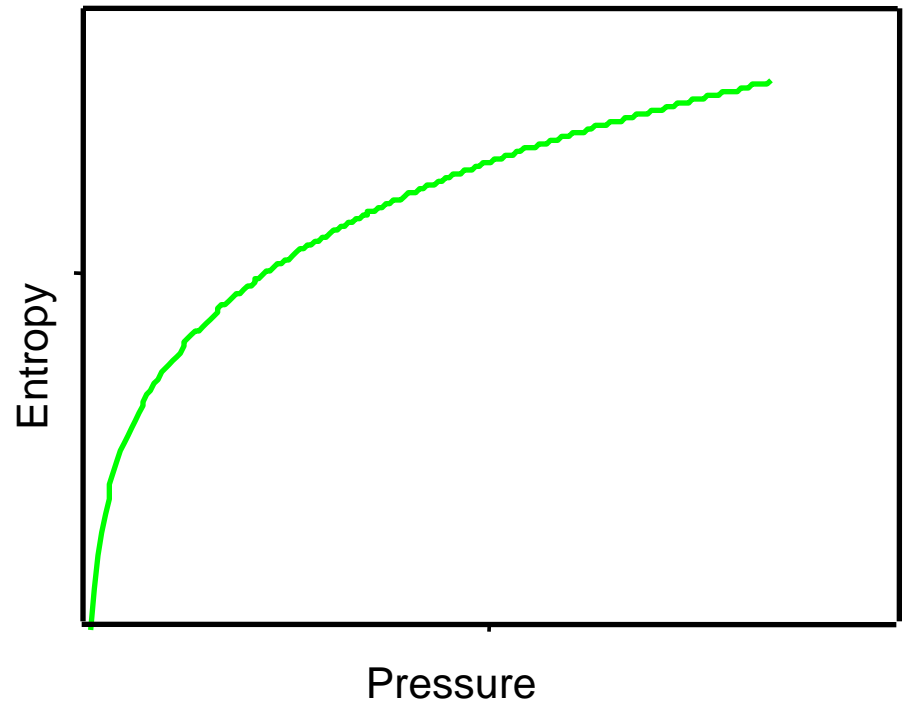
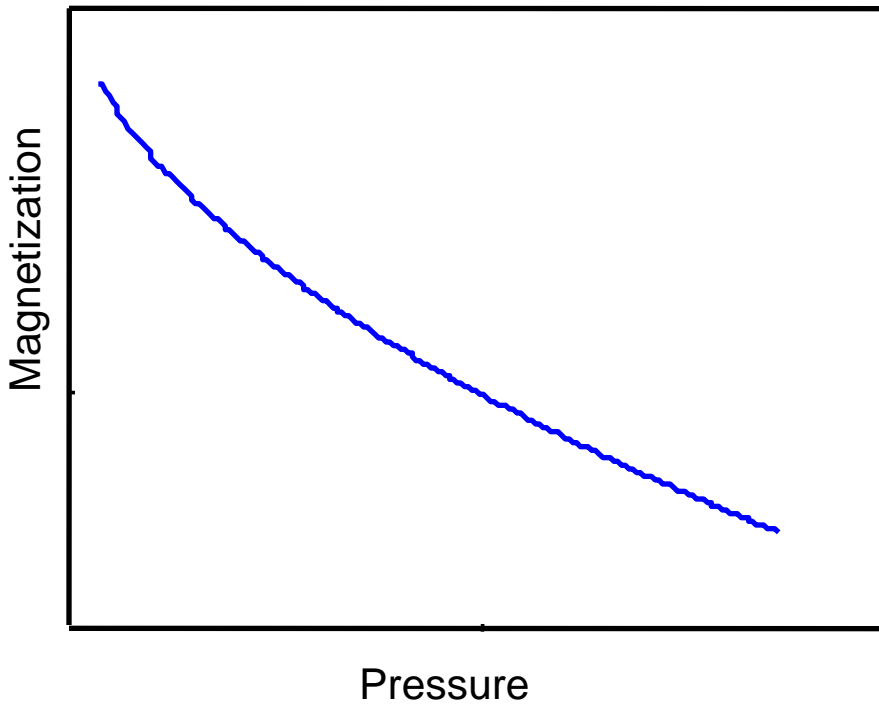
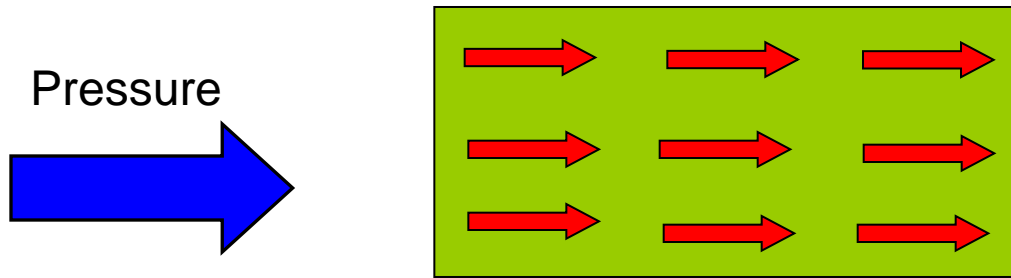
CONVENTIONAL REFRIGERATOR



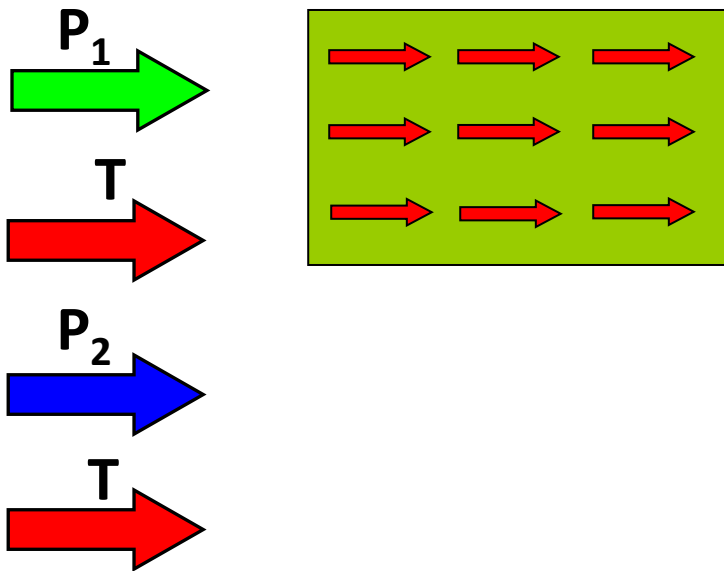
Drawback: Harmful fluids

MAGNETIC BAROCALORIC EFFECT

MAGNETIC BAROCALORIC EFFECT



MAGNETIC BAROCALORIC EFFECT

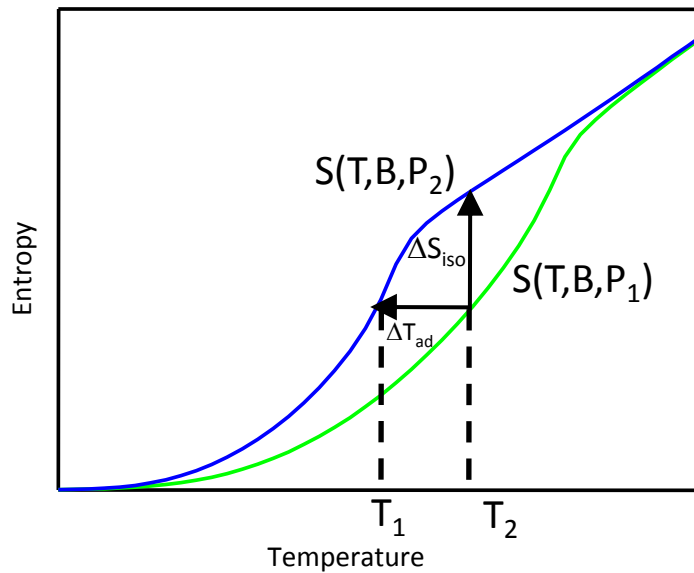


BAROCALORIC QUANTITIES

$$\Delta S_{iso}^{bar}(T, B, \Delta P) = S(T, B, P_2) - S(T, B, P_1)$$

$$\Delta T_{ad}^{bar}(T, B, \Delta P) = T_1 - T_2$$

$$S(T, B, P_2) = S(T, B, P_1)$$



THERMODYNAMICS OF THE BCE (ΔS_{iso}^{bar})

ENTROPY CHANGE

$$dS(T, B, P) = \left[\frac{\partial S(T, B, P)}{\partial T} + \frac{\delta S(T_C, B, P)}{\delta T} \right]_{B, P} dT + \left[\frac{\partial S(T, B, P)}{\partial B} + \frac{\delta S(T, B_C, P)}{\delta B} \right]_{T, P} dB + \left[\frac{\partial S(T, B, P)}{\partial P} + \frac{\delta S(T, B, P_C)}{\delta P} \right]_{T, B} dP$$

ISOTHERMAL PROCESS

ISOFIELD PROCESS

$$\Delta S_{iso}^{bar}(T, B, \Delta P) = \int_{P_1}^{P_2} \left[\frac{\partial S(T, B, P)}{\partial P} + \frac{\delta S(T, B_C, P)}{\delta P} \right]_{T, B} dP$$

SECOND ORDER TRANSITION

$$\Delta S_{iso}^{bar}(T, B, \Delta P) = \int_{P_1}^{P_2} \left[\frac{\partial S(T, B, P)}{\partial P} \right]_{T, B} dP$$

THERMODYNAMICS OF THE BCE (ΔT_{ad}^{bar})

ENTROPY CHANGE

$$0 = \left[\frac{\partial S(T, B, P)}{\partial T} + \frac{\delta S(T_C, B, P)}{\delta T} \right]_{B, P} dT + \left[\frac{\partial S(T, B, P)}{\partial B} + \frac{\delta S(T, B_C, P)}{\delta B} \right]_{T, P} dB + \left[\frac{\partial S(T, B, P)}{\partial P} + \frac{\delta S(T, B, P_C)}{\delta P} \right]_{T, B} dP$$

ADIABATIC PROCESS

ISOFIELD PROCESS

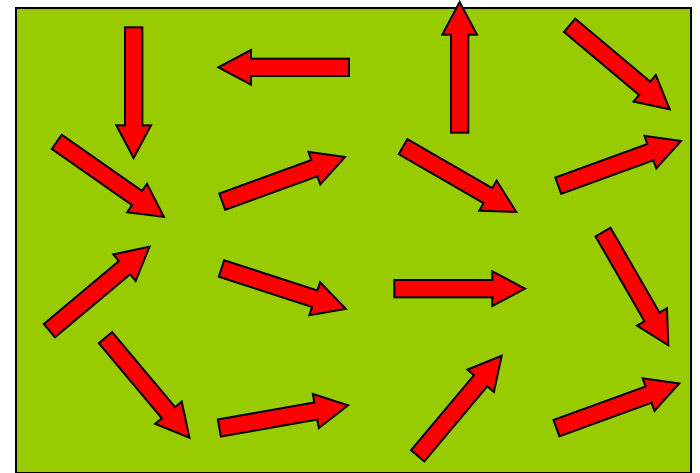
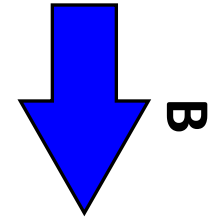
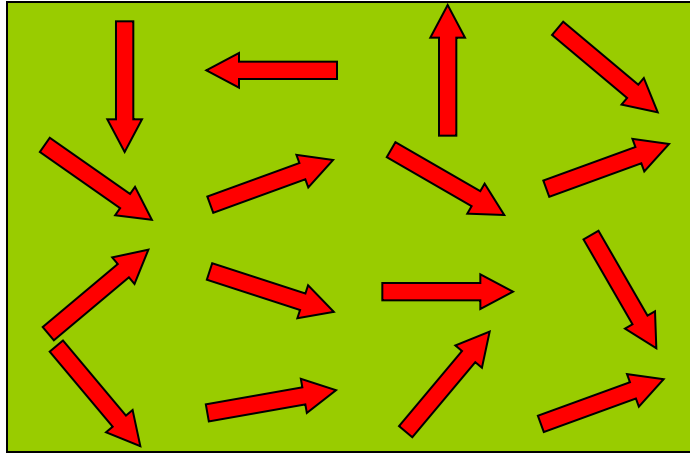
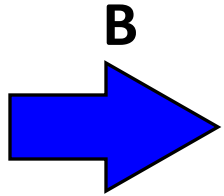
$$\Delta T_{ad}^{bar}(T, B, \Delta P) = - \int_{P_1}^{P_2} \frac{1}{C(T, B, P)} \left[\frac{\partial S(T, B, P)}{\partial P} + \frac{\delta S(T, B, P_C)}{\delta P} \right]_{T, B} dP$$

SECOND ORDER TRANSITION

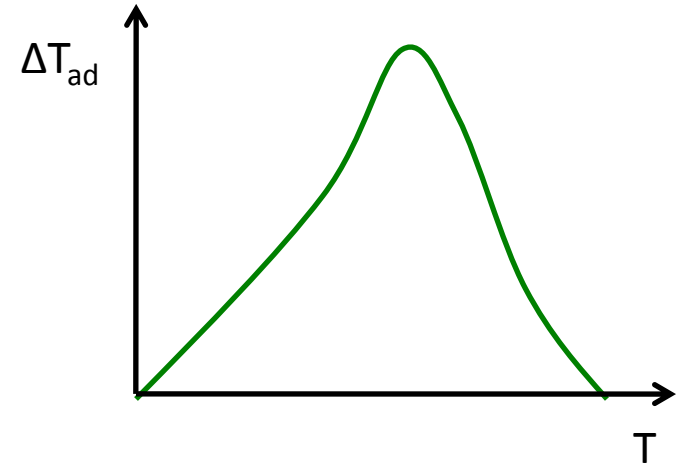
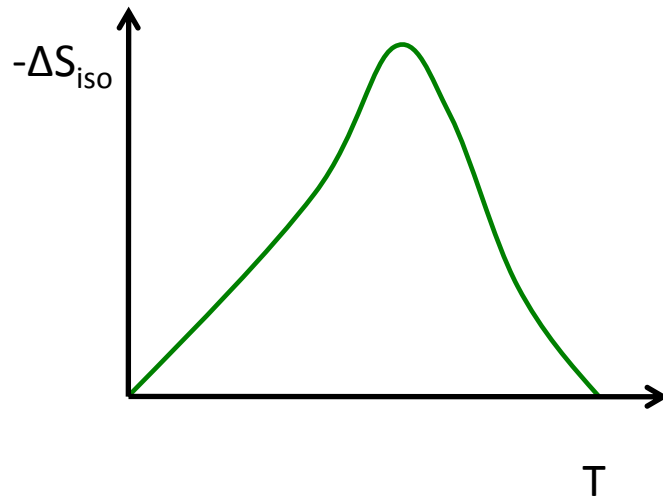
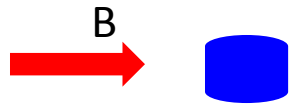
$$\Delta T_{ad}^{bar}(T, B, \Delta P) = - \int_{P_1}^{P_2} \frac{1}{C(T, B, P)} \left[\frac{\partial S(T, B, P)}{\partial P} \right]_{T, B} dP$$

ANISOTROPIC MAGNETOCALORIC EFFECT

ANISOTROPIC MCE



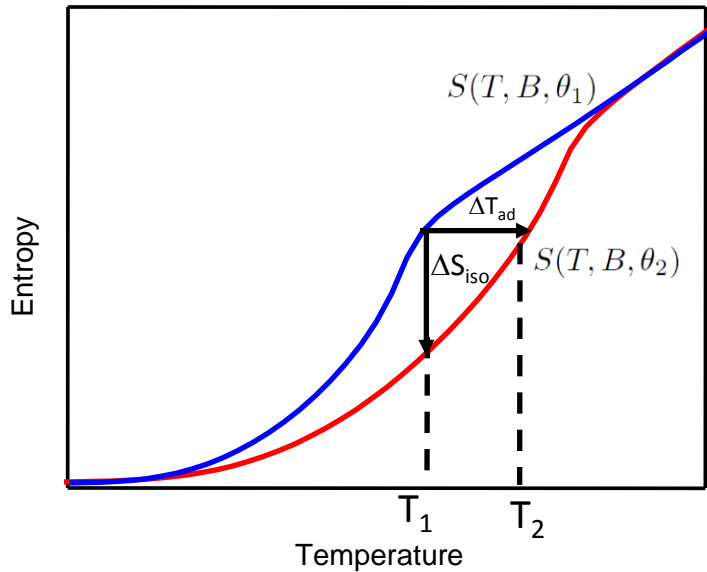
ANISOTROPIC MCE



Niktin et al, Phys. Rev. Lett. 105 (2010)137205

von Ranke et al, J. Appl. Phys. 104 (2008)093906

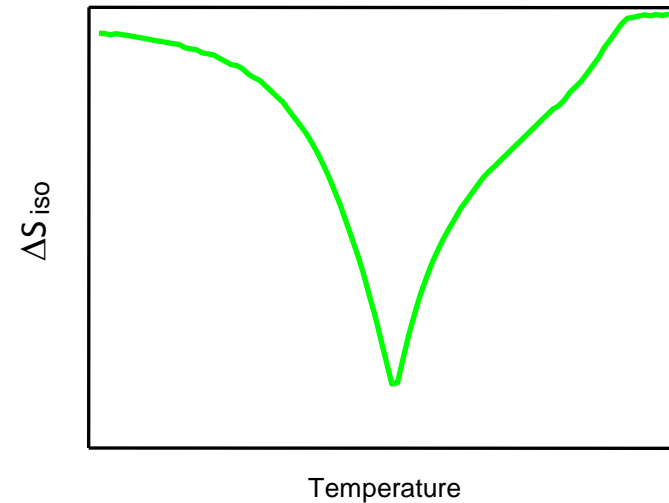
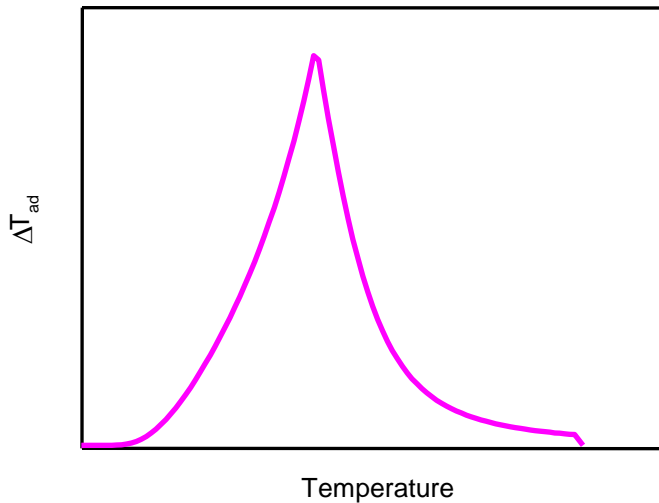
ANISOTROPIC MAGNETOCALORIC EFFECT



$$\Delta S_{iso}^{ani}(T, B, \Delta\theta) = S(T, B, \theta_2) - S(T, B, \theta_1)$$

$$\Delta T_{ad}^{ani}(T, B, \Delta\theta) = T_2 - T_1$$

$$S(T, B, \theta_2) = S(T, B, \theta_1)$$



THERMODYNAMICS OF THE AMCE (ΔS_{iso}^{ani})

ENTROPY CHANGE

$$dS(T, B, \theta) = \left[\frac{\partial S(T, B, \theta)}{\partial T} + \frac{\delta S(T_C, B, \theta)}{\delta T} \right]_{B, \theta} dT + \left[\frac{\partial S(T, B, \theta)}{\partial B} + \frac{\delta S(T, B_C, \theta)}{\delta B} \right]_{T, \theta} dB + \left[\frac{\partial S(T, B, \theta)}{\partial \theta} + \frac{\delta \theta(T, B, \theta_C)}{\delta \theta} \right]_{T, B} d\theta$$

ISOTHERMAL PROCESS

ISOFIELD PROCESS

$$\Delta S_{iso}^{ani}(T, B, \Delta\theta) = \int_{\theta_1}^{\theta_2} \left[\frac{\partial S(T, B, \theta)}{\partial \theta} + \frac{\delta S(T, B, \theta_C)}{\delta \theta} \right]_{T, B} d\theta$$

SECOND ORDER TRANSITION

$$\Delta S_{iso}^{ani}(T, B, \Delta\theta) = \int_{\theta_1}^{\theta_2} \left[\frac{\partial S(T, B, \theta)}{\partial \theta} \right]_{T, B} d\theta$$

THERMODYNAMICS OF THE AMCE (ΔT_{ad}^{ani})

ENTROPY CHANGE

$$0 = \left[\frac{\partial S(T, B, \theta)}{\partial T} + \frac{\delta S(T_C, B, \theta)}{\delta T} \right]_{B, \theta} dT + \left[\frac{\partial S(T, B, \theta)}{\partial B} + \frac{\delta S(T, B_C, \theta)}{\delta B} \right]_{T, \theta} dB + \left[\frac{\partial S(T, B, \theta)}{\partial \theta} + \frac{\delta \theta(T, B, \theta_C)}{\delta \theta} \right]_{T, B} d\theta$$

ADIABATIC PROCESS

ISOFIELD PROCESS

$$\Delta T_{ad}^{ani}(T, B, \Delta\theta) = - \int_{\theta_1}^{\theta_2} \frac{1}{C(T, B, \theta)} \left[\frac{\partial S(T, B, \theta)}{\partial \theta} + \frac{\delta S(T, B, \theta_C)}{\delta \theta} \right]_{T, B} d\theta$$

SECOND ORDER TRANSITION

$$\Delta T_{ad}^{ani}(T, B, \Delta\theta) = - \int_{\theta_1}^{\theta_2} \frac{1}{C(T, B, \theta)} \left[\frac{\partial S(T, B, \theta)}{\partial \theta} \right]_{T, B} d\theta$$

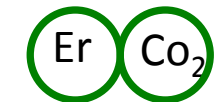
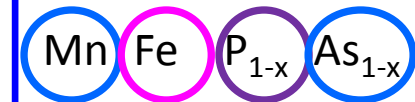
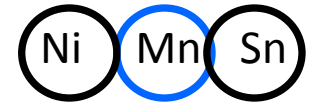
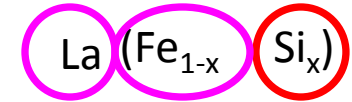
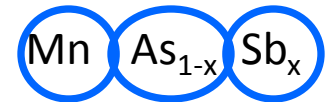
MAGNETIC MATERIALS

MAGNETIC MATERIALS

Periodic Table of the Elements

- hydrogen
- alkali metals
- alkali earth metals
- transition metals
- poor metals
- nonmetals
- noble gases
- rare earth metals

1	H	2	He																																
3	Li	4	Be																																
5	B	6	C	7	N	8	O	9	F	10	Ne																								
11	Na	12	Mg	13	Al	14	Si	15	P	16	S	17	Cl	18	Ar																				
19	K	20	Ca	21	Sc	22	Ti	23	V	24	Cr	25	Mn	26	Fe	27	Co	28	Ni	29	Cu	30	Zn	31	Ga	32	Ge	33	As	34	Se	35	Br	36	Kr
37	Rb	38	Sr	39	Y	40	Zr	41	Nb	42	Mo	43	Tc	44	Ru	45	Rh	46	Pd	47	Ag	48	Cd	49	In	50	Sn	51	Sb	52	Te	53	I	54	Xe
55	Cs	56	Ba	57	La	72	Hf	73	Ta	74	W	75	Re	76	Os	77	Ir	78	Pt	79	Au	80	Hg	81	Tl	82	Pb	83	Bi	84	Po	85	At	86	Rn
87	Fr	88	Ra	89	Ac	104	Unq	105	Unp	106	Unh	107	Uns	108	Uno	109	Une	110	Uun																
58	Ce	59	Pr	60	Nd	61	Pm	62	Sm	63	Eu	64	Gd	65	Tb	66	Dy	67	Ho	68	Er	69	Tm	70	Yb	71	Lu								
90	Th	91	Pa	92	U	93	Np	94	Pu	95	Am	96	Cm	97	Bk	98	Cf	99	Es	100	Fm	101	Md	102	No	103	Lr								



A. M. Tishin and Y. Spickin, The magnetocaloric effect and its application.

Gschneidner et al, Rep. Prog. Phys 68(2005)1479.

THEORY AND CALCULATIONS

GENERAL INTRODUCTION

THERMODYNAMIC VIEW

SYSTEMS OF LOCALIZED MAGNETIC MOMENTS

APROXIMATIONS

CALORIC QUANTITIES

MAGNETISM

Localized

4f Spins - Localized levels

Rare earth metals and alloys

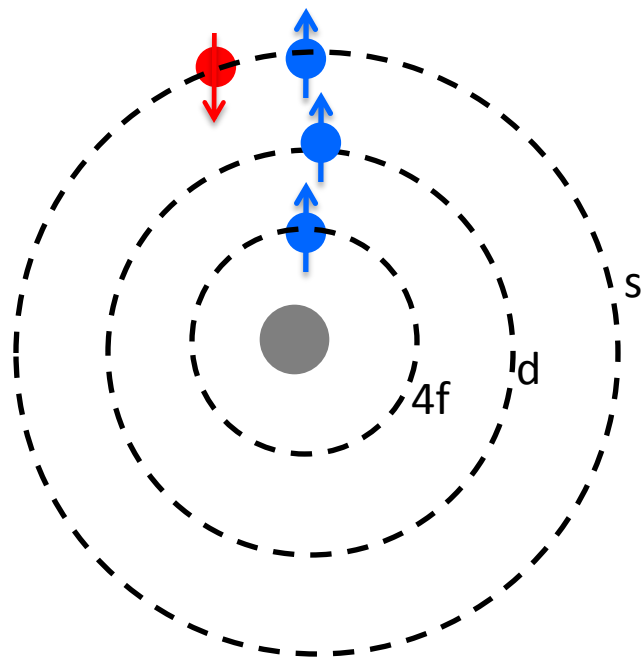
Itinerant

3d electrons - Energy band

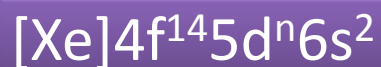
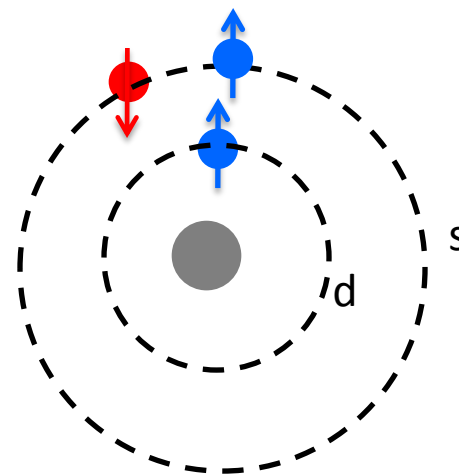
Transition metals and alloys

ATOMIC MAGNETISM

Rare earth



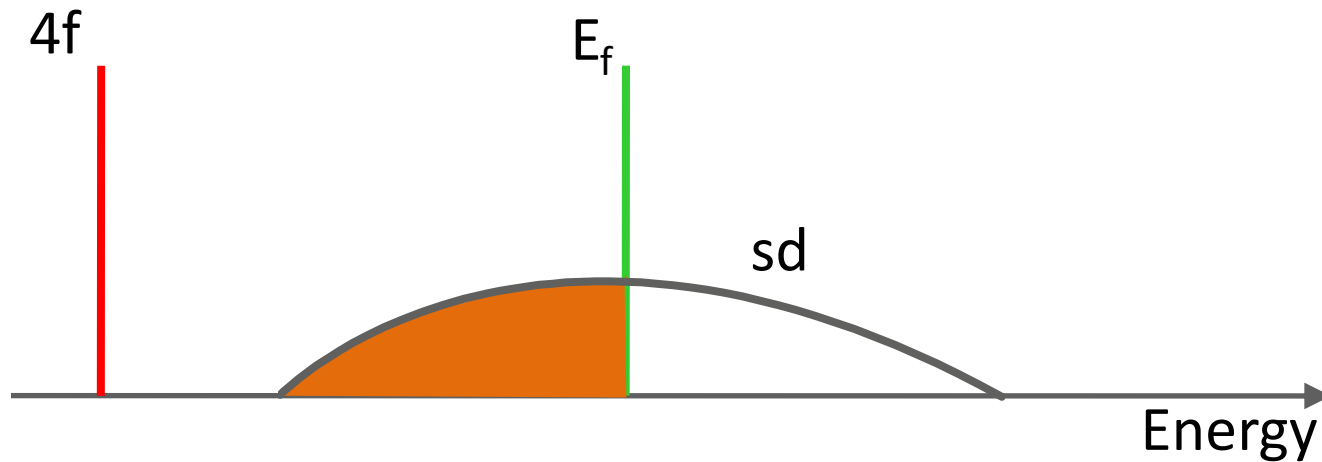
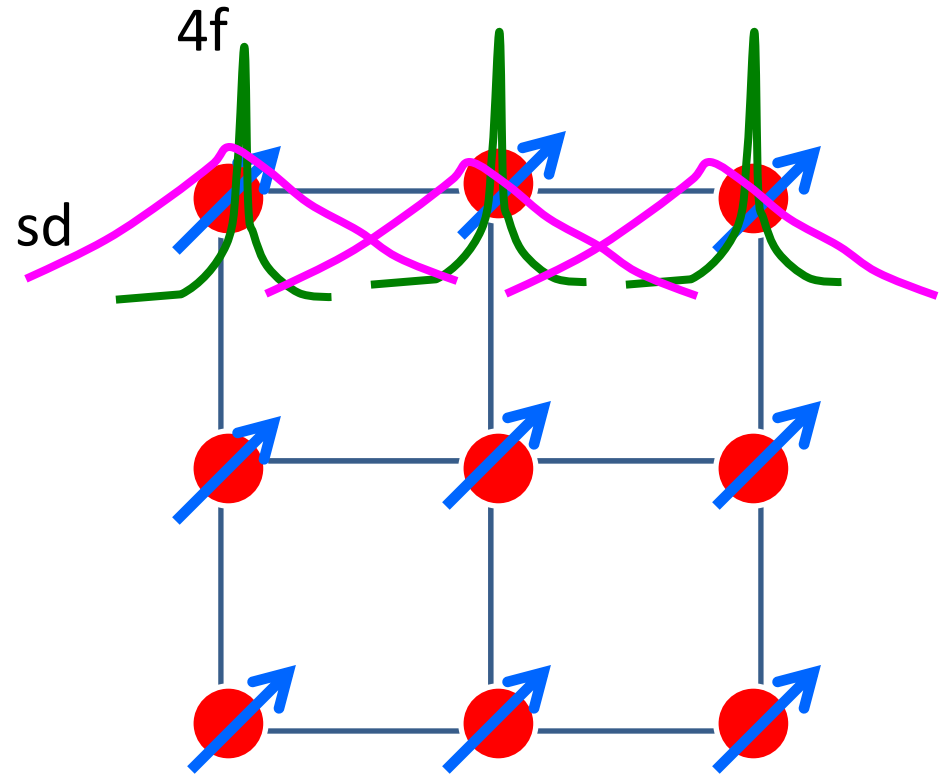
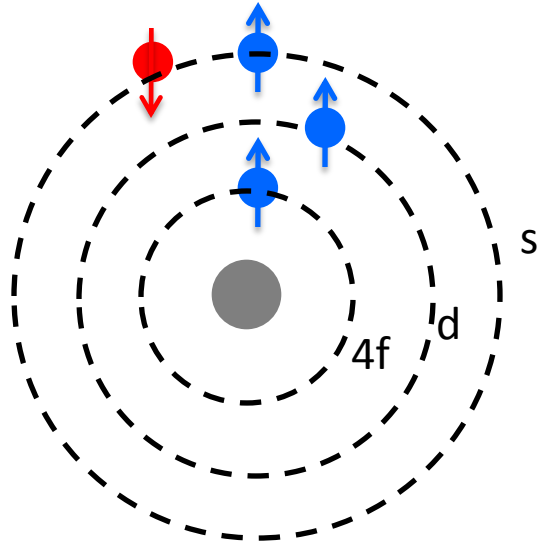
Transition metals





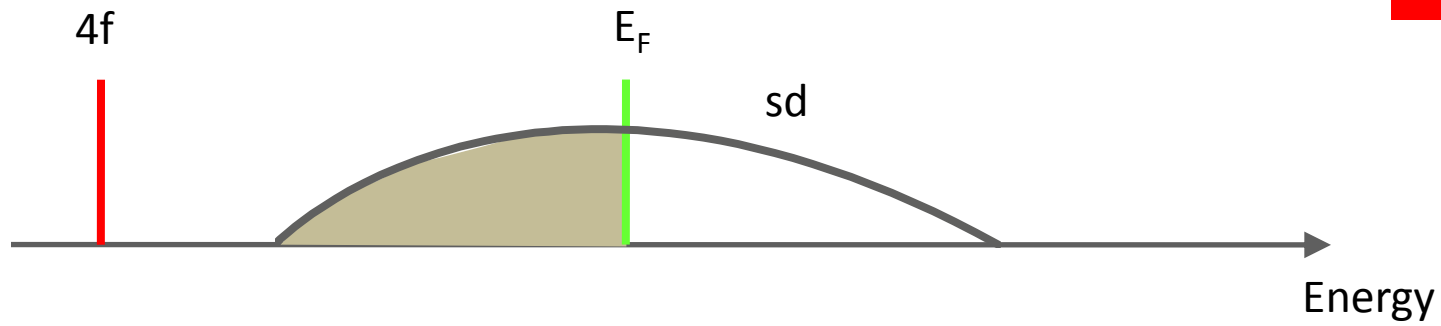
RARE EARTH METALS AND THEIR ALLOYS

RARE EARTH METALS

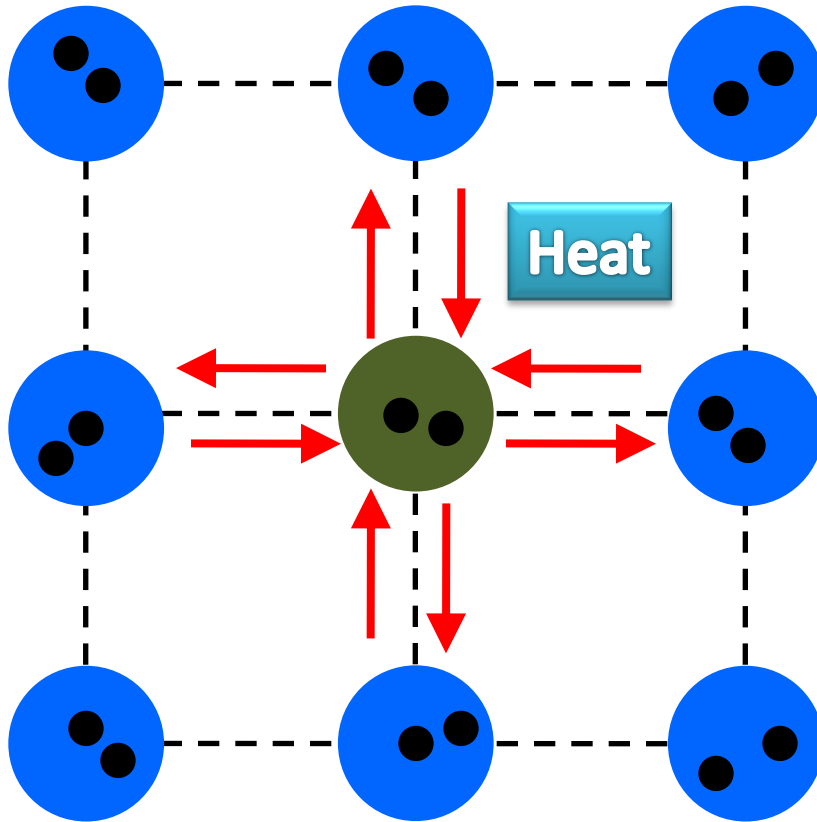


REGIMES

RKKY



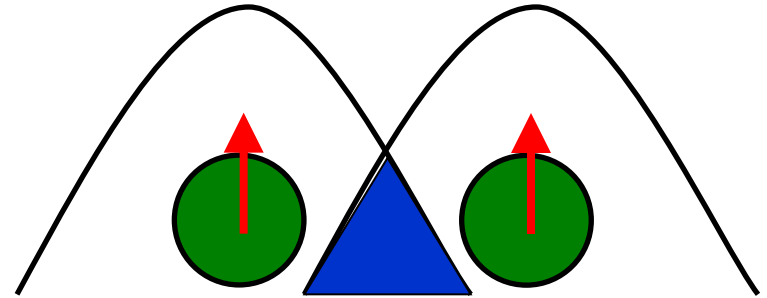
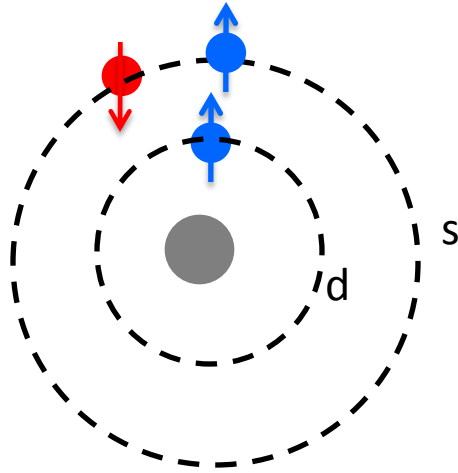
CANONICAL ENSEMBLE



$$Z = \sum_{i=1}^n e^{-\beta E_i}$$

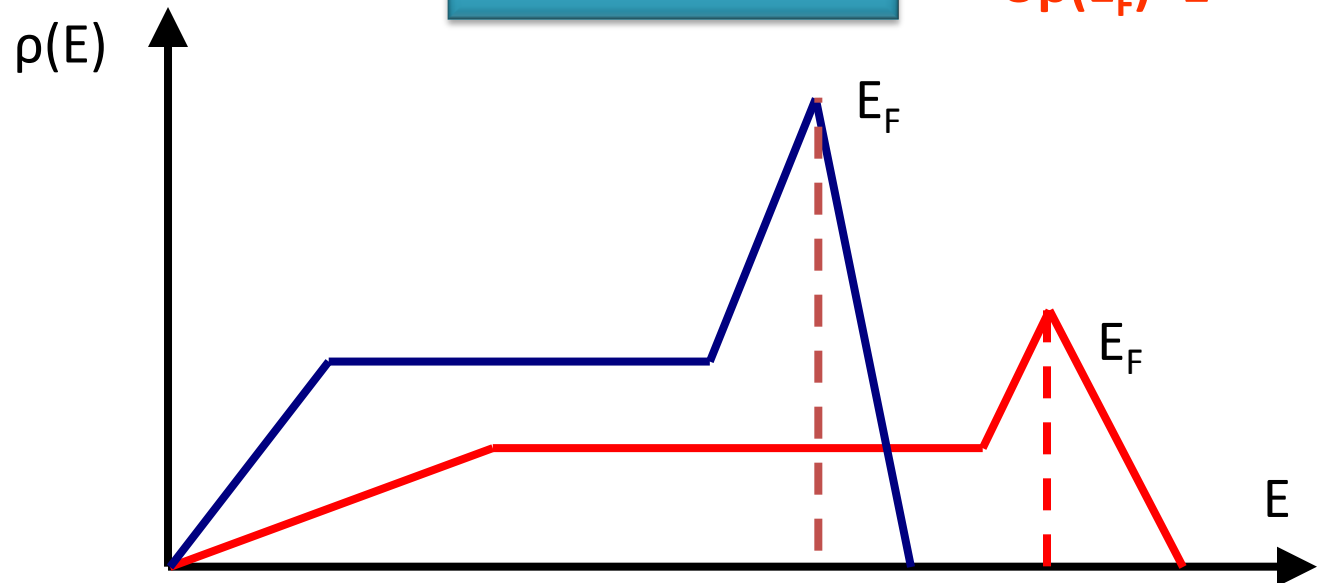
TRANSITION METALS AND THEIR ALLOYS

ITINERANT MAGNETISM

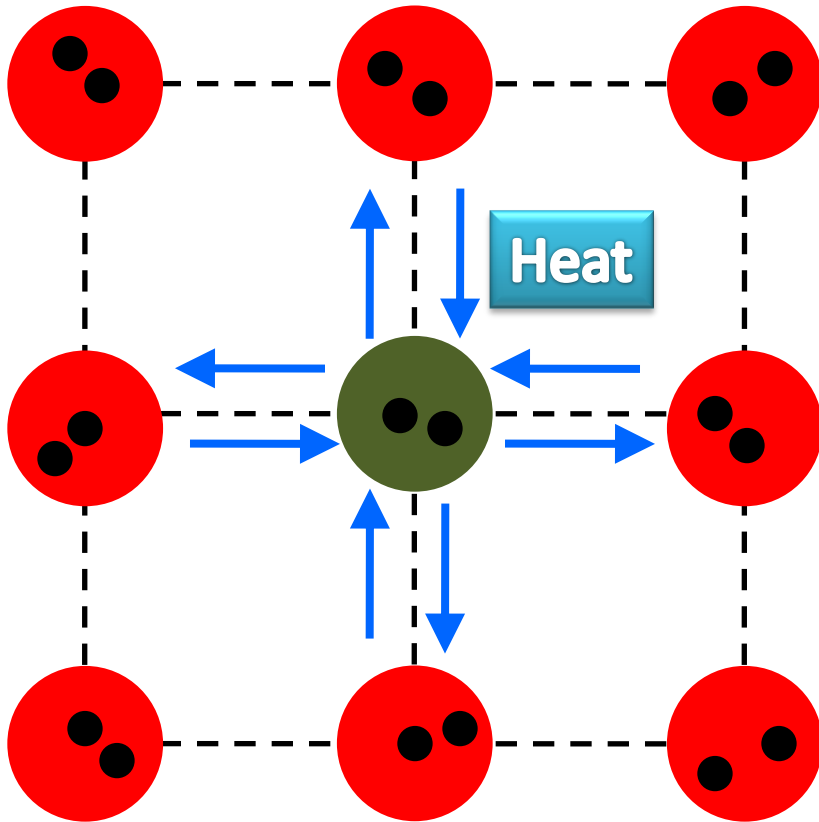


Stoner criterion

$$U\rho(E_F) > 1$$



GRAN-CANONICAL ENSEMBLE



$$Z = \prod_{k\sigma} \{1 + \exp[-\beta(\varepsilon_{k\sigma} - \mu)]\}$$

RARE EARTH METALS AND ALLOYS

MICROSCOPIC DESCRIPTION

HAMILTONIAN

Lattice

$$\mathcal{H}_{lat} = \sum_q \hbar\omega_q a_q^\dagger a_q$$

Non magnetic electrons (spd)

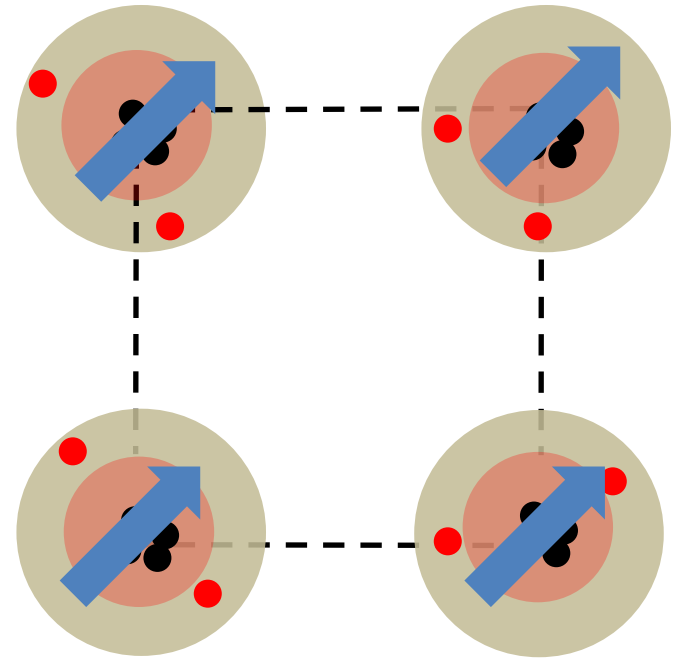
$$\mathcal{H}_{el}^{spd} = \sum_k \varepsilon_k c_k^\dagger c_k$$

Magnetic electrons (4f)

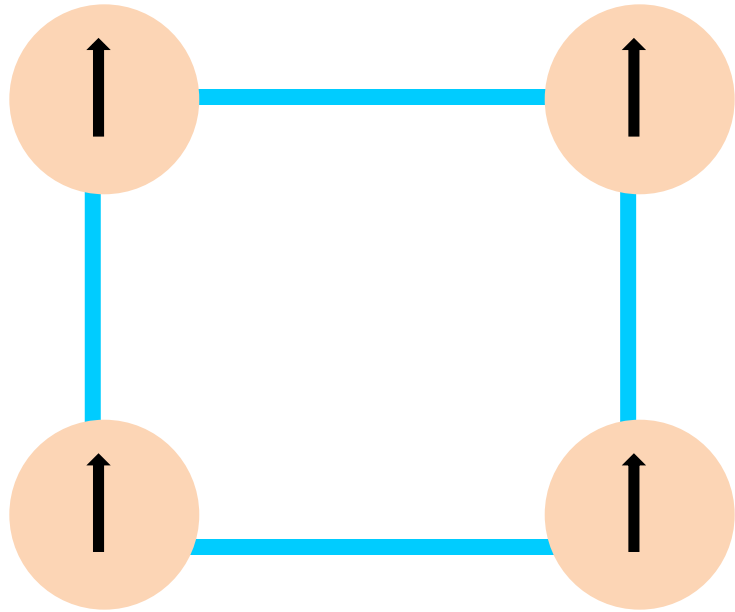
$$\mathcal{H}_{mag}^{4f} = - \sum_{i,j} \mathcal{J}_{ij}(r) \vec{J}_i \cdot \vec{J}_j - \sum_i g\mu_B \vec{B} \cdot \vec{J}_i$$

Total

$$\mathcal{H} = \mathcal{H}_{lat} + \mathcal{H}_{el}^{spd} + \mathcal{H}_{mag}^{4f}$$



MAGNETIC HAMILTONIAN



$$\mathcal{H}_{mag}^{4f} = - \sum_{i,j} \mathcal{J}_{ij}(r) \vec{J}_i \cdot \vec{J}_j - \sum_i g\mu_B \vec{B} \cdot \vec{J}_i$$

Approximation: Translation invariance

$$\mathcal{H}_{mag}^{4f} = - \sum_{ij} \tilde{\mathcal{J}}_0 J_i \cdot J_j - \sum_i g\mu_B B J_i$$

MEAN FIELD THEORY

Hamiltonian

$$\mathcal{H}_{mag}^{4f} = - \sum_{i,j} \mathcal{J}_0 \vec{J}_i \cdot \vec{J}_j - \sum_i g\mu_B \vec{B} \cdot \vec{J}_i$$

Mean field

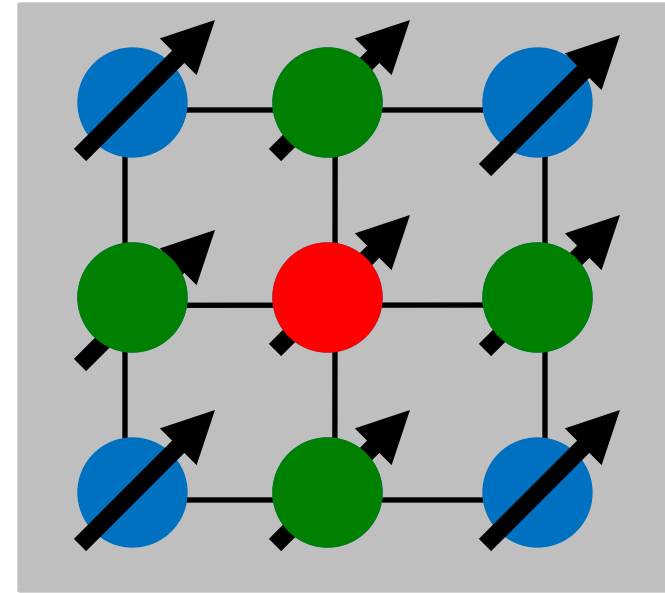
Molecular field

$$\mathcal{H}_{mag}^{4f} = -\mathcal{J}_0 \sum_i \vec{J}_i \cdot \left[\sum_j \langle \vec{J}_j \rangle \right] - \sum_i g\mu_B \vec{B} \cdot \vec{J}_i$$

Isotropic system

$$\mathcal{H}_{mag}^{4f} = - \sum_i g\mu_B B^{eff} \cdot J_i^z$$

$$B^{eff} = B + \frac{\mathcal{J}_0 \langle J^z \rangle}{g\mu_B}$$



MEAN FIELD THEORY

Equation of motion

$$\mathcal{H}_{mag}^{4f} |\psi\rangle = E |\psi\rangle$$

$$\mathcal{H}_{mag}^{4f} = - \sum_i g\mu_B B^{eff} \cdot J_i^z$$

Mean field

$$-g\mu_B B^{eff} \cdot \left(J_i^z |\psi\rangle \right) = E |\psi\rangle$$

$$\left(J_i^z |\psi\rangle = m |\psi\rangle \right)$$

Energy

$$E_m = -g\mu_B B^{eff} m$$

$$-J \leq m \leq J$$

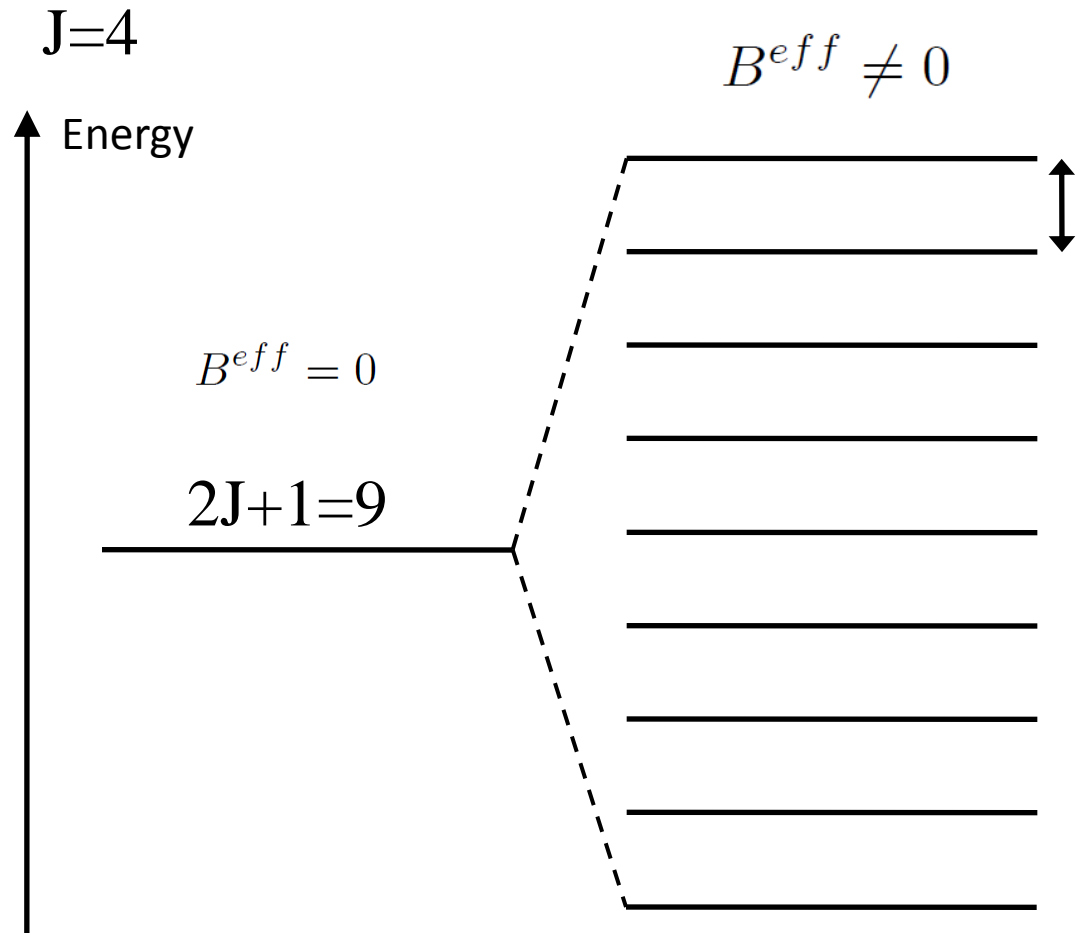
$$B^{eff} = B + \frac{\mathcal{J}_0 \langle J^z \rangle}{g\mu_B}$$

ENERGY LEVELS

$$E_m = -g\mu_B B^{eff} m$$

$$B^{eff} = B + \frac{\mathcal{J}_0 \langle J^z \rangle}{g\mu_B}$$

$$-J \leq m \leq J$$



MAGNETIZATION

Partition function

$$Z_{mag}^{4f}(T, B, P) = \sum_{m=-J}^{m=J} e^{-\beta E_m}$$

Magnetic free energy

$$F_{mag}^{4f}(T, B, P) = -k_B T \ln \left[\sum_{m=-J}^{m=J} e^{-\beta E_m} \right]$$

$$E_m = -g\mu_B B^{eff} m$$

Magnetization

$$M(T, B, P) = - \left[\frac{\partial F_{mag}^{4f}(T, B, P)}{\partial B^{eff}} \right]_{T, P}$$

$$M(T, B, P) = \frac{\sum_{m=-J}^{m=J} [\partial E_m / \partial B^{eff}] e^{-\beta E_m}}{\sum_{m=-J}^{m=J} e^{-\beta E_m}}$$

$$B^{eff} = B + \frac{\mathcal{J}_0 \langle J^z \rangle}{g\mu_B}$$

Self-consistency

FREE ENERGY AND MAGNETIC ENTROPY

Magnetic free energy

$$F_{mag}^{4f}(T, B, P) = -k_B T \ln \left[\sum_{m=-J}^{m=J} e^{-\beta E_m} \right]$$

Magnetic entropy

$$S_{mag}^{4f}(T, B, P) = - \left[\frac{\partial F_{mag}^{4f}(T, B, P)}{\partial T} \right]_{B, P}$$

$$S_{mag}^{4f}(T, B, P) = N_m \Re \left[\ln \sum_{m=-J}^{m=J} e^{-\beta E_m} + \frac{1}{k_B T} \frac{\sum_{m=-J}^{m=J} E_m e^{-\beta E_m}}{\sum_{m=-J}^{m=J} e^{-\beta E_m}} \right]$$

ALTERNATIVE CALCULATION

PARTITION FUNCTION AND FREE ENERGY

Partition function

$$Z_{mag}^{4f}(T, B, P) = \sum_{m=-J}^{m=J} e^{-\beta E_m}$$

$$E_m = -g\mu_B B^{eff} m$$

Geometric series

$$Z_{mag}^{4f}(T, B, P) = \sum_{m=-J}^{m=J} e^{-\beta E_m} = e^{\beta g\mu_B B^{eff} J} + e^{\beta g\mu_B B^{eff} (J-1)} + \dots + e^{-\beta g\mu_B B^{eff} J}$$

$$Z_{mag}^{4f}(T, B, P) = \frac{\sinh \left[J + \frac{1}{2} \right] y}{\sinh \left[\frac{y}{2} \right]}$$

$$y = \frac{g\mu_B B^{eff}}{k_B T}$$

Magnetic free energy

$$F_{mag}^{4f}(T, B, P) = -k_B T \ln \left\{ \frac{\sinh \left[J + \frac{1}{2} \right] y}{\sinh \left[\frac{y}{2} \right]} \right\}$$

MAGNETIZATION

Magnetic free energy

$$F_{mag}^{4f}(T, B, P) = -k_B T \ln \left\{ \frac{\sinh \left[J + \frac{1}{2} \right] y}{\sinh \left[\frac{y}{2} \right]} \right\}$$

Magnetization

$$M(T, B, P) = - \left[\frac{\partial F_{mag}^{4f}(T, B, P)}{\partial B^{eff}} \right]_{T, P}$$

$$M(T, B, P) = g\mu_B B_J(y)$$

$$B_J(y) = \frac{1}{J} \left\{ \left(\frac{2J+1}{2} \right) \coth \left[\left(\frac{2J+1}{2} \right) y \right] - \frac{1}{2} \coth \left(\frac{y}{2} \right) \right\}$$

$$y = \frac{g\mu_B B^{eff}}{k_B T}$$

Brillouin function

Self-consistency

$$B^{eff} = B + \frac{\mathcal{J}_0 \langle J^z \rangle}{g\mu_B}$$

MAGNETIC ENTROPY

Magnetic free energy

$$F_{mag}^{4f}(T, B, P) = -k_B T \ln \left\{ \frac{\sinh \left[\left(J + \frac{1}{2} \right) y \right]}{\sinh \left[\frac{y}{2} \right]} \right\} \quad y = \frac{g\mu_B B^{eff}}{k_B T}$$

Magnetic entropy

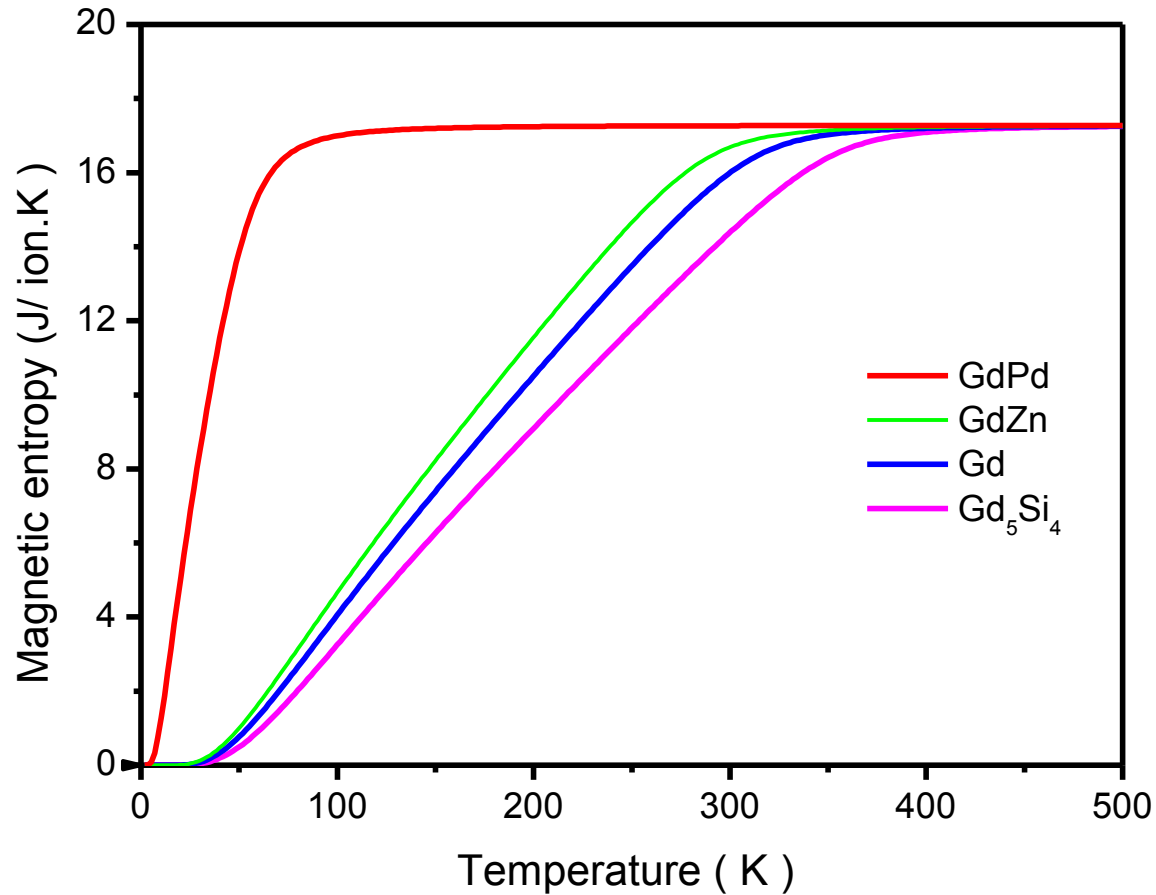
$$S_{mag}^{4f}(T, B, P) = - \left[\frac{\partial F_{mag}^{4f}(T, B, P)}{\partial T} \right]_{B, P}$$

$$S_{mag}^{4f}(T, B, P) = N_m \Re \left\{ \ln \left[\frac{\sinh \left(\left(J + \frac{1}{2} \right) y \right)}{\sinh \left(\frac{y}{2} \right)} \right] - \frac{g\mu_B B^{eff}}{k_B T} \left(J + \frac{1}{2} \right) \coth \left[\left(J + \frac{1}{2} \right) y \right] - \frac{1}{2} \coth \left[\left(\frac{y}{2} \right) \right] \right\}$$

Saturation value

$$S_{mag}^{4f}(T, B, P) |_{T \rightarrow \infty} = N_m \Re \ln (2J + 1)$$

MAGNETIC ENTROPY CURVES



$$S_{mag}^{4f}(T, B, P) |_{T \rightarrow \infty} = N_m \Re \ln(2J + 1)$$

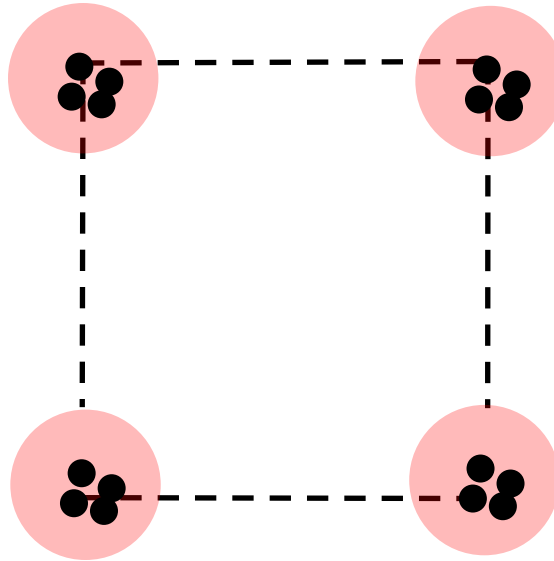
CRYSTAL LATTICE HAMILTONIAN

LATTICE HAMILTONIAN

Lattice Hamiltonian

$$\mathcal{H}_{lat} = \sum_q \hbar\omega_q a_q^\dagger a_q$$

$$\varepsilon_q = \sum_q \left(n_q + \frac{1}{2}\right) \hbar\omega_q$$



LATTICE : CANONICAL ENSEMBLE

Partition function

$$Z = \sum_{i=1}^n e^{-\beta E_i}$$

$$\varepsilon_q = \sum_q (n_q + \frac{1}{2}) \hbar \omega_q$$

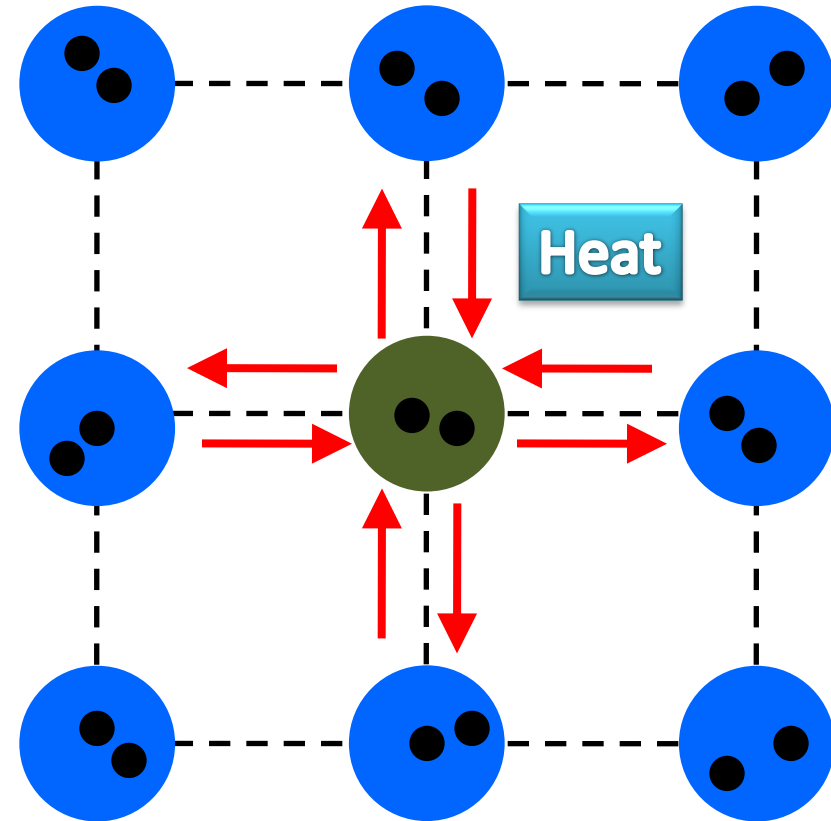
$$Z_{lat}(T, B, P) = \sum_{n_q = n_1, n_2, \dots} e^{-\beta \sum_q (n_q + \frac{1}{2}) \hbar \omega_q}$$

Bosons

$$Z_{lat}(T, B, P) = \prod_q \frac{1}{(1 - e^{-\beta \hbar \omega_q})}$$

Lattice free energy

$$F_{lat}(T, B, P) = 3N_A k_B T \sum_q \ln (1 - e^{-\beta \hbar \omega_q})$$



LATTICE ENTROPY

Lattice free energy

$$F_{lat}(T, B, P) = 3N_A k_B T \sum_q \ln (1 - e^{-\beta \hbar \omega_q})$$

Density of phonons

$$F_{lat}(T, B, P) = \Re T \int \ln (1 - e^{-\beta \hbar \omega}) \rho^{ph}(\omega) d\omega$$

Lattice entropy

$$S_{lat}(T, B, P) = - \left[\frac{\partial F_{lat}(T, B, P)}{\partial T} \right]_B$$

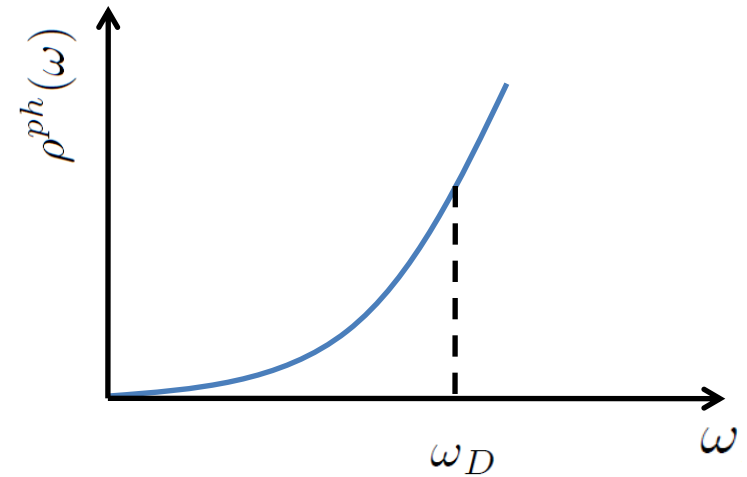
$$S_{lat}(T, B, P) = N_i \Re \left[- \int \ln (1 - e^{-\beta \hbar \omega}) \rho^{ph}(\omega) d\omega + \frac{1}{k_B T} \int \frac{\hbar \omega}{(e^{\beta \hbar \omega} - 1)} \rho^{ph}(\omega) d\omega \right]$$

LATTICE ENTROPY (Debye approximation)

$$S_{lat}(T, B, P) = N_i \Re \left[- \int \ln (1 - e^{-\beta \hbar \omega}) \rho^{ph}(\omega) d\omega + \frac{1}{k_B T} \int \frac{\hbar \omega}{(e^{\beta \hbar \omega} - 1)} \rho^{ph}(\omega) d\omega \right]$$

Debye approximation

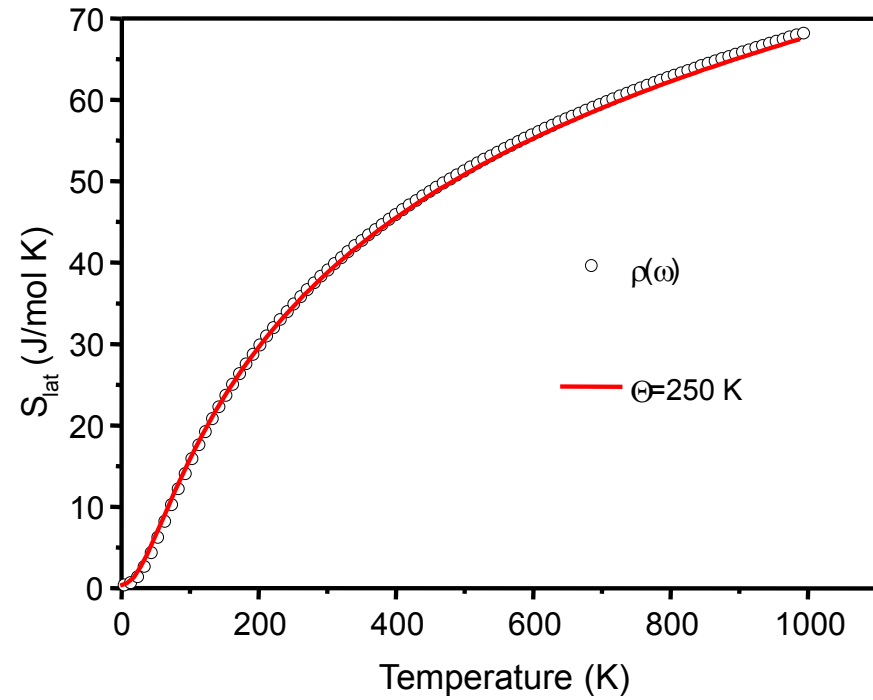
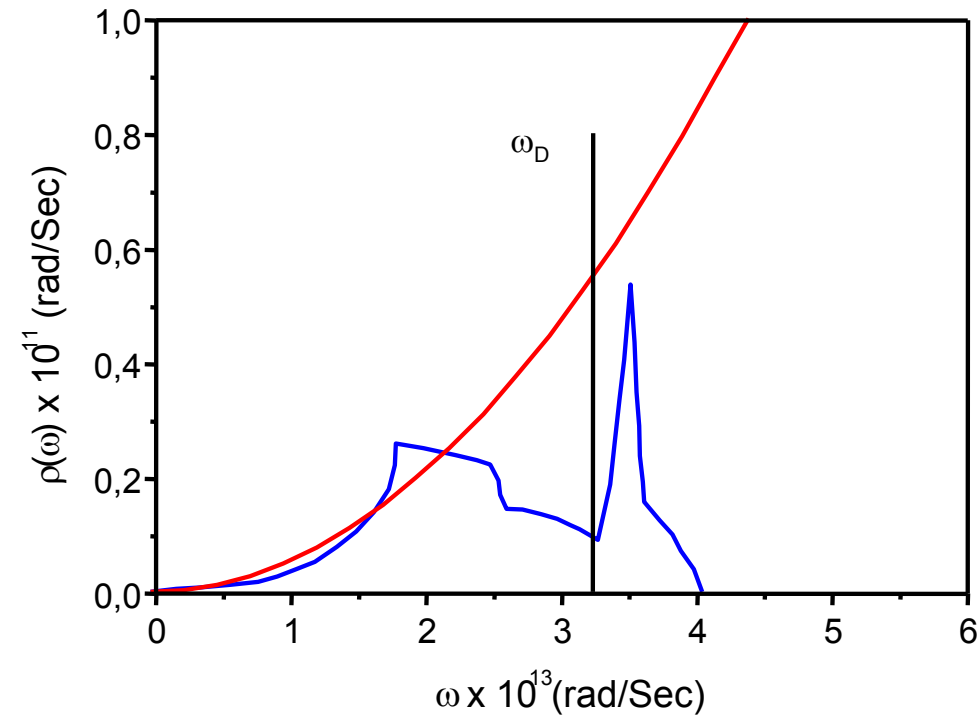
$$\begin{cases} \rho^{ph}(\omega) = \left[\frac{3V}{2\pi^2 v^3} \right] \omega^2 \\ \omega_D = \left[\frac{6\pi^2 v^3 N_A}{V} \right]^{1/3} \end{cases}$$



Lattice entropy

$$S_{lat}(T, B, P) = N_i \left[-3 \Re \ln \left(1 - e^{-\frac{\Theta_D}{T}} \right) + 12 \Re \left(\frac{T}{\Theta_D} \right)^3 \int_0^{\Theta_D/T} \frac{x^3}{e^x - 1} dx \right] \quad \Theta_D = \frac{\hbar \omega_D}{k_B}$$

LATTICE ENTROPY



$$S_{lat}(T, B, P) = N_i \Re \left[- \int \ln (1 - e^{-\beta \hbar \omega}) \rho^{ph}(\omega) d\omega + \frac{1}{k_B T} \int \frac{\hbar \omega}{(e^{\beta \hbar \omega} - 1)} \rho^{ph}(\omega) d\omega \right]$$

$$S_{lat}(T, B, P) = N_i \left[-3 \Re \ln \left(1 - e^{-\frac{\Theta_D}{T}} \right) + 12 \Re \left(\frac{T}{\Theta_D} \right)^3 \int_0^{\Theta_D/T} \frac{x^3}{e^x - 1} dx \right]$$

NON MAGNETIC CONDUCTION ELECTRONS

Hamiltonian

$$\mathcal{H}_{el}^{spd} = \sum_k \varepsilon_k c_k^+ c_k$$

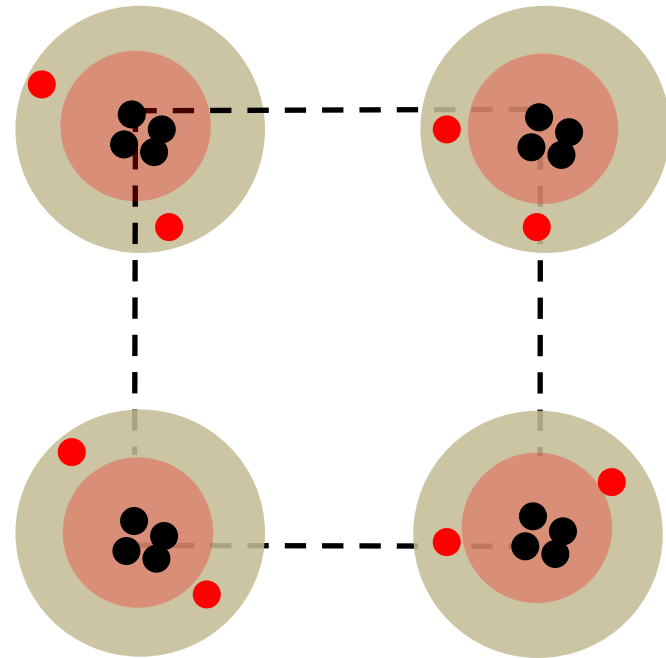
Partition function

$$Z_{el}^{spd}(T, B, P) = \prod_{k\sigma} \left\{ 1 + \exp \left[-\beta \left(\varepsilon_k^{spd} - \mu \right) \right] \right\}$$

Free energy

$$F_{el}^{spd}(T, B, P) = -\frac{1}{\beta} \sum_{l=1}^5 \sum_{k\sigma} \ln \left\{ 1 + \exp \left[-\beta \left(\varepsilon_k^{spd} - \mu \right) \right] \right\}$$

Gran-canonical ensemble



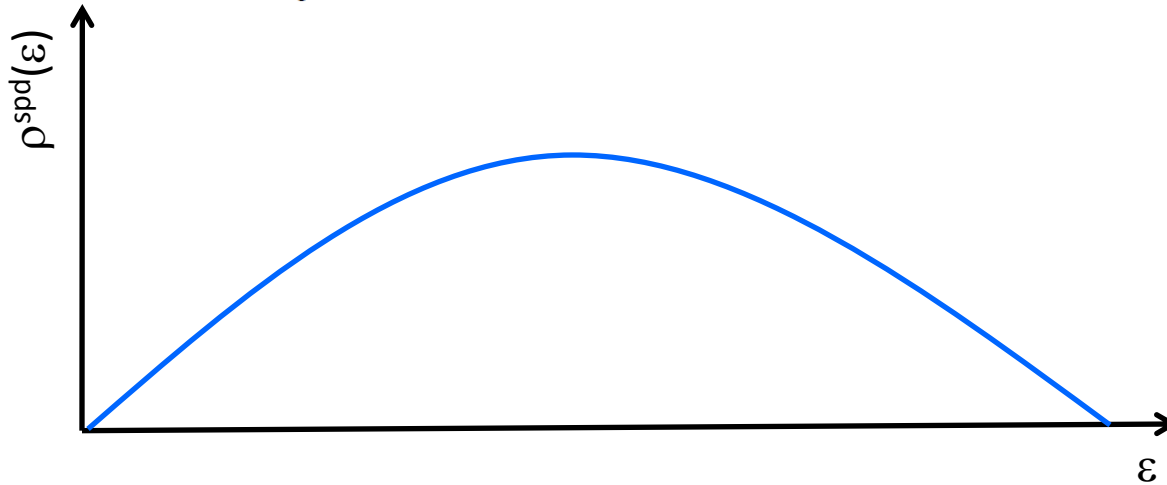
FREE ENERGY

Free energy

$$F_{el}^{spd}(T, B, P) = -\frac{1}{\beta} \sum_{l=1}^5 \sum_{k\sigma} \ln \left\{ 1 + \exp \left[-\beta \left(\varepsilon_k^{spd} - \mu \right) \right] \right\}$$

Density of phonons

$$F_{el}^{spd}(T, B, P) = -\mathcal{R}T \sum_{\sigma} \int_{-\infty}^{\infty} \ln \left\{ 1 + \exp \left[-\beta \left(\varepsilon_{\sigma} - \mu \right) \right] \right\} \rho_{\sigma}^{spd}(\varepsilon) d\varepsilon$$



ELECTRONIC ENTROPY

Free energy

$$F_{el}^{spd}(T, B, P) = -\mathfrak{R}T \sum_{\sigma} \int_{-\infty}^{\infty} \ln \{1 + \exp[-\beta(\varepsilon_{\sigma} - \mu)]\} \rho_{\sigma}^{spd}(\varepsilon) d\varepsilon$$

Electronic entropy

$$S_{el}^{spd} = - \left[\frac{\partial F_{el}^{spd}(T, B, P)}{\partial T} \right]_B$$

$$S_{el}^{spd}(T, B, P) = N_{el} \mathfrak{R} \left\{ \sum_{\sigma} \int_{-\infty}^{\infty} \ln \{1 + \exp[-\beta(\varepsilon - \mu)]\} \rho_{\sigma}^{spd}(\varepsilon) d\varepsilon + \frac{1}{k_B T} \sum_{\sigma} \int_{-\infty}^{\infty} (\varepsilon - \mu) f(\varepsilon) \rho_{\sigma}^{spd}(\varepsilon) d\varepsilon \right\}$$

Sommerfeld approximation

$$S_{el}^{spd}(T) = \gamma T$$

$$\gamma = \frac{\pi^2 k_B^2 \rho_{\sigma}^{spd}(\varepsilon_f)}{3}$$

MCE QUANTITIES

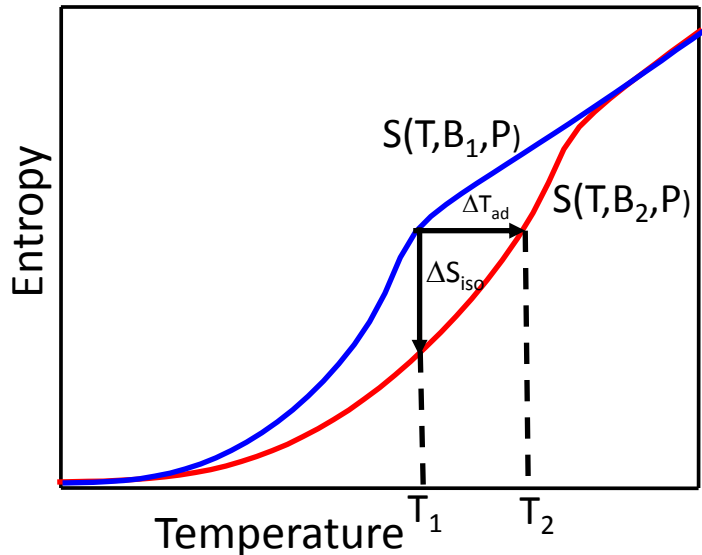
Total entropy

$$S(T, B, P) = S_{mag}^{4f}(T, B, P) + S_{lat}(T, B, P) + S_{el}^{spd}(T, B, P)$$

$$S(T, B, P) = N_m \mathfrak{R} \left[\ln \sum_{m=-J}^{m=J} e^{-\beta E_m} + \frac{1}{k_B T} \frac{\sum_{m=-J}^{m=J} E_m e^{-\beta E_m}}{\sum_{m=-J}^{m=J} e^{-\beta E_m}} \right]$$

$$+ N_i \left[-3\mathfrak{R} \ln \left(1 - e^{-\frac{\Theta_D}{T}} \right) + 12\mathfrak{R} \left(\frac{T}{\Theta_D} \right)^3 \int_0^{\Theta_D/T} \frac{x^3}{e^x - 1} dx \right] + \gamma T$$

Magnetocaloric quantities

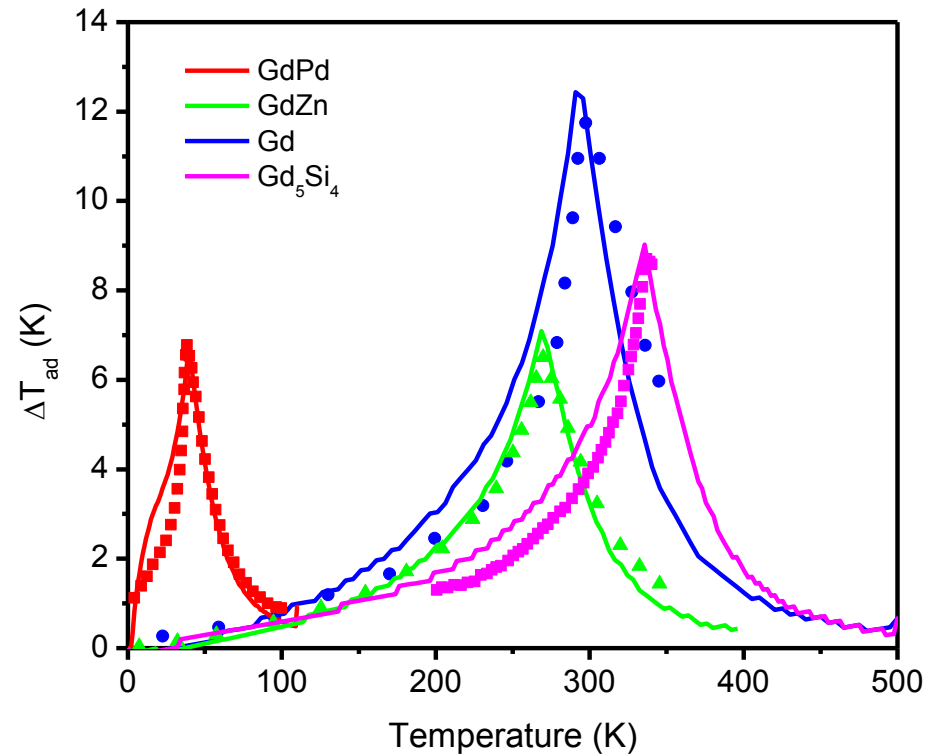
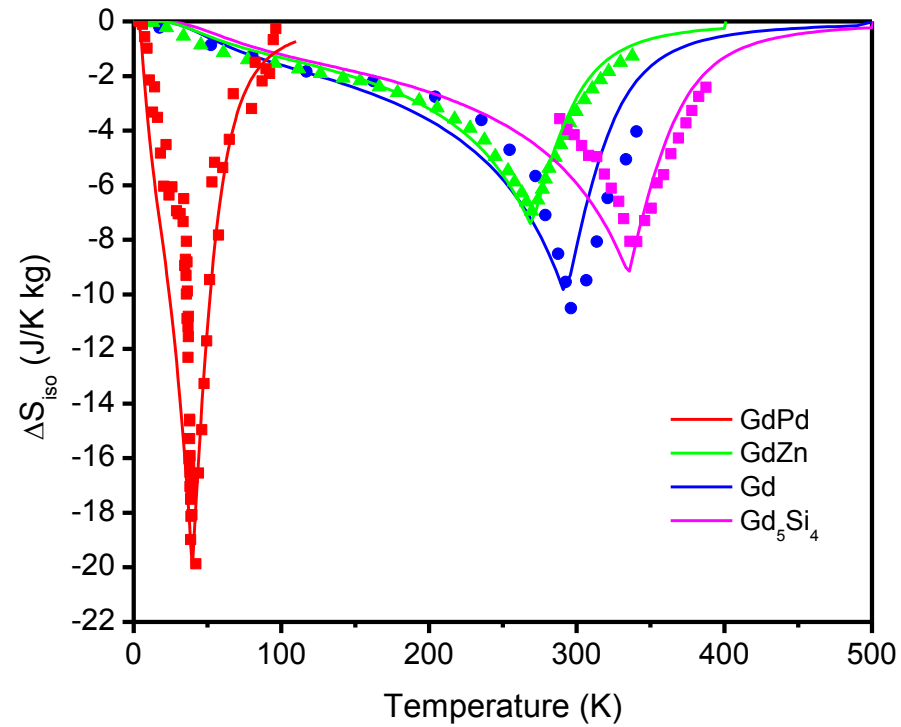


$$\Delta S_{iso}(T, \Delta B, P) = S(T, B_2, P) - S(T, B_1, P)$$

$$\Delta T_{ad}(T, \Delta B, P) = T_2 - T_1$$

$$S(T, B_2, P) = S(T, B_1, P)$$

APPLICATION: Gd COMPOUNDS



ANISOTROPY

HAMILTONIAN

Total

$$\mathcal{H} = \mathcal{H}_{lat} + \mathcal{H}_{el}^{spd} + \mathcal{H}_{mag}^{4f}$$

Lattice

$$\mathcal{H}_{lat} = \sum_a \hbar \omega_q a_q^+ a_q$$

Non magnetic electrons

$$\mathcal{H}_{el}^{spd} = \sum_k \varepsilon_k c_k^+ c_k$$

Magnetic electrons

$$\mathcal{H}_{mag}^{4f} = - \sum_{i,j} \mathcal{J}_{ij}(r) \vec{J}_i \cdot \vec{J}_j - \sum_i g \mu_B \vec{B} \cdot \vec{J}_i - \sum_i D J_{iz}^2$$

MEAN FIELD APPROXIMATION

Hamiltonian

$$\mathcal{H}_{mag}^{4f} = - \sum_{i,j} \mathcal{J}_{ij}(r) \vec{J}_i \cdot \vec{J}_j - \sum_i g\mu_B \vec{B} \cdot \vec{J}_i - \sum_i DJ_{iz}^2$$

Mean field approximation

$$\mathcal{H}_{mag}^{4f} = -\mathcal{J}_0 \sum_i \vec{J}_i \cdot \langle \vec{J} \rangle - \sum_i g\mu_B \vec{B} \cdot \vec{J}_i - \sum_i DJ_{iz}^2$$

Magnetic electrons

$$\mathcal{H}_{mag}^{4f} = - \sum_i g\mu_B B_x^{eff} J_{ix} + g\mu_B B_y^{eff} J_{iy} + g\mu_B B_z^{eff} J_{iz}$$

$$B_x^{eff} = B \cos \theta_x + \frac{\mathcal{J}_0 \langle J_x \rangle}{g\mu_B}$$

$$B_y^{eff} = B \cos \theta_y + \frac{\mathcal{J}_0 \langle J_y \rangle}{g\mu_B}$$

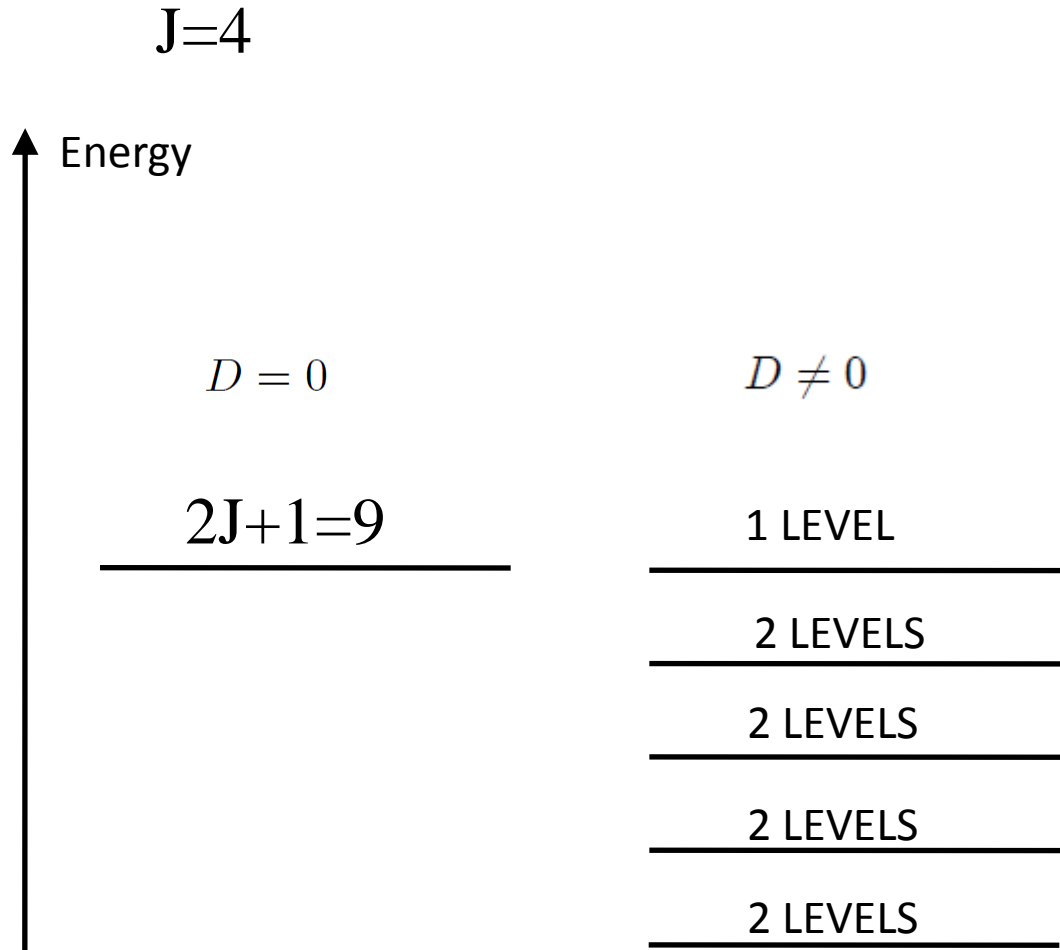
$$B_z^{eff} = B \cos \theta_z + \frac{\mathcal{J}_0 \langle J_z \rangle + DJ_z}{g\mu_B}$$

ENERGY LEVELS

$$B_x^{eff} = B \cos \theta_x + \frac{\mathcal{J}_0 \langle J_x \rangle}{g\mu_B}$$

$$B_y^{eff} = B \cos \theta_y + \frac{\mathcal{J}_0 \langle J_y \rangle}{g\mu_B}$$

$$B_z^{eff} = B \cos \theta_z + \frac{\mathcal{J}_0 \langle J_z \rangle + DJ_z}{g\mu_B}$$



MATRIX HAMILTONIAN

Total

$$\mathcal{H}_{mag}^{4f} = - \sum_i g\mu_B B_x^{eff} J_{ix} + g\mu_B B_y^{eff} J_{iy} + g\mu_B B_z^{eff} J_{iz}$$

Mean field approximation

$$\mathcal{H}_{mag}^{4f} = \begin{bmatrix} \mathcal{H}_{11} & \mathcal{H}_{12} & \cdots & \mathcal{H}_{1(2J+1)} \\ \mathcal{H}_{21} & \mathcal{H}_{22} & \cdots & \mathcal{H}_{2(2J+1)} \\ \vdots & \vdots & \ddots & \vdots \\ \mathcal{H}_{(2J+1)1} & \mathcal{H}_{(2J+1)2} & \cdots & \mathcal{H}_{(2J+1)(2J+1)} \end{bmatrix}$$

$$\mathcal{H}_{ij} = \langle \psi_i | \mathcal{H}_{mag}^{4f} | \psi_j \rangle$$

Energy eigenvalues and eigenvectors

$$E_m \quad | \psi_m \rangle$$

MAGNETIZATION

Magnetization

$$M(T, B, P) = \hat{i}M_x(T, B, P) + \hat{j}M_y(T, B, P) + \hat{k}M_z(T, B, P)$$

Magnetization components

$$M_x(T, B, P) = g\mu_B \langle J_x \rangle$$

$$M_y(T, B, P) = g\mu_B \langle J_y \rangle$$

$$M_z(T, B, P) = g\mu_B \langle J_z \rangle$$

Average values

$$\langle J_x \rangle = \frac{\sum \langle \psi_m | J_x | \psi_m \rangle}{\sum_m e^{-\beta E_m}}$$

$$\langle J_y \rangle = \frac{\sum \langle \psi_m | J_y | \psi_m \rangle}{\sum_m e^{-\beta E_m}}$$

$$\langle J_z \rangle = \frac{\sum \langle \psi_m | J_z | \psi_m \rangle}{\sum_m e^{-\beta E_m}}$$

PARTITION FUNCTION AND FREE ENERGY

Partition function

$$Z_{mag}(T, B, P) = \sum_{m=-J}^{m=J} e^{-\beta E_m}$$

Magnetic free energy

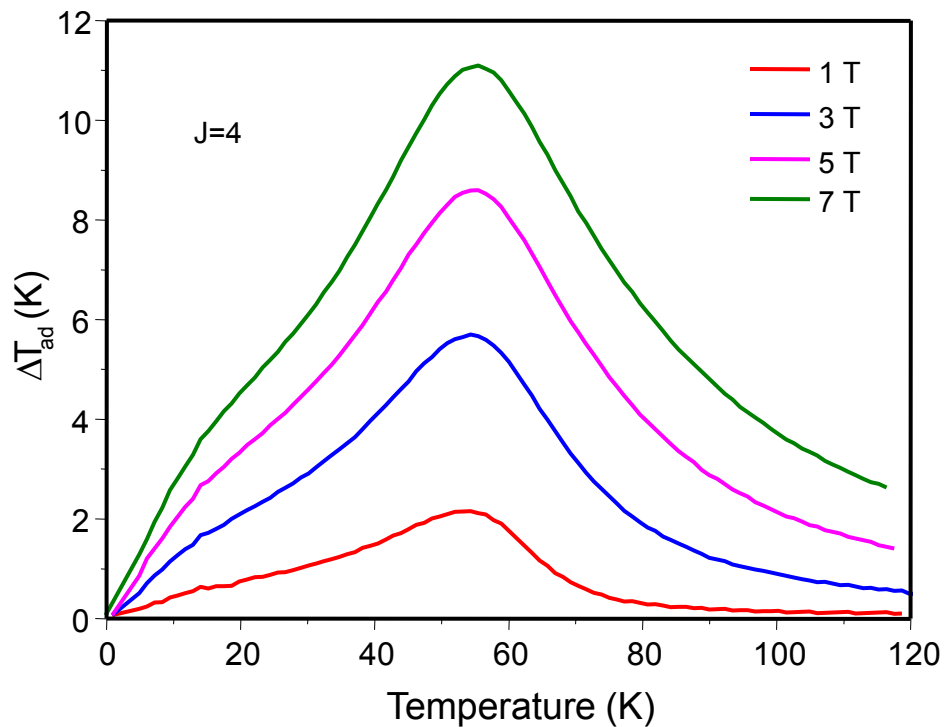
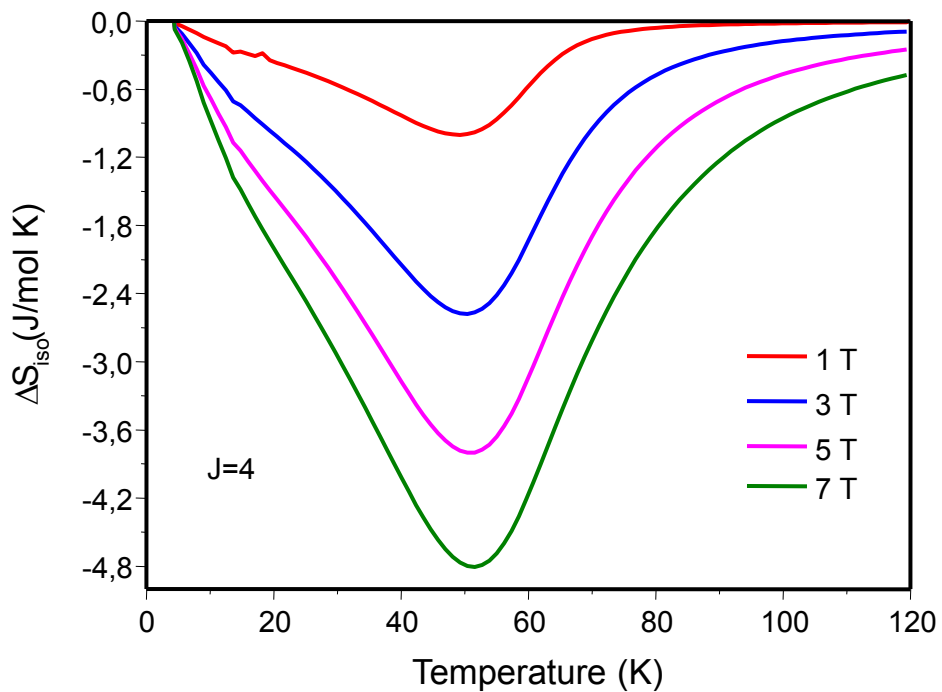
$$F_{mag}^{4f} = -k_B T \ln \sum_m e^{-\beta E_m}$$

Magnetic entropy

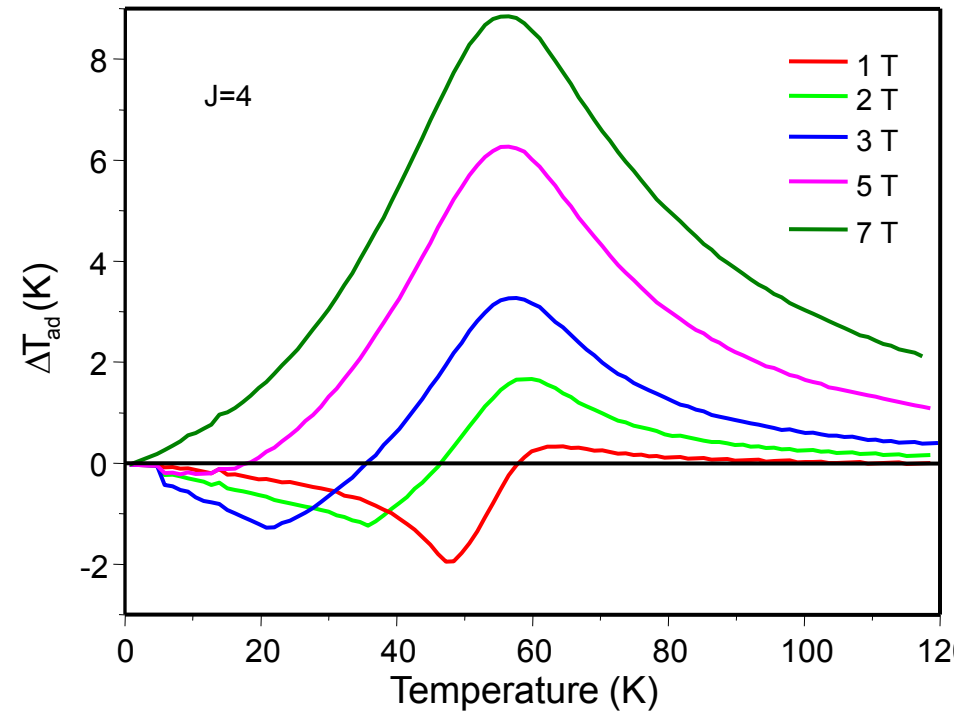
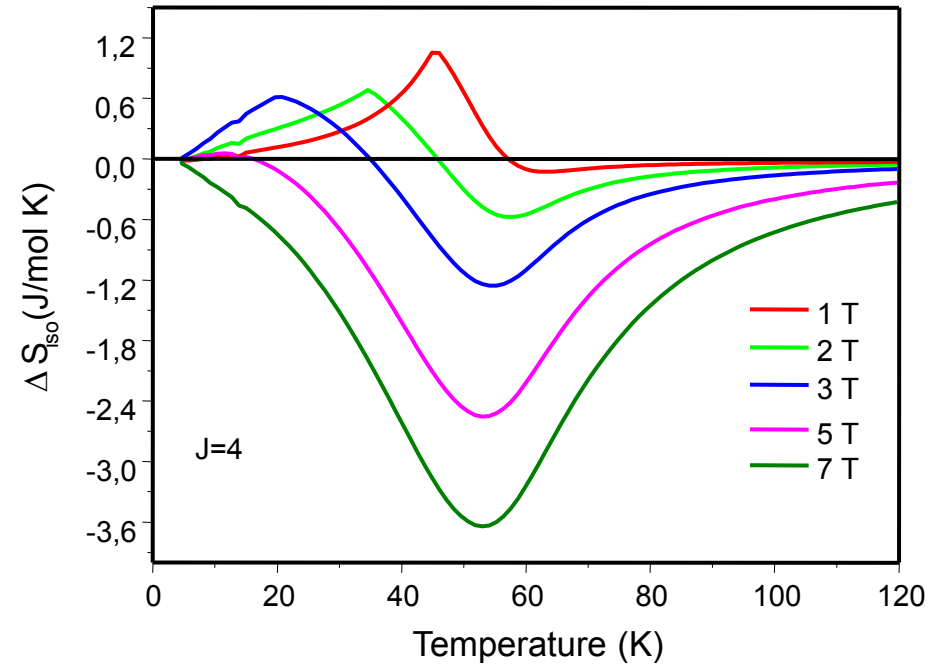
$$S_{mag}(T, B, P) = - \left[\frac{\partial F_{mag}(T, B, P)}{\partial T} \right]_{B, P}$$

$$S_{mag}(T, B, P) = N_m \Re \left[\ln \sum_{m=-J}^{m=J} e^{-\beta E_m} + \frac{1}{k_B T} \frac{\sum_{m=-J}^{m=J} E_m e^{-\beta E_m}}{\sum_{m=-J}^{m=J} e^{-\beta E_m}} \right]$$

ANISOTROPIC SYSTEM ($B=B_z$)

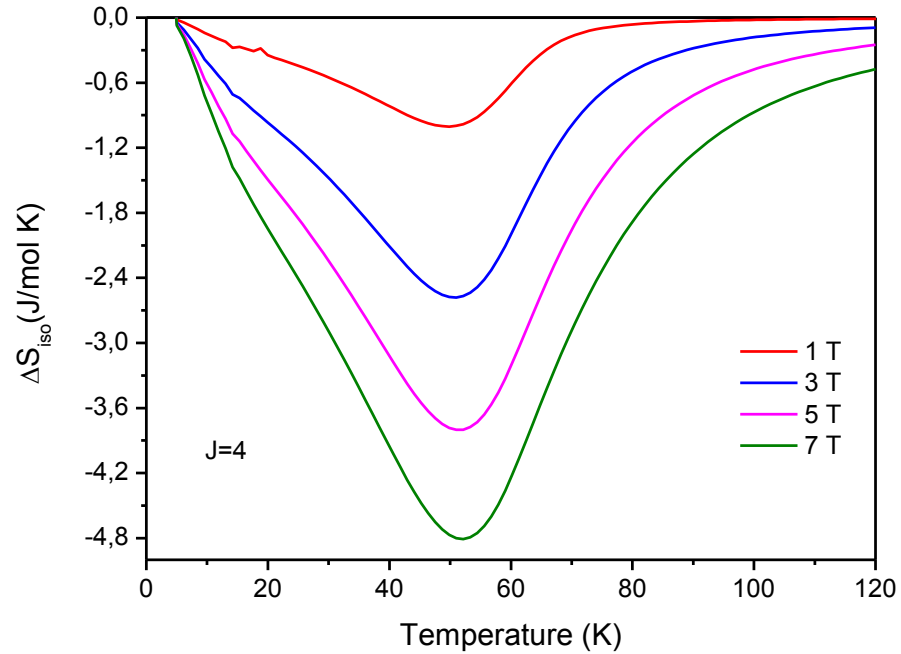


ANISOTROPIC SYSTEM ($B=B_x$)

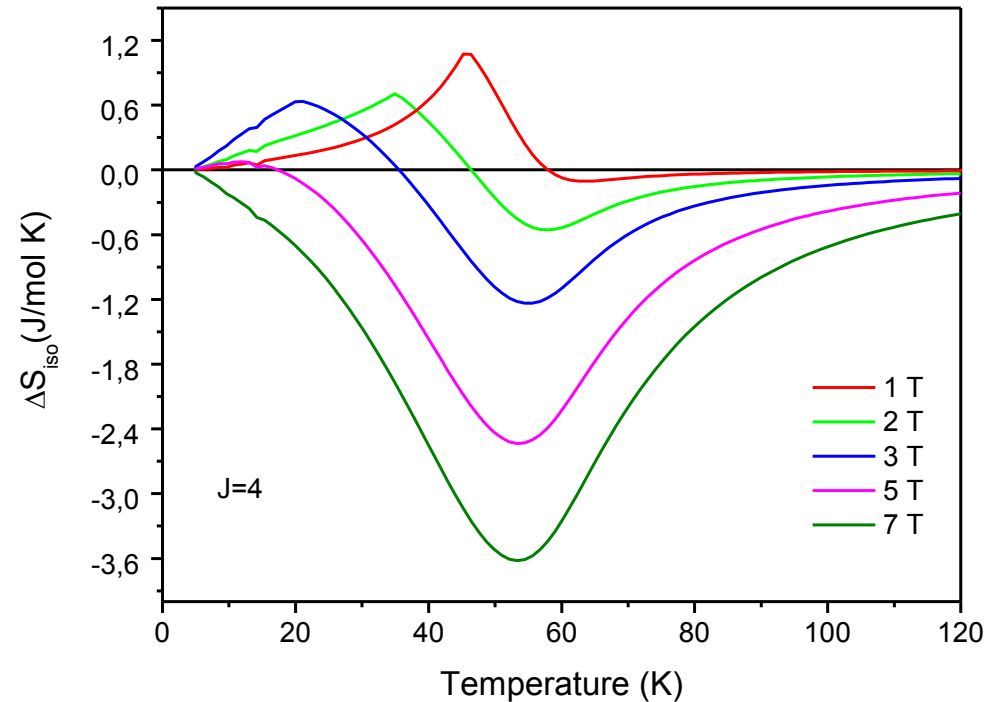


COMPARISON: ΔS_{iso}

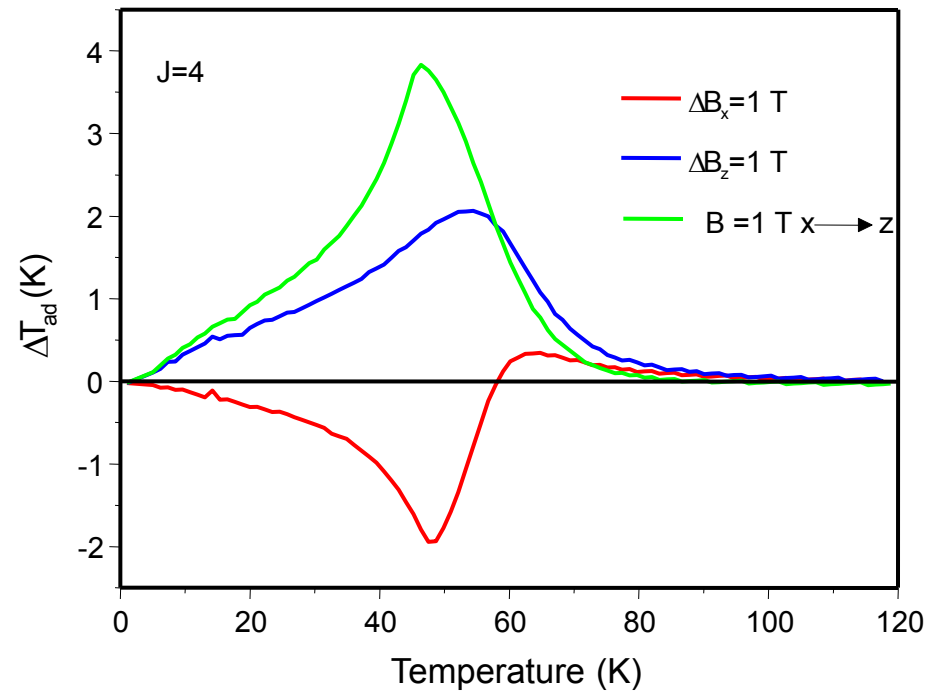
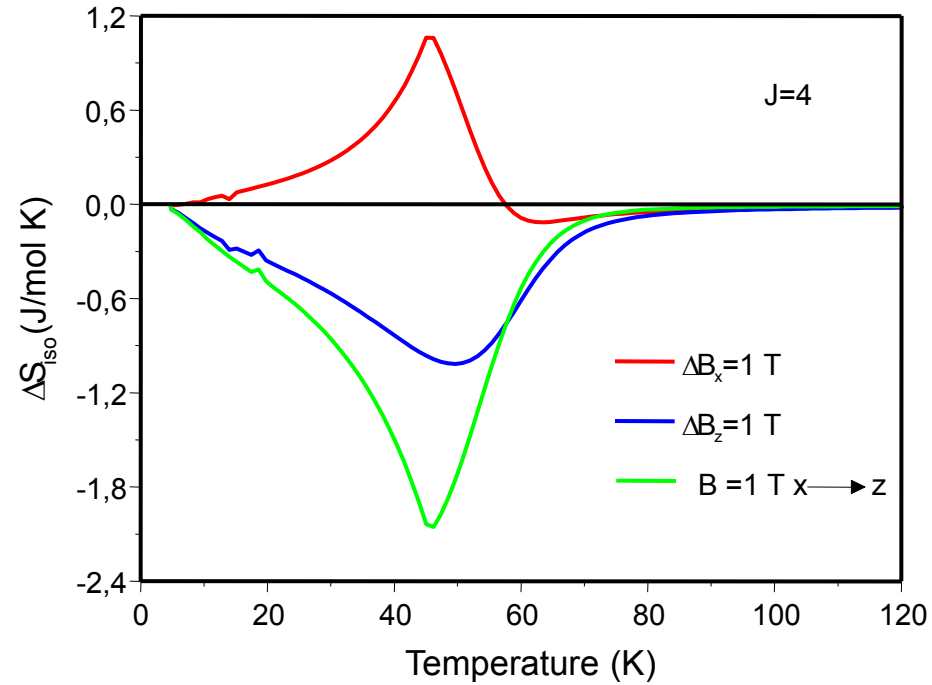
B=B_z



B=B_x

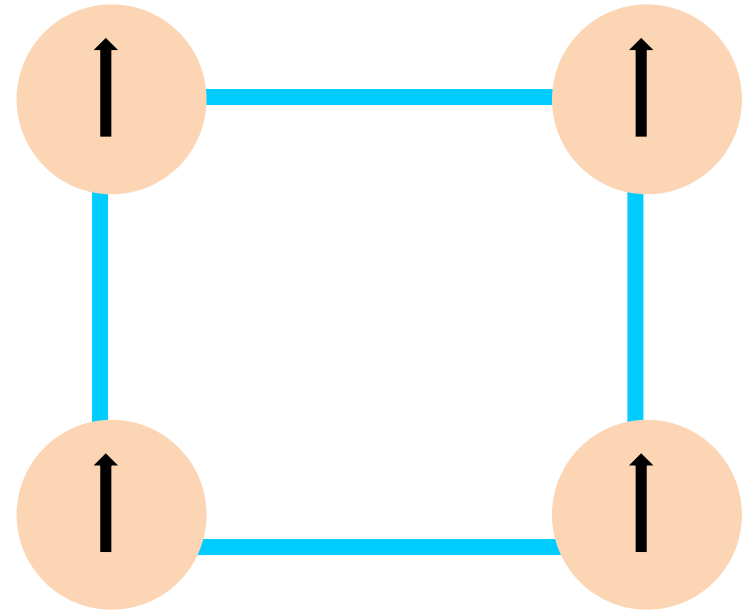
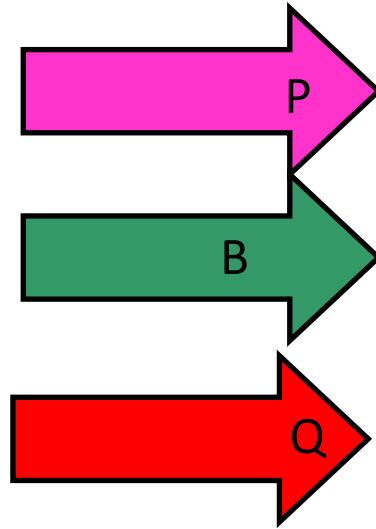


ANISOTROPIC MCE $B_x \rightarrow B_z$



FIRST ORDER MAGNETIC PHASE TRANSITION

MAGNETOELASTIC COUPLING



$$\mathcal{H}_{mag}^{4f} = - \sum_{ij} \tilde{\mathcal{J}}(r) J_i \cdot J_j - \sum_i g \mu_B B J_i$$

$$\mathcal{J}_{ij}(r) = \mathcal{J}_0(r_0) + \mathcal{J}_1(r_0) J_i J_j$$

$$\mathcal{J}_1 = \left[\frac{d\mathcal{J}(r)}{dr} \right]_{r=r_0}^2$$

$$\mathcal{H}_{mag}^{4f} = - \sum_{ij} \tilde{\mathcal{J}}_0 J_i \cdot J_j - \sum_{ij} \tilde{\mathcal{J}}_1 (J_i J_j)^2 - \sum_i g \mu_B B J_i$$

MEAN FIELD THEORY

$$\mathcal{H}_{mag}^{4f} = - \sum_{i,j} \mathcal{J}_0 J_i J_j - \sum_{i,j} \mathcal{J}_1 (J_i J_j)^2 - \sum_i g \mu_B B J_i$$

$$\mathcal{H}_{mag}^{4f} = - \sum_i \left[\mathcal{J}_0 \langle J \rangle + \mathcal{J}_1 \langle J \rangle^3 + g \mu_B \vec{B} \right] \cdot J_i$$

$$\mathcal{H}_{mag}^{4f} = -g \mu_B \sum_i B^{eff} J_i$$

$$B^{eff} = B + \frac{\mathcal{J}_0 \langle J \rangle + \mathcal{J}_1 \langle J \rangle^3}{g \mu_B}$$

LATTICE ENTROPY (REVISITED)

Debye approximation

$$S_{lat}(T, B, P) = N_i \left[-3\Re \ln \left(1 - e^{-\frac{\Theta_D}{T}} \right) + 12\Re \left(\frac{T}{\Theta_D} \right)^3 \int_0^{\Theta_D/T} \frac{x^3}{e^x - 1} dx \right] \quad \Theta_D = \frac{\hbar\omega_D}{k_B}$$

Renormalized Debye frequency

$$\omega_D = \left[\frac{6\pi^2 v^3 N_A}{V} \right]^{1/3} \quad \tilde{\omega}_D = \left[\frac{6\pi^2 v^3 N_A}{V_0 + \Delta V} \right]^{1/3} \quad \tilde{\omega}_D = \omega_D \left[1 - \frac{1}{3} \frac{\Delta V}{V_0} \right]$$
$$\tilde{\Theta}_D = \frac{\hbar\tilde{\omega}_D}{k_B} \quad \tilde{\Theta}_D = \frac{\hbar\omega_D}{k_B} \left[1 - \frac{1}{3} \frac{\Delta V}{V_0} \right] \quad \tilde{\Theta}_D = \Theta_D [1 - \alpha M^2]$$

Renormalized lattice entropy

$$S_{lat}(T, B, P) = N_i \left[-3\Re \ln \left(1 - e^{-\frac{\tilde{\Theta}_D}{T}} \right) + 12\Re \left(\frac{T}{\tilde{\Theta}_D} \right)^3 \int_0^{\tilde{\Theta}_D/T} \frac{x^3}{e^x - 1} dx \right]$$

MCE QUANTITIES

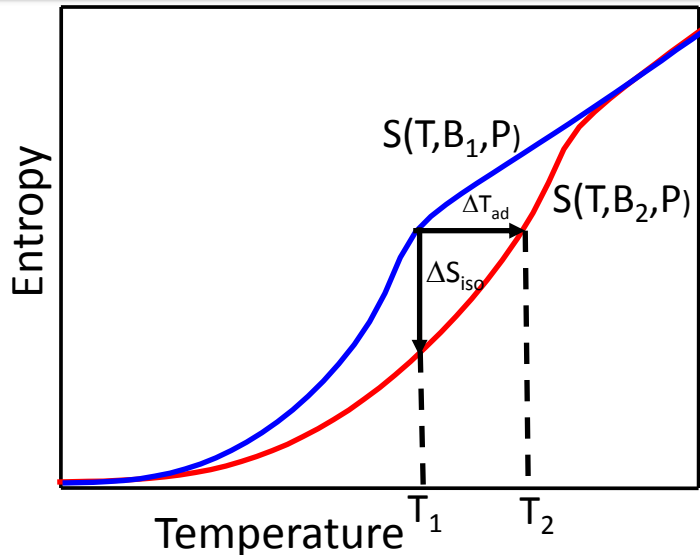
Total entropy

$$S(T, B, P) = S_{mag}^{4f}(T, B, P) + S_{lat}(T, B, P) + S_{el}^{spd}(T, B, P)$$

$$S(T, B, P) = N_m \mathfrak{R} \left[\ln \sum_m e^{-\beta E_m} + \frac{1}{k_B T} \frac{\sum_m E_m e^{-\beta E_m}}{\sum_m e^{-\beta E_m}} \right]$$

$$+ N_i \left[-3 \mathfrak{R} \ln \left(1 - e^{-\frac{\tilde{\Theta}_D}{T}} \right) + 12 \mathfrak{R} \left(\frac{T}{\tilde{\Theta}_D} \right)^3 \int_0^{\frac{\tilde{\Theta}_D}{T}} \frac{x^3}{e^x - 1} dx \right] + \gamma T$$

Magnetocaloric quantities

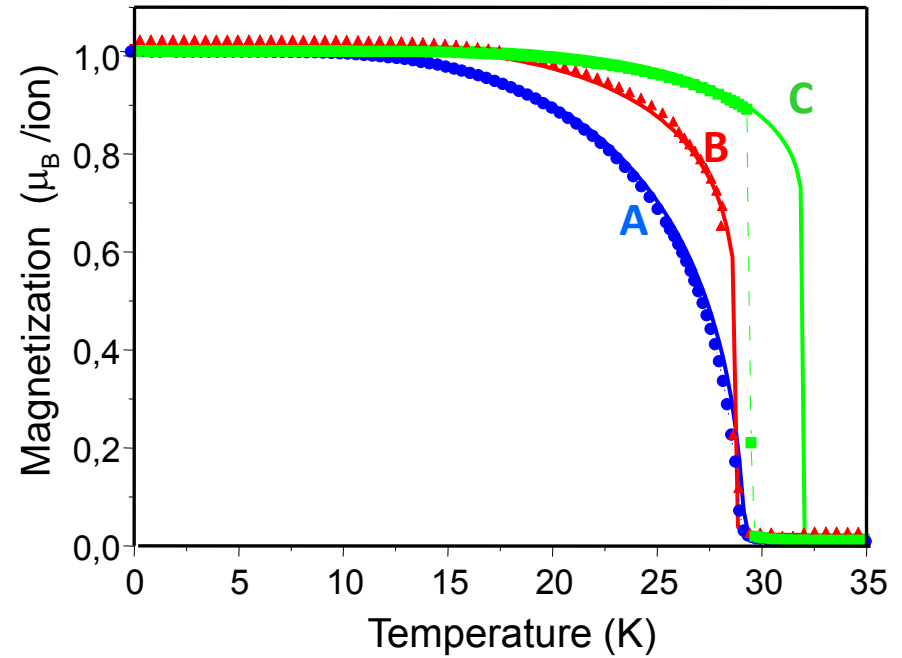
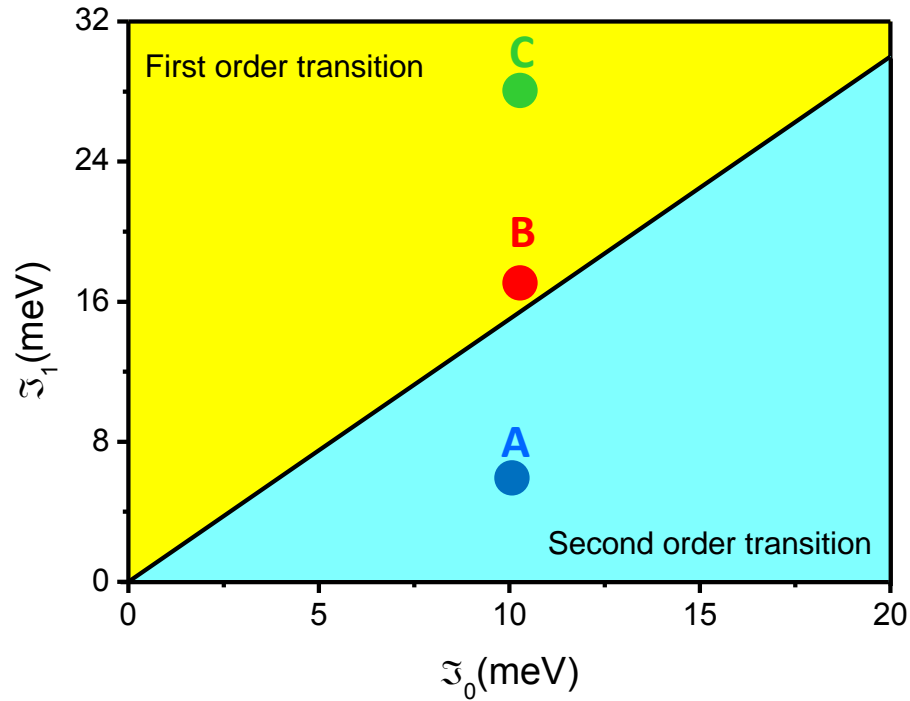


$$\Delta S_{iso}(T, \Delta B, P) = S(T, B_2, P) - S(T, B_1, P)$$

$$\Delta T_{ad}(T, \Delta B, P) = T_2 - T_1$$

$$S(T, B_2, P) = S(T, B_1, P)$$

SYSTEMATIC STUDY (J=1/2)



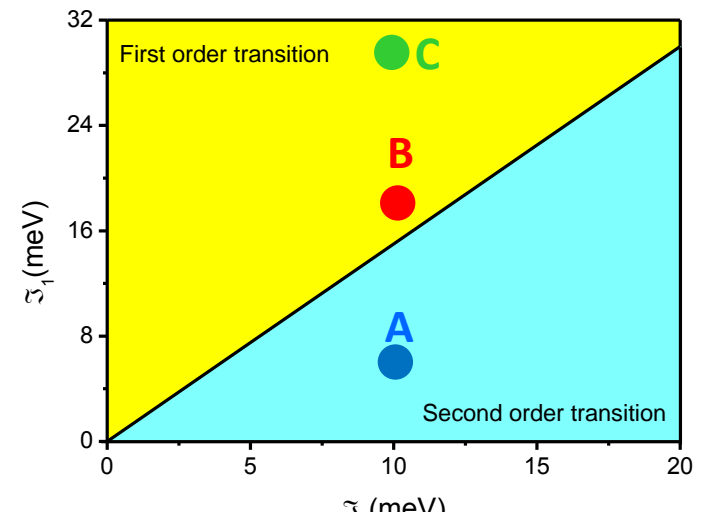
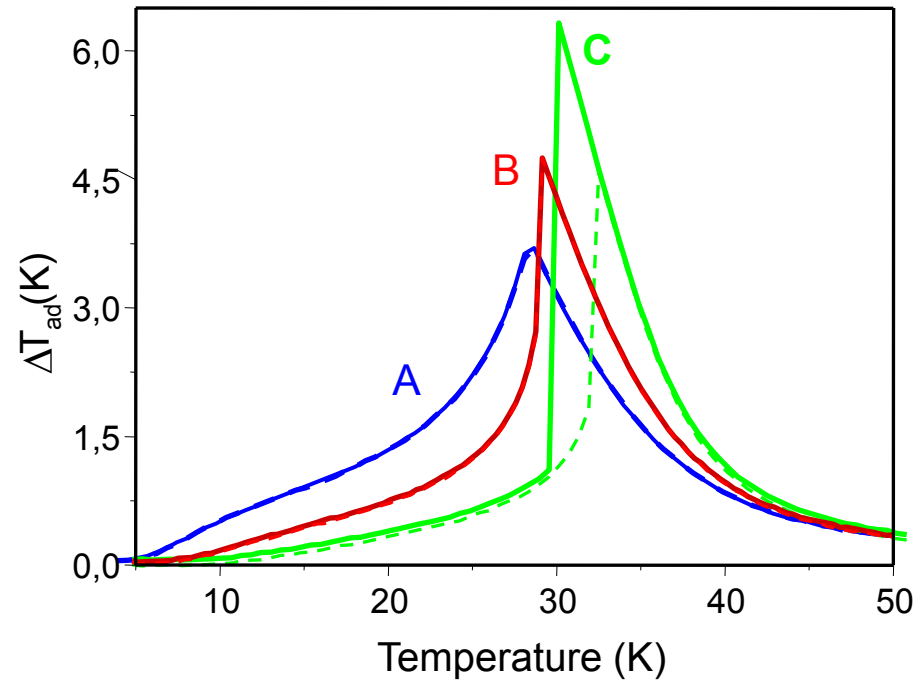
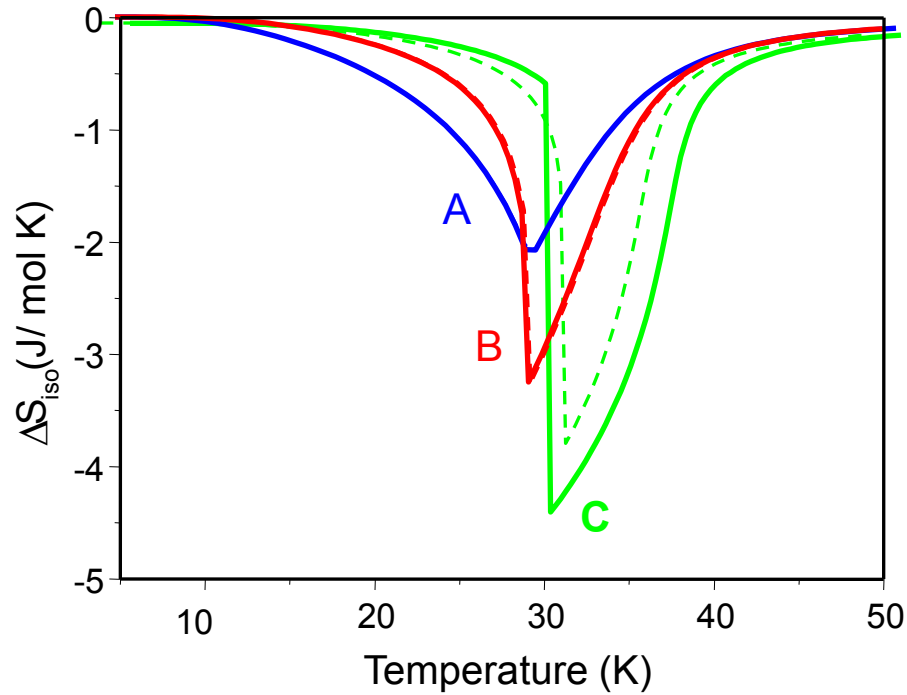
$$\mathcal{H}_{mag}^{4f} = - \sum_{ij} \tilde{\mathfrak{J}}_0 J_i \cdot J_j - \sum_{ij} \tilde{\mathfrak{J}}_1 (J_i J_j)^2 - \sum_i g \mu_B B J_i$$

$$\mathcal{J}_1 / \mathcal{J}_0$$

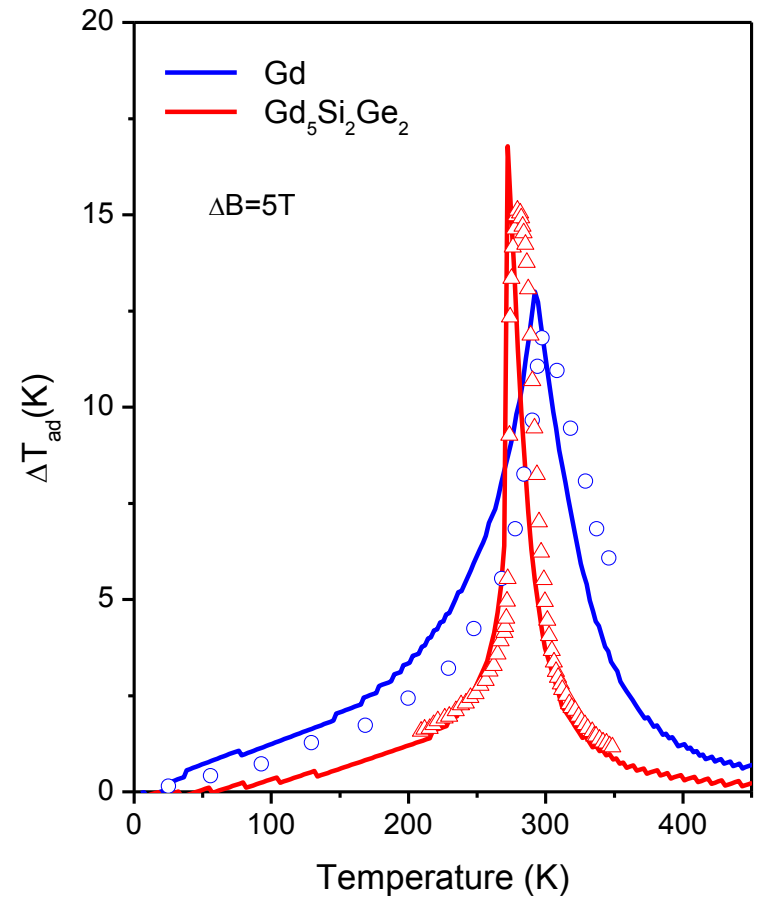
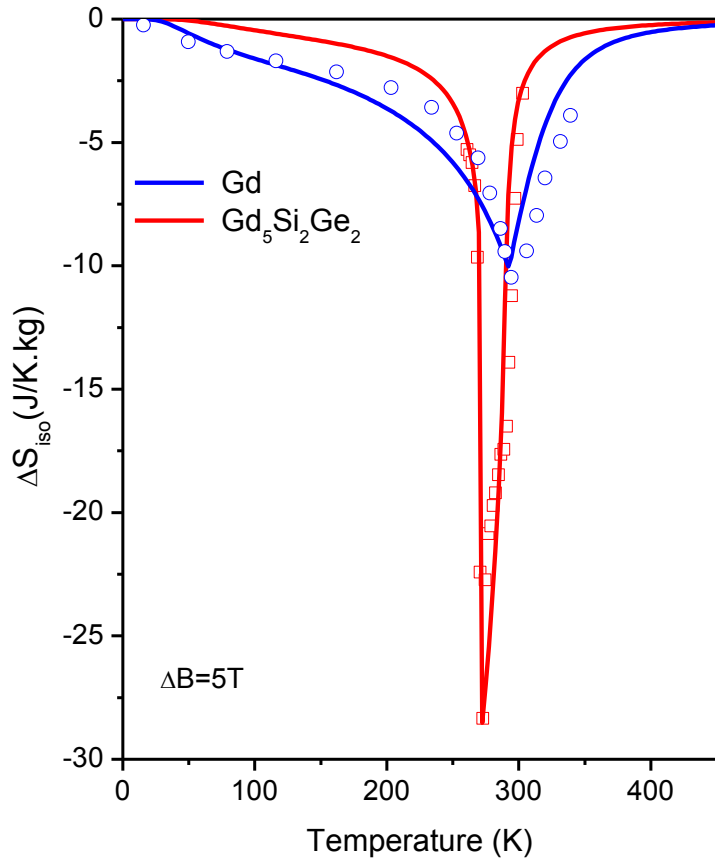
2nd order

1st order

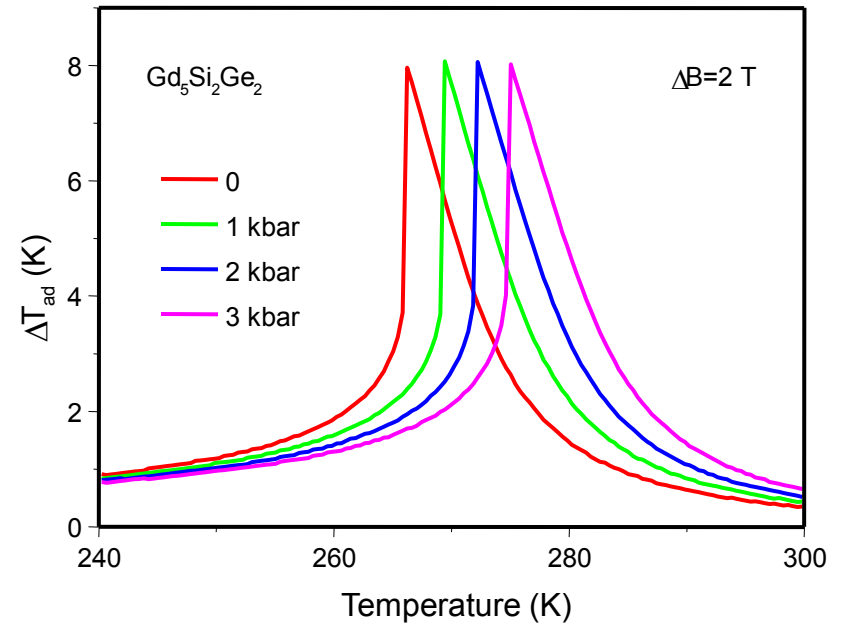
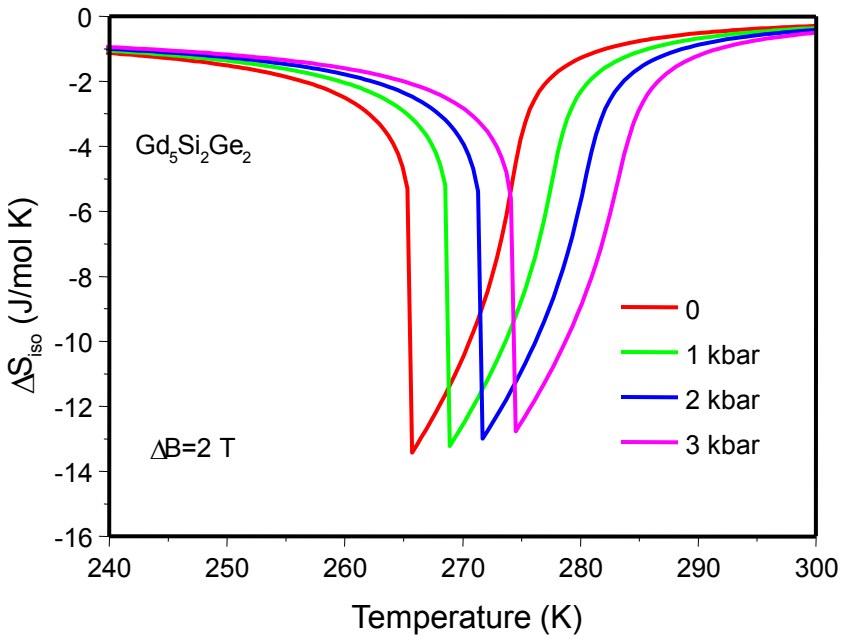
SYSTEMATIC STUDY (J=1/2)



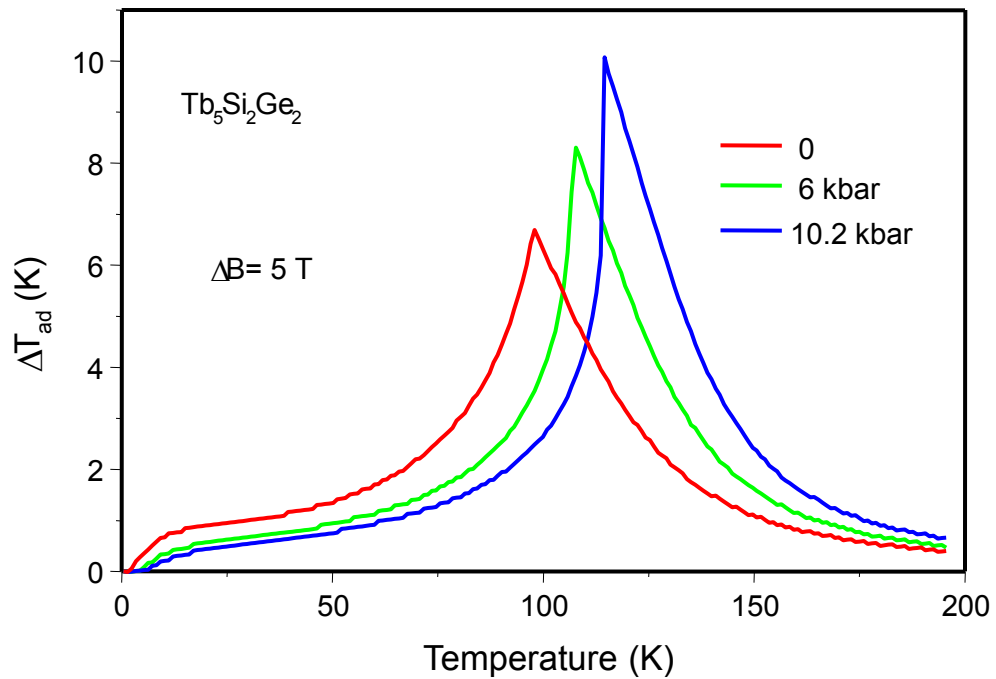
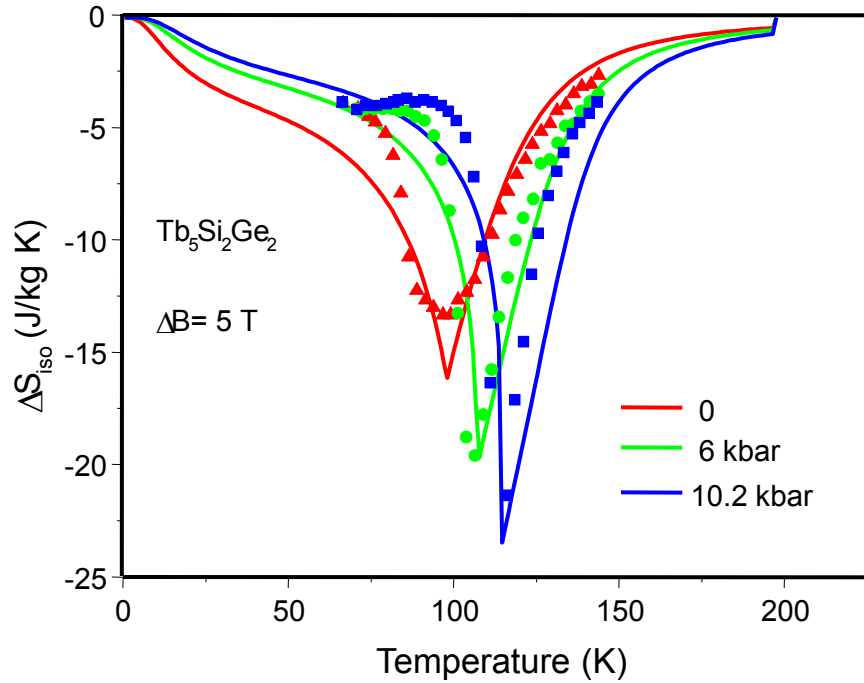
APPLICATION: $\text{Gd}_5\text{Si}_2\text{Ge}_2$



PRESSURE EFFECTS: $\text{Gd}_5\text{Si}_2\text{Ge}_2$



PRESSURE EFFECTS: $\text{Tb}_5\text{Si}_2\text{Ge}_2$



Morelon et al, Phys. Rev. Let. 93 (2004) 137201

N. A. de Oliveira, Journ. Appl. Phys. 113 (2013) 033910

BCE QUANTITIES

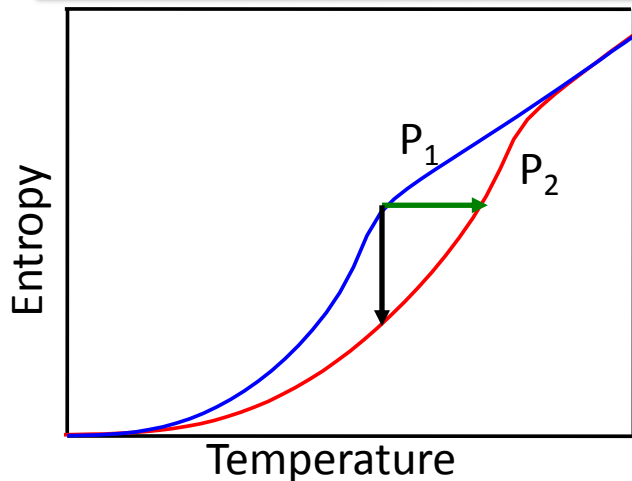
Total entropy

$$S(T, B, P) = S_{mag}^{4f}(T, B, P) + S_{lat}(T, B, P) + S_{el}^{spd}(T, B, P)$$

$$S(T, B, P) = N_m \mathfrak{R} \left[\ln \sum_m e^{-\beta E_m} + \frac{1}{k_B T} \frac{\sum_m E_m e^{-\beta E_m}}{\sum_m e^{-\beta E_m}} \right]$$

$$+ N_i \left[-3 \mathfrak{R} \ln \left(1 - e^{-\frac{\tilde{\Theta}_D}{T}} \right) + 12 \mathfrak{R} \left(\frac{T}{\tilde{\Theta}_D} \right)^3 \int_0^{\frac{\tilde{\Theta}_D}{T}} \frac{x^3}{e^x - 1} dx \right] + \gamma T$$

Barocaloric quantities

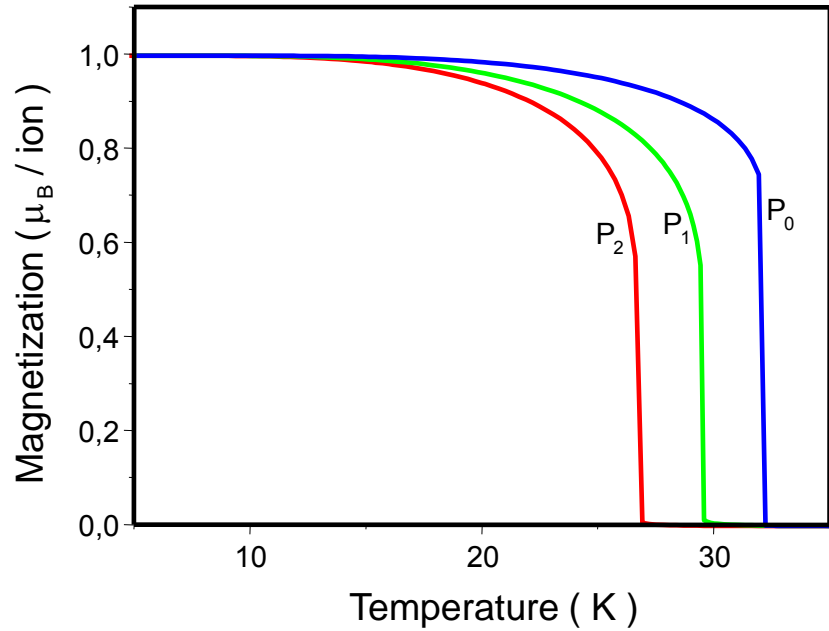


$$\Delta S_{iso}^{bar}(T, B, \Delta P) = S(T, B, P_2) - S(T, B, P_1)$$

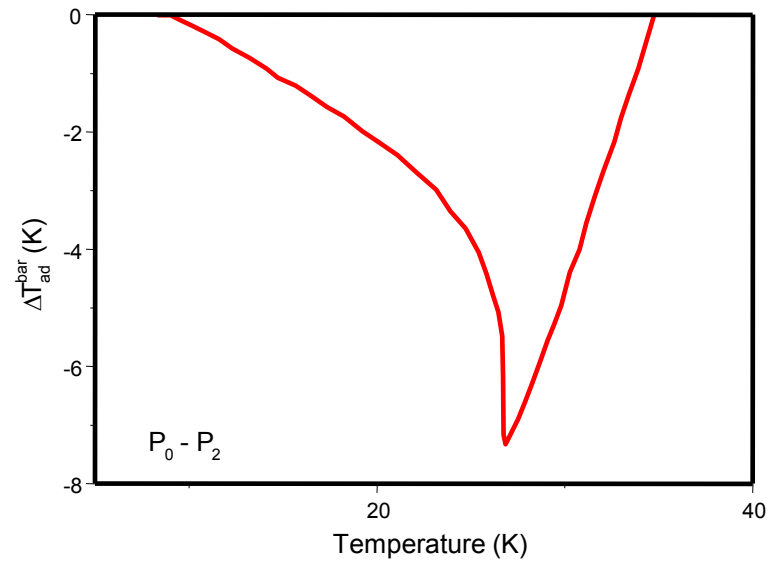
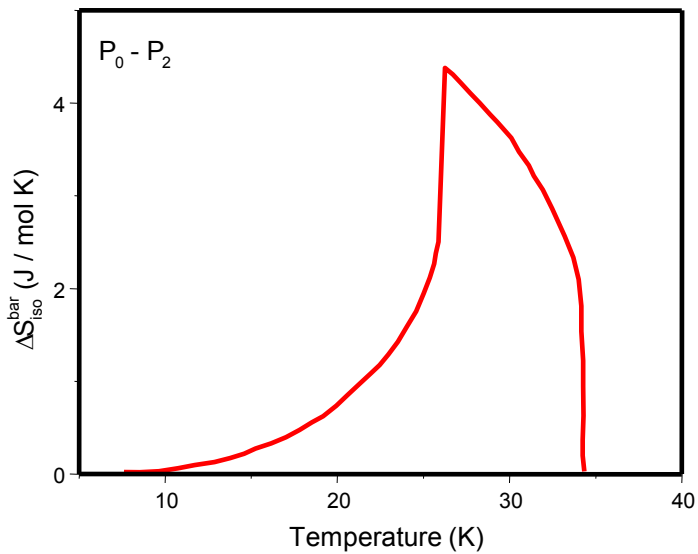
$$\Delta T_{ad}^{bar}(T, B, \Delta P) = T_1 - T_2$$

$$S(T, B, P_2) = S(T, B, P_1)$$

Systematic analysis: Scenario 1

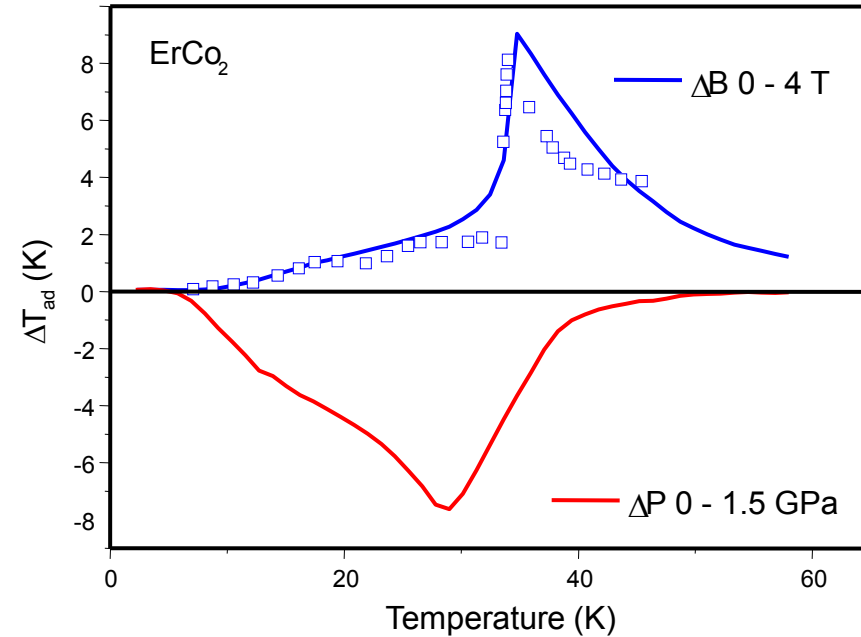
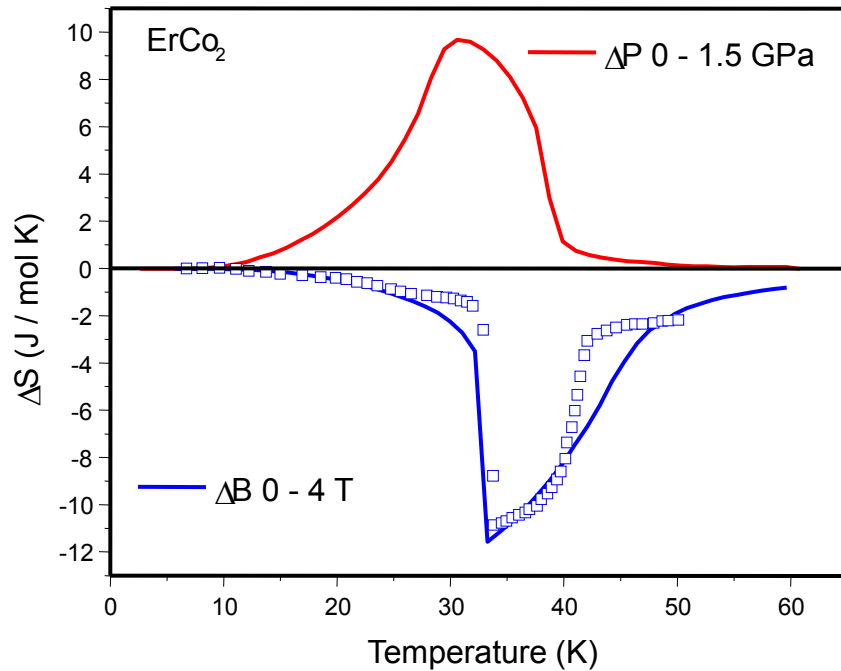


INVERSE BAROCALORIC EFFECT



BCE - ErCo₂

INVERSE BAROCALORIC EFFECT

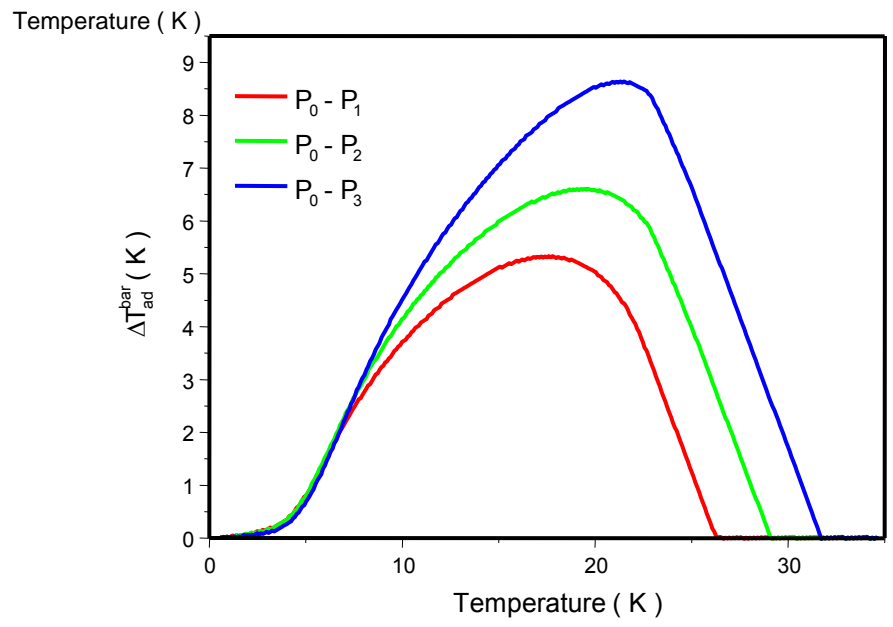
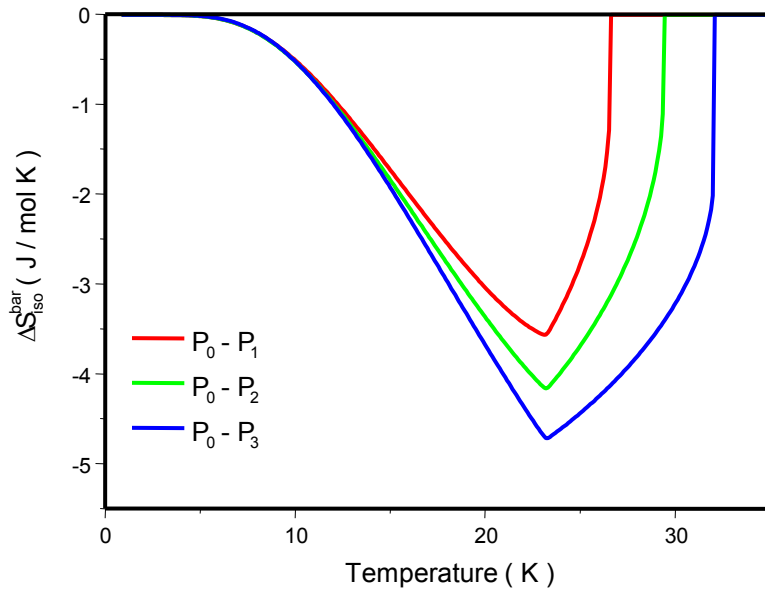
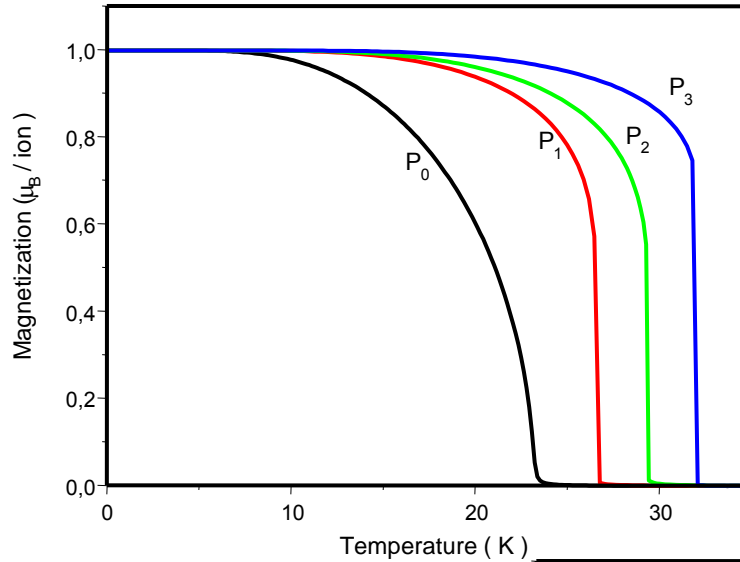


NORMAL MAGNETOCALORIC EFFECT

Wada et al, *Cryogenics* 39 (1999) 915

N. A. de Oliveira, *Journ. Appl. Phys.* 70 (2007) 052501

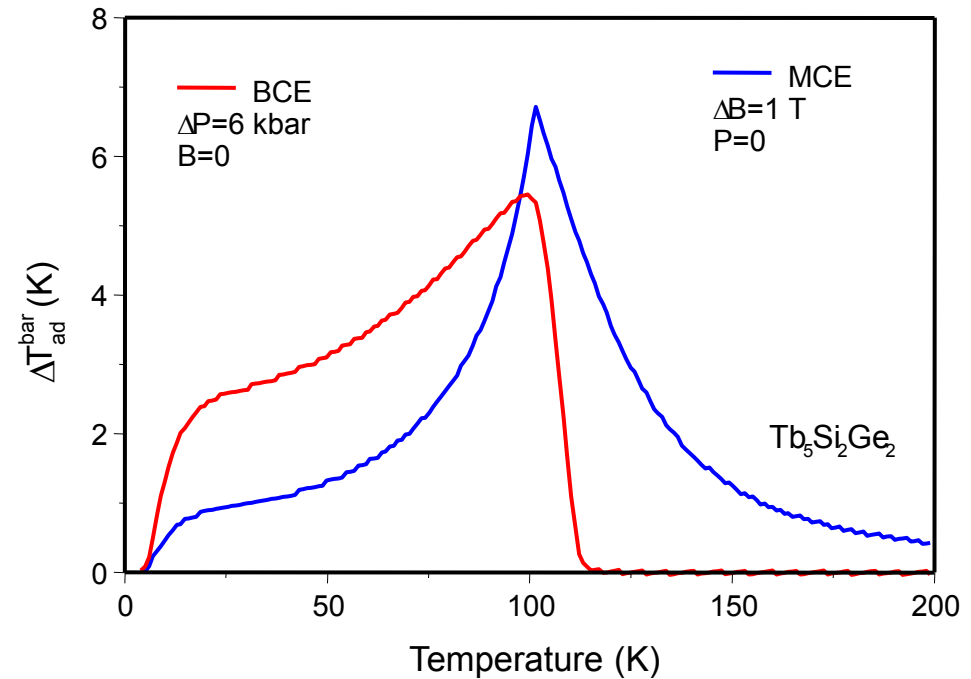
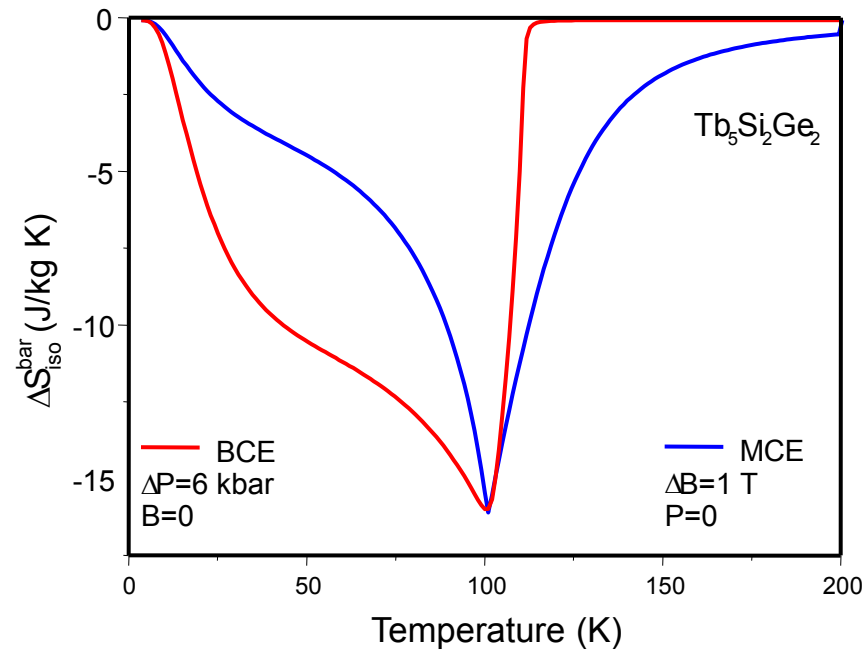
Systematic analysis: Scenario 2



NORMAL BAROCALORIC EFFECT

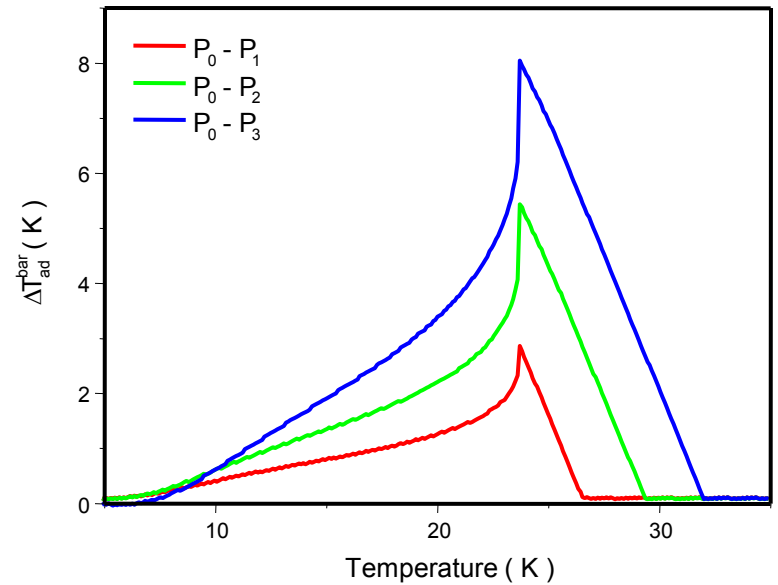
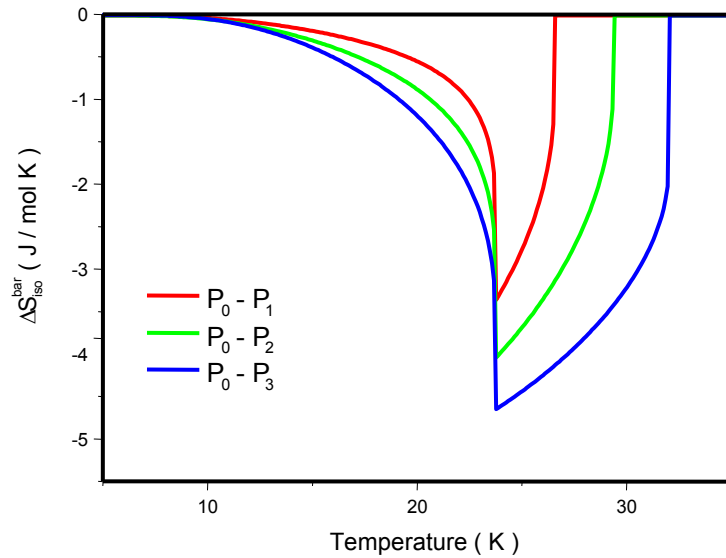
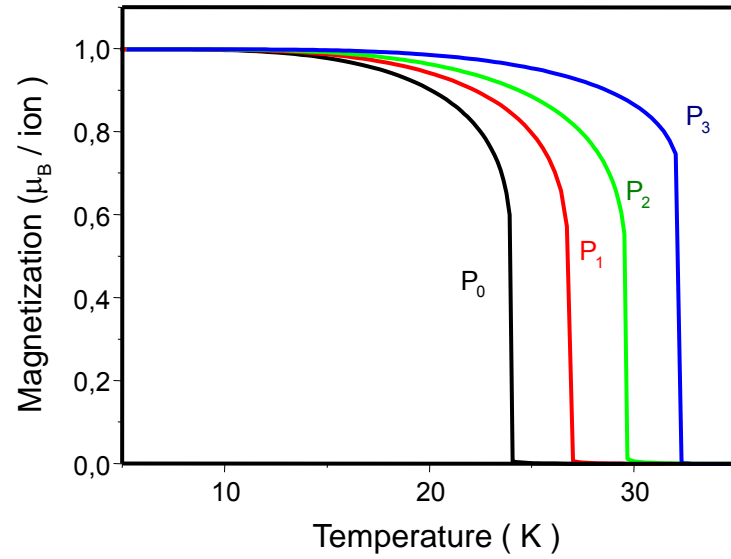
BCE - $\text{Tb}_5\text{Si}_2\text{Ge}_2$

NORMAL BAROCALORIC EFFECT

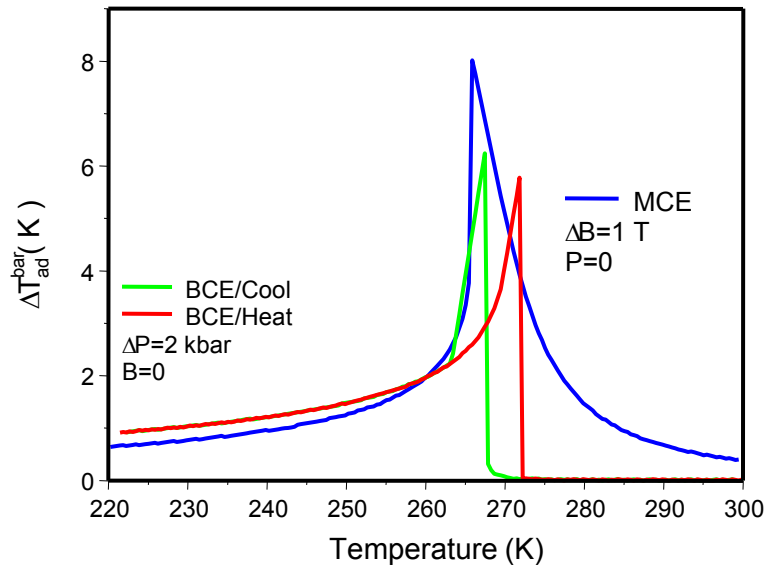
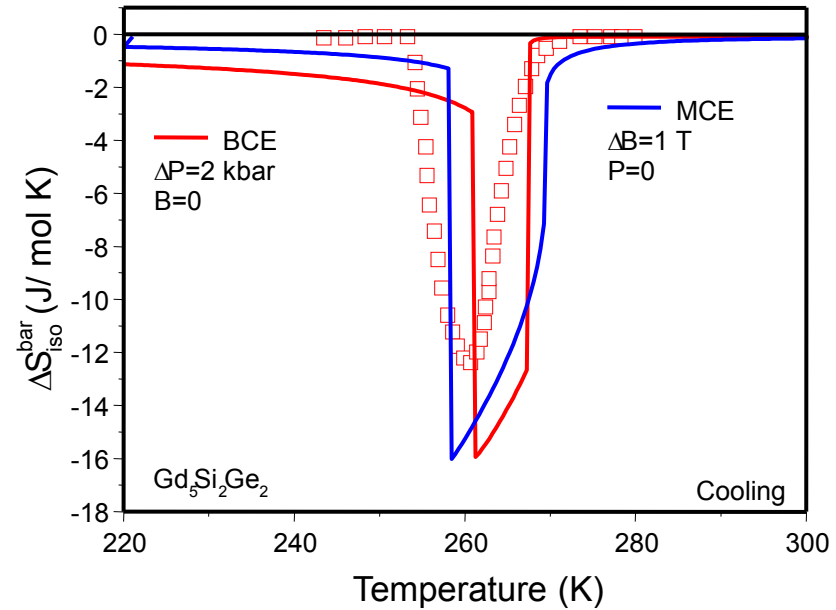
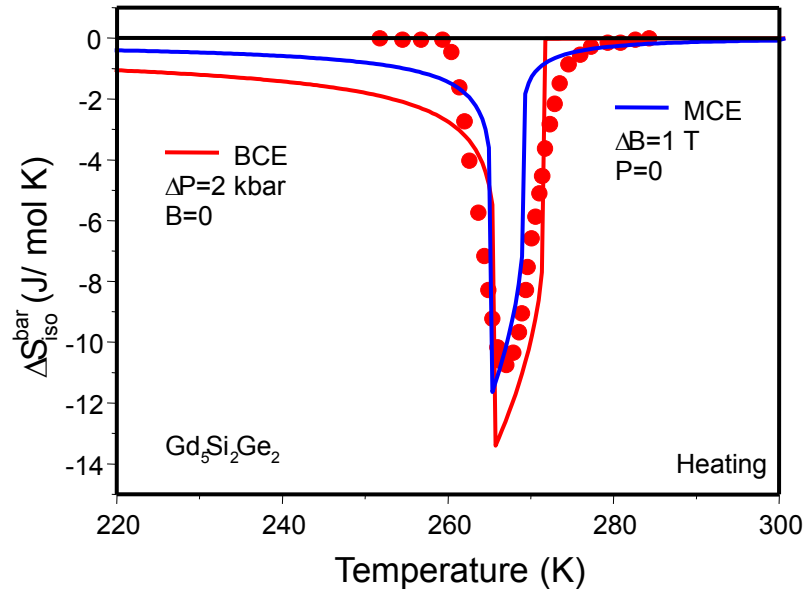


NORMAL MAGNETOCALORIC EFFECT

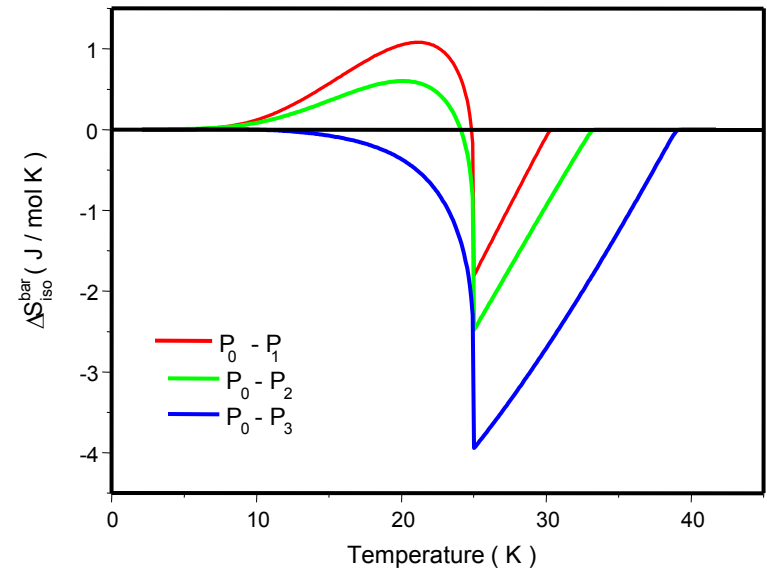
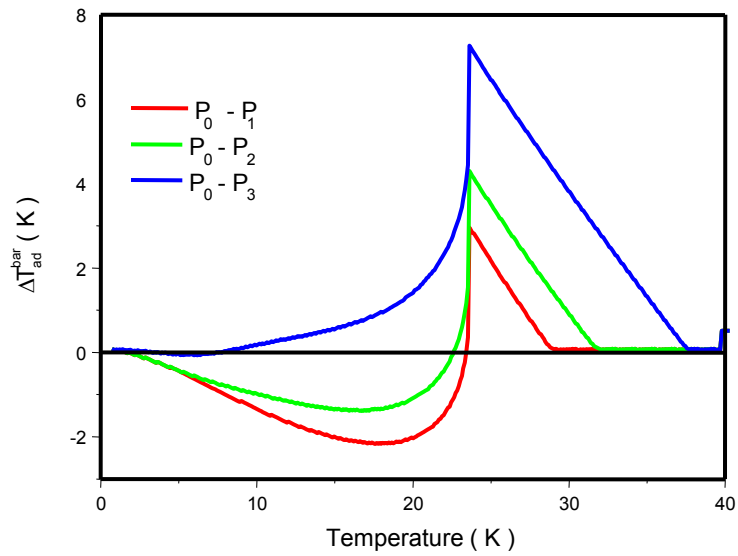
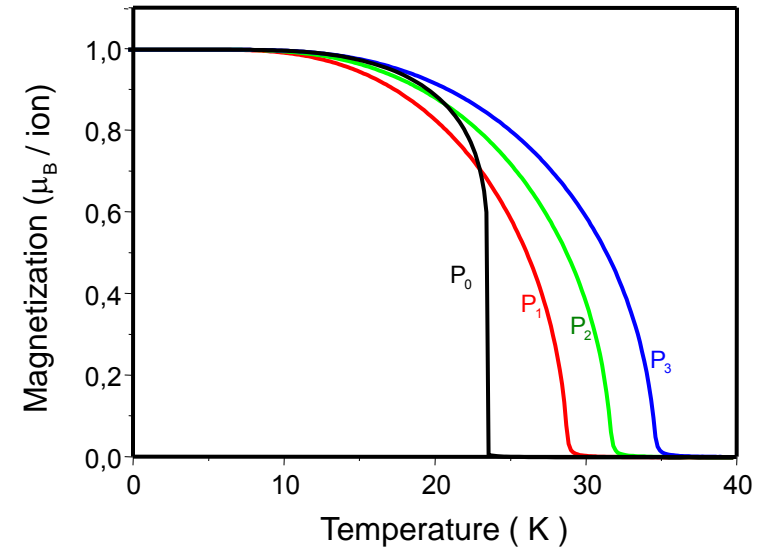
Systematic analysis: Scenario 3



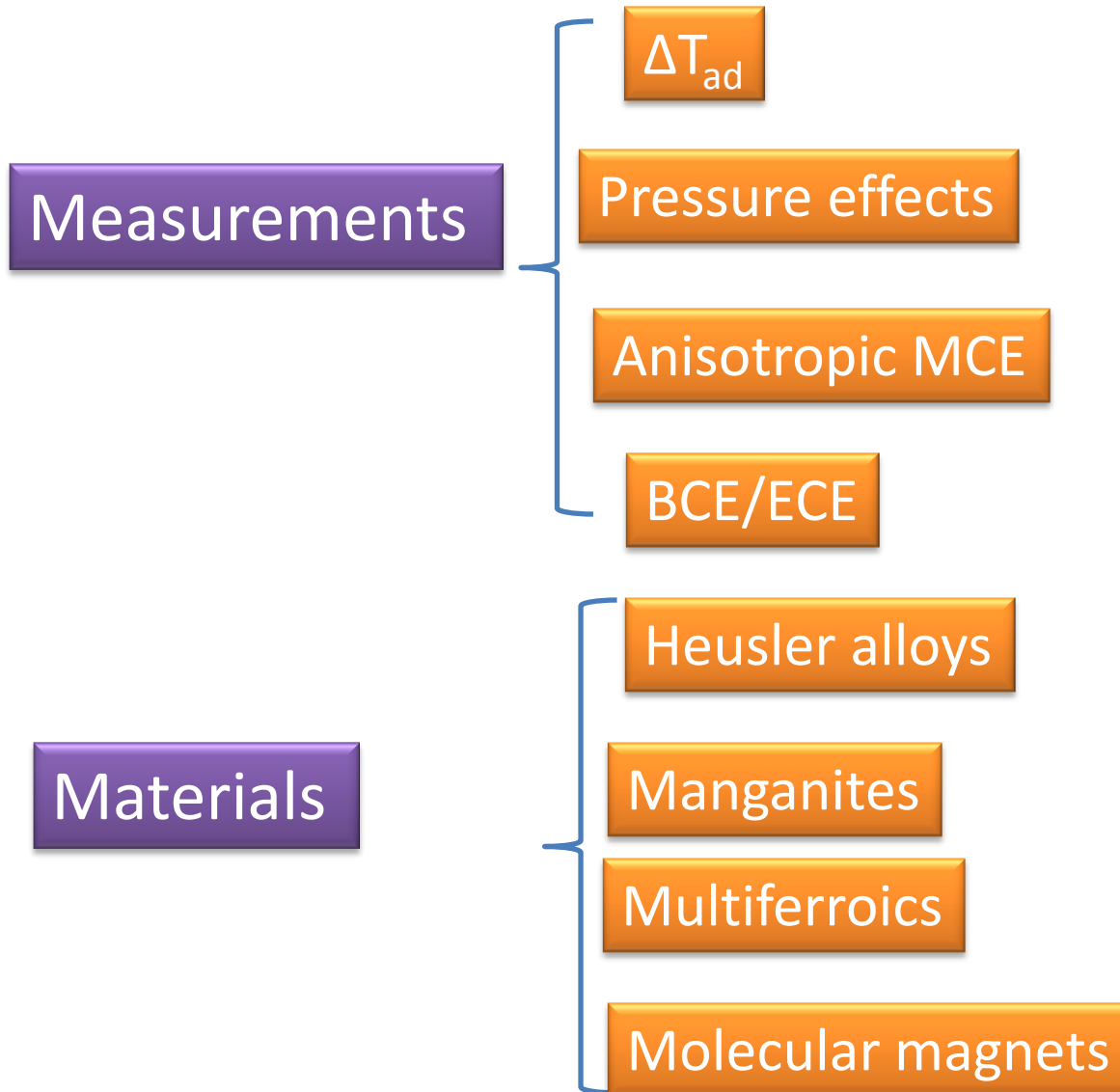
APPLICATION: $\text{Gd}_5\text{Si}_2\text{Ge}_2$



Systematic analysis: Scenario 4



PROSPECTS (Experimental)



PROSPECTS (THEORY)

Localized systems

Electron-phonon coupling

Monte Carlo simulations

Disordered effects

First order transition

Anisotropy

Itinerant systems

Electron-phonon coupling

Multi-band calculations

Antiferromagnetic systems

Beyond mean field

Field theory

QMC

ACKNOWLEDGEMENTS



FIRST ORDER MAGNETIC PHASE TRANSITION

ALTERNATIVE CALCULATION

ALTERNATIVE CALCULATION

Gibbs energy

$$G = F_{mag} + E_{mag}$$

$$E_{mag} = \int M \cdot dB^{eff}$$

Effective field

$$B^{eff} = B + \frac{\mathcal{J}_0 \langle J \rangle + \mathcal{J}_1 \langle J \rangle^3}{g\mu_B}$$

$$M(T, B, P) = g\mu_B \langle J \rangle$$

$$B^{eff} = B + \left[\frac{\mathcal{J}_0}{g^2\mu_B^2} \right] M + \left[\frac{\mathcal{J}_1}{g^4\mu_B^4} \right] M^3$$

$$B^{eff} = B_0 + \lambda_0 M + \lambda_1 M^3$$

$$dB^{eff} = (\lambda_0 + 3\lambda_1 M^2) dM$$

Magnetic energy

$$E_{mag} = \frac{\lambda_0 M^2}{2} + \frac{3\lambda_1 M^4}{4} + F_0$$

$$G = F_{mag} + \frac{\lambda_0 M^2}{2} + \frac{3\lambda_1 M^4}{4} + F_0$$

ALTERNATIVE CALCULATION

Energy

$$G = F_{mag} + \frac{\lambda_0 M^2}{2} + \frac{3\lambda_1 M^4}{4} + F_0$$

Derivative

$$B^{eff} = B_0 + \lambda_0 M + \lambda_1 M^3$$

$$\left(\frac{\partial G}{\partial M}\right)_T = \left(\frac{\partial F_{mag}}{\partial B^{eff}}\right)_T \overbrace{\left(\frac{\partial B^{eff}}{\partial M}\right)_T}^{\lambda_0 + 3\lambda_1 M^2} + \lambda_0 M + 3\lambda_1 M^3$$

Algebra

$$\overbrace{\left(\frac{\partial G}{\partial M}\right)_T}^0 = \left(\frac{\partial B^{eff}}{\partial M}\right)_T \overbrace{\left[\left(\frac{\partial F_{mag}}{\partial B^{eff}}\right)_T + M\right]}^0$$

$$M = - \left(\frac{\partial F_{mag}}{\partial B^{eff}}\right)_T$$

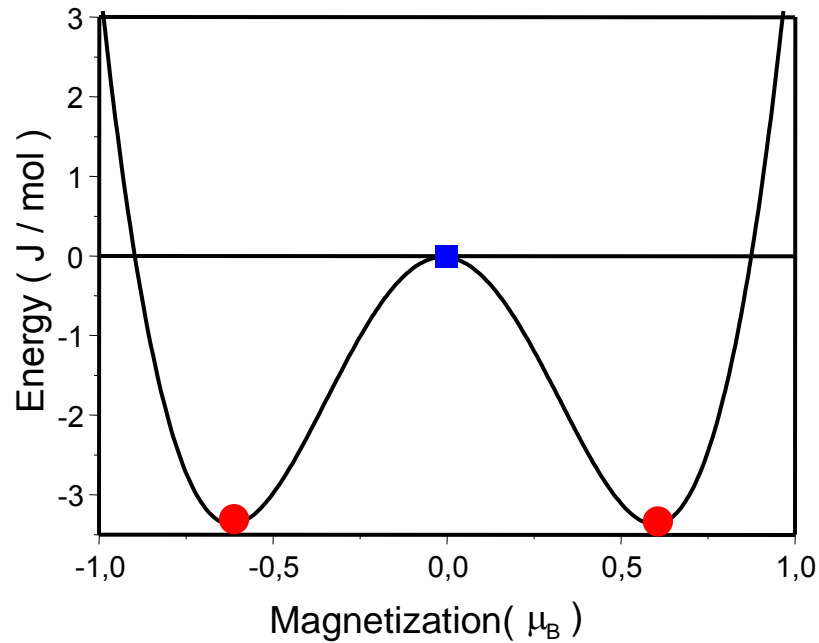
ALTERNATIVE CALCULATION

Free energy

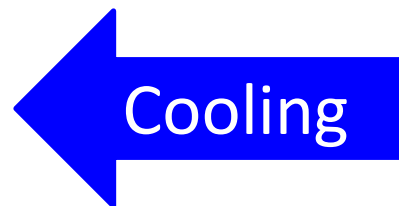
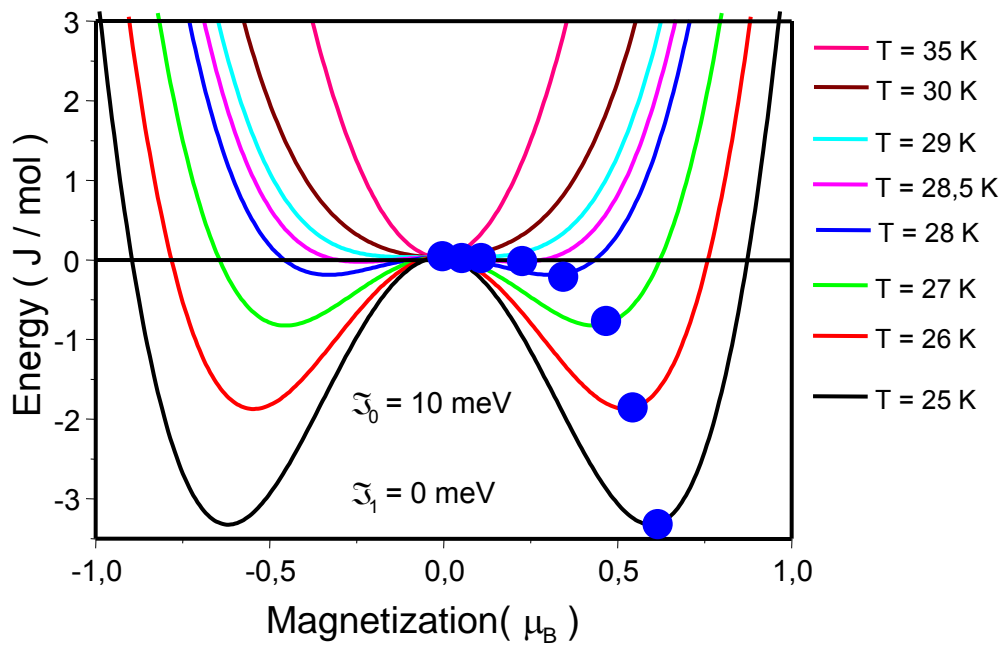
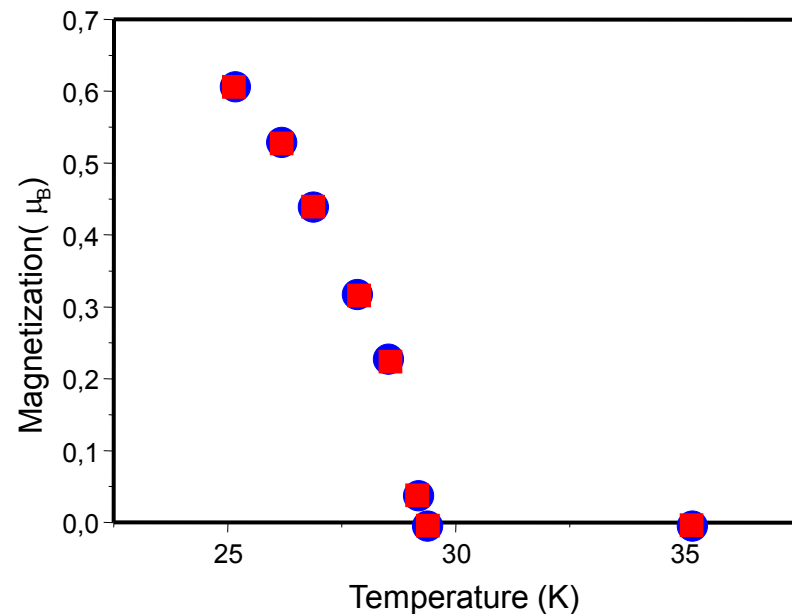
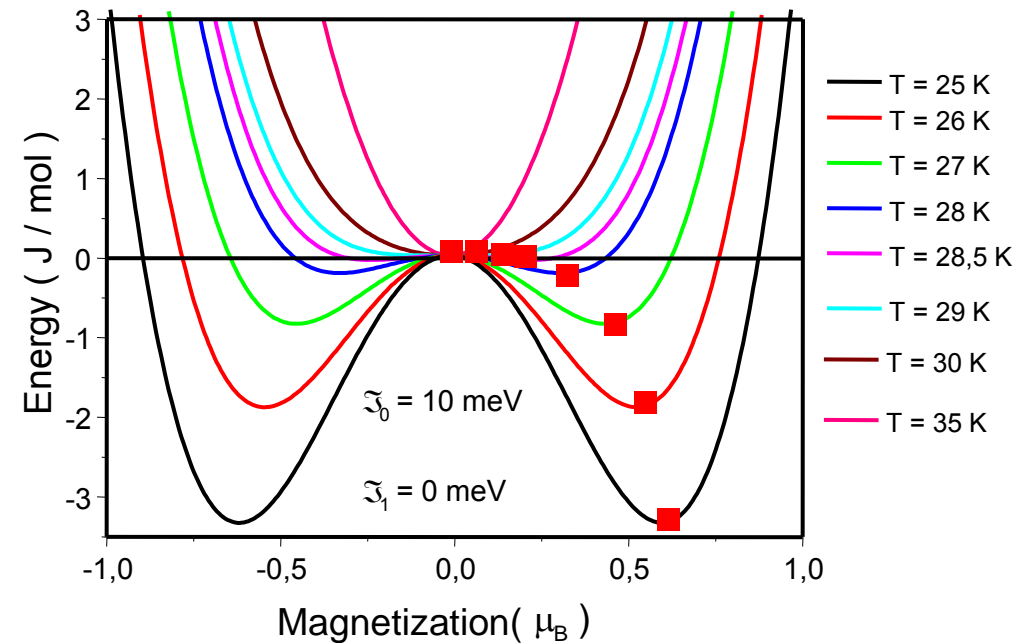
$$G = F_{mag} + \frac{\lambda_0 M^2}{2} + \frac{3\lambda_1 M^4}{4} + F_0$$

$$\frac{\partial G}{\partial M} = 0$$

$$M = - \left(\frac{\partial F_{mag}}{\partial B^{eff}} \right)_T$$



SECOND ORDER TRANSITION



FIRST ORDER TRANSITION

