## Polarized Raman spectra of L-arginine hydrochloride monohydrated single crystal

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Polarized Raman spectra of *L*-arginine hydrochloride monohydrated single crystal in nine different scattering geometries of the two irreducible representations of factor group  $C_2$  were studied at room temperature. The experimental wavenumber values are compared with those obtained from ab-initio calculation and the assignment of the Raman bands to the respective molecular vibrations is also given. Finally, a discussion related to a previously reported phase transition undergone by *L*-arginine hydrochloride monohydrated single crystal at low temperature is furnished.

Keywords:

# 1. INTRODUCTION

Amino acids  $(NH_2 - CH - COOH - R)$ , where *R* is a radical) are the basic units of proteins and peptides of all living beings. For an unknown reason Nature has choose 20 of these special molecules, differing in the *R* part, to form the impressive number of proteins found in our planet. *L*arginine, among other amino acids, and its salts, including *L*arginine·*HCl*, is known to protect protein from inactivation in frozen solutions, during freeze-drying [1], during spraydrying [2] and in the storage of lyophilized solids [3]. Beyond this, *L*-arginine and acid combinations have been extensively used in the last years to assist in the recovery of chemically unfolded proteins and recombinant proteins expressed in inclusion bodies [4].

Beyond these biological aspects related to amino acid in general, and *L*-arginine in particular, it is also possible to found interesting physical properties related to this amino acid. For example, the search for new non-linear optical materials with high non-linear coefficients and high damage threshold has lead to the discovery of *L*-arginine phosphate monohydrate (LAP) [5]. These studies also allowed the discovery of new materials having *L*-arginine as the main substance. Among these material is *L*-arginine hydrochloride monohydrate (LAHW) whose several physical properties as thermal expansion, elastic properties and dielectric characteristics were presented in Ref. [6].

Some works deal with the temperature behavior of *L*-arginine- $HCl \cdot H_2O$  monocrystals. It was observed that under high temperature conditions LAHW begins to lost water of crystallization at 70°*C* and at 200°*C* about two-third of it are eliminated [7]. Under low temperature conditions, on the other hand, it was observed through Raman scattering technique, evidence of a phase transition undergone by LAHW between 100 and 110 *K* [8]. Although in the paper of Ref. [8] some Raman spectra are presented, up to now, there is no

complete investigation of the polarized Raman spectra for LAHW, as well as, an assignment of the observed modes. Such assignment is very important because it helps us to understand what vibrations are related to the crystal modifications, in particular under temperature and pressure changes.

The objective of this paper is twofold: (i) to present the Raman spectra in the spectral region  $20 - 3700 cm^{-1}$  for nine scattering geometries of the two irreducible representations of the  $C_2$  factor group; (ii) to give a tentative assignment of the observed Raman bands based on *ab-initio* calculations.

#### 2. EXPERIMENTAL

Single crystals of *LAHCL* were grown from aqueous solution containing *L*-arginine hydrochloride powder,  $C_6H_{14}N_4O_2 \cdot HCl$ , from Sigma by the slow evaporation method at controlled temperature (293*K*). In order to characterize and to do the orientation of the crystals we used X-ray diffraction patterns obtained from a Rigaku DMAX diffractometer using *Cu K*\alpha radiation monochromated with a graphite crystal. The Raman spectra were obtained with a Jobin Yvon T64000 micro-Raman system equipped with an  $N_2$ -cooled charge coupled device system. The samples were polished with diamond paste with granulations with 10, 3 and 1µm. The slits were set for a  $2cm^{-1}$  spectral resolution. Excitation was effected with the 514.5*nm* radiation from an argon ion laser. The incident laser had a power less than 5*mW* on the surface of the sample.

## 3. RESULTS AND DISCUSSION

Fig. 1 presents the zwitterion form of an arginine molecule (the number associated to the atoms will be used in the assignment of the normal modes). Single crystal data of LAHW crystal confirms that the compound grows with a monoclinic lattice belonging to the  $P2_1(C_2^2)$  space group with two molecules of arginine ( $C_6H_{14}O_2N_4$ ), two units of *HCl* and two units of  $H_2O$  per unit cell. In the primitive cell

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**Table 1.** Wavenumbers  $(cm^{-1})$  and tentative assignment for the bands appearing in the Raman spectra of LAHW crystal at room temperature. The notation a(bc)d is the conventional Porto notation, where *a* and *d* represent the directions of the incident and the scattered light and *b* and *c* represent the directions of the incident and the scattered light.

v (cm-1) Calculate	Base Set		scattering geometries								
	6-311++G**	6-311G	x(yy)x	z(yy)z	y(xx)y	z(xx)z	x(zz)x	y(zz)y	x(yz)z	z(yx)z	y(zx)y
44					43	43	41	43	42	43	
48	τ (N17-C19)										
57			59	59	60	59	58	59	58	59	61
61	r(NH <sub>3</sub> <sup>+</sup> ) ip		63	63		63	61	63	62	63	
			68	68	71	71	69	70	68	70	71
78		r(H16-C15-H18)	75	75	85	75	74	75	74	75	77
86	r(H16-C15-H18)			87		83	85	87		86	
			94		96	96	95	96	95	96	90
			105	102	102	102	101	102	105	102	100
						109	106	106			106
			120	120	118	118	117	119	118	119	119
131		r(H13-C12-H14)									132
137	r(H13-C12-H14)		142	142	140	141	138	139	141	142	
							149	151			152
163	w(C15-N17-C19)	r(N1-H7-H26)	158	159	159	162	156	158	158	162	
			185	190	189	189	189	192		189	191
237	$\tau (\mathrm{NH_3^+})$		244	245	243	245	243	243	243	244	
261	r(H10-C9-H11)										
262		r(H10-C9-H11)	268	273	269	274	267	270		273	
305	$\delta$ (C9-C12-C15)		307	307	306	307	303	303	305	307	303
308											
326											321
327	τ (skel.)										
334		$\tau (\mathrm{NH_{3}^{+}})$			331		331	332	335		
			340	340	341	340				340	340
389			377	378	378	379	377	377	380	378	
394	$\delta$ (C3-C2-N1)										
415										415	
			419	417			419	418	422		425
446	r <sub>1</sub> (NH <sub>2</sub> )	r <sub>1</sub> (NH <sub>2</sub> )	446	446		448	445	447	446	447	
457	r <sub>2</sub> (NH <sub>2</sub> )		466	466	465	465				465	470
513		δ (N17-H20)		529	531		532	533	533	530	535
545	δ (N17-C19-N22)		539			537				539	
557	v (C2-C3-N1)		551	553	550	551			552	552	
558											
567	∂ (C2-C9-C12)										
585	∂ (N17-C19-N21)		621	620			618	619	619	621	619
625									629		
670	v (C2-C3)		678	679		679	675	678	677	678	

of LAHW, all the atoms occupy sites with symmetry  $C_1$  and the 186 vibrations can be decomposed into the irreducible representations of the factor group  $C_2$  as  $\Gamma = 93A + 93B$ , with one *A* and two *B* belonging to the acoustic branch and all the other being Raman and infra red active. In Figs. 2 – 5, *Z* represents an axis parallel to the [001] direction, *X* represents an axis parallel to the [100] direction and *Y* is an axis perpendicular to the *X* and *Z* axes.

In order to give insights about the assignment of the normal modes of the crystal we have performed *ab initio* calculations. All calculations were carried out using Gaussian98 (by computational resources in CENAPAD's facilities) and the results were viewed by MOLEKEL programme packages. Geometry optimization and frequency calculation for *L*-arginine in the gas phase were performed with Self-Consistent methods and Hartree-Fock (HF) level of theory, with 6-31+G(d,p) and 6-311++G(d,p) basis sets. The molecular framework was designed in Gchempaint (gnome chemical software packages) and the atomic coordenates of isolated *l*-arginine was used in input file. This structure was optimized using  $\delta E < 10^{-8}$  as convergence parameter and the vibrational wavenumbers were then calculated. The out-

v (cm-1)	Base	scattering geometries									
Calculate	6-311++G**	6-311G	x(yy)x	z(yy)z	y(xx)y	z(xx)z	x(zz)x	y(zz)y	x(yz)z	z(yx)z	y(zx)y
692			701	702		702	700	701	700	702	
717	w(N21-C19-N17-										
732	w(N17-C19-N22)		751	753	751				748	753	
786	τ (skel.)		784	784	782	782			783	783	777
799											
829											
831				831	831	830				830	
846	$\delta(CO_2)$		843	843	843	843	842	843	842	843	841
885	• (• • 2)										
890		δ (N21-H23)		892	897	898	890	892	890	892	891
894	r(NH <sub>2</sub> <sup>+</sup> ) op	0 (1121 1123)	897	898	0,7,7	0,0	897	898	897	898	898
900	1(1113) / Op		0,7,7	0,0			0,7,7	0,0	0,7,7	0,0	0,0
902											
904											
929	w(NH <sub>2</sub> )		930	931	931	930	929	930	930	930	930
)2)	w(1112)		930	940	030	030	030	930	930	940	940
050			055	740	,3,	,3,	053	055	,37	240	740
950			955	060	050	050	955	955	058	050	057
937	# '(NILL <sup>+</sup> ) on		900	900	939	939	938	939	938	939	937
1001	1 <sub>1</sub> ( <b>NH</b> <sub>3</sub> ) Op										
1006			1020	1020	1027	1029	1028	1020	1020	1021	1020
			1029	1029	1027	1028	1028	1029	1029	1021	1029
1052	v(C2 N1)		1049	1049	1049	1049	1049	1049	1049	1050	
1033	V(C2-N1)		1032	1055	1039	1052	1000	1000	1032	1057	1064
1073	- (NIL <sup>+</sup> )		1020	1002	1002	1000	1000	1000	1002	1001	1004
1096	$r_1(NH_3)$ op		1089	1092	1092	1090	1090	1090	1093	1091	1101
1112		S AN 1120	1099	1100	1099	1099	1100	1101	1099	1099	1101
1112		∂ <sub>1</sub> (N1-H26)		1122	1124	1125				1126	1122
1120				1132	1134	1135				1136	1133
1120	$r_2'(NH_3')$ op										
1144		w(N1-H5-H26)	1139				1138	1138	1140		1142
1163	v(C19-N22)										
1172	a ma h	∂ 2(N1-H26)									
1180	$r_2(NH_3)$ op		1179	1181	1177	1179	1176	1178	1179	1180	1179
1211	r <sub>1</sub> "(NH <sub>3</sub> <sup>-</sup> ) op	w ( $NH_3^{-}$ )									
1231							1231	1232			
1247	$r_2$ "(NH <sub>3</sub> <sup>-</sup> ) op		1253				1250	1250	1253		1253
1257	$t_1(NH_2)$										
1259		r <sub>2</sub> (NH <sub>2</sub> )									
1275		_									
1282		$\delta$ (N21-H23)									1291
1294	$t_2(NH_2)$			1200	1297	1298	1297	1298	1297	1298	
			1212	1509	1212	1212	1212	1212	1212	1212	1212
			1312	1226	1312	1512	1312	1313	1312	1313	1313
			1326	1326	1325	1526	1524	1325	1326	1526	1326

put file contained the optimized structure, the Raman and IR intensities and frequencies, and the atomic displacements for each mode. At the optimized structure of the molecule, no imaginary frequency was obtained, proving that a true minimum of the potential energy surface was found.

In all spectra presented in this paper we use the conventional Porto notation a(bc)d, where a and d represent the directions of the incident and the scattered light and b and c represent the directions of polarizations of the incident and the scattered light. The scattering geometries in all figures are related to the irreducible representations of the factor group  $C_2$  as follows: x(yy)x, z(yy)z, y(xx)y, z(xx)z, x(zz)x, y(zz)y and z(yx)z are from *A* irreducible representation and x(yz)x and y(zx)y are from the *B* irreducible representation.

Figure 2 presents the Raman spectra of LAHW for nine different scattering geometries in the spectral region  $25 - 225 \ cm^{-1}$  at room temperature. This region is known to have the lattice modes of the crystal ( $\omega < 200 \ cm^{-1}$ ), although some internal modes can also be observed with low wavenumber, as occurs with the torsional vibration of  $CO_2$ ,  $\tau(CO_2)$ , for *L*-asparagine monohydrated [9], *L*-valine [10] and *L*-isoleucine [11].

v (cm-1) Calculate	Base Set		scattering geometries								
	6-311++G**	6-311G	x(yy)x	z(yy)z	y(xx)y	z(xx)z	x(zz)x	y(zz)y	x(yz)z	z(yx)z	y(zx)y
1337	t (CH <sub>2</sub> )		1341	1342	1341	1342	1340	1341	1340	1342	
1344		t (CH <sub>2</sub> )	1349	1349	1349	1349	1348	1349	1348	1349	
											1358
			1371	1371	1370	1371			1370	1371	1389
1388		v (C2–C3)	1388	1389	1389		1390	1391		1389	
1404	$\delta$ (CH <sub>2</sub> )		1406		1406	1406	1405	1406	1405	1406	1409
			1439	1440	1440	1440	1438	1440	1439	1440	1440
			1450	1450	1449	1450	1449	1450	1449	1450	
			1456	1456	1456	1456	1455	1456	1455	1456	1457
1474	r(CH <sub>2</sub> )		1472	1474		1472	1471		1472	1473	
1534		w(CH <sub>2</sub> )	1535	1535							1535
1572		$\delta$ (NH <sub>3</sub> )op		1567	1565	1568	1564			1564	1566
1588	$\delta$ (H-N-C)	$\delta$ (H-N-C)			1591						
				1605		1605				1604	1607
1613	$\delta$ (H10-C9-H11)						1614	1615			
1623	δ (H14-C12-		1619		1622	1624			1621	1625	1621
1.550	H13) δ (H18-C15-	δ(H14-C12-	1654	1655		1655	1655	1657	1652	1655	
1650	H16)	H13)	2853	2854	2855	2854	2854	2854	2854	2854	2855
			2879	2879	2879	2879	2879	2878	2879	2879	2878
			2890	2801	2891	2801	2890	2890	2890	2801	2891
			2000	2001	2001	2001	2005	2000	2000	2001	2004
			2917	2918	2917	2917	2918	2917	2918	2918	2917
			2944	2947	2943	2945	2940	2942	2944	2944	2944
			2953	2954	2952	2745	2953	2742	2953	2953	2744
			2962	2962	2962	2962	2962	2962	2962	2962	2960
			2970	2970	2970	2971	2970	2970	2970	2970	2971
			2991	2991	2991	2991	2991	2991	2991	2991	2991
21(0			2001	2777	2001	3176	2001	2001	2001	2001	2//1
5109		v (C12–H14)	3179			5170	3183	3183	3181	3183	3183
		v (H16-C15-		3188	3193						
3194		H18)									
			3335	3332			3335	3334	3336	3335	3332
					3347	3344					
				3392	3391					3390	3398
				3482	3482	3481				3483	3483

Legend:  $\delta$  = bending,  $\nu$  = stretching,  $\tau$  = torsion, w = wagging, t = twist, r = rocking, ip = in plane, op = out of plane, skel. = framework vibration; subscripts indicate degenerence; superscript indicate same vibration for different atoms in same cluster. Atom numbering is shown in Figure 1.



FIG. 1: Molecule of arginine in the zwitterion form.

In fact, the ab initio calculations performed in the present work show the occurrence of vibrations with  $\omega < 200 cm^{-1}$ that can be associated to internal modes. As an example, using the 6-311++G\*\* base set, it is possible to assign the in-plane torsion of  $NH_3^+$ ,  $\tau(NH3^+)ip$ , as the mode appearing with wavenumber of 63  $cm^{-1}$ .

Related to the torsional vibration of  $CO_2^-$ ,  $\tau(CO_2)$ , our calculations were not able to identify it precisely. However, it is possible to assign the band observed at ~ 189  $cm^{-1}$ as the  $\tau(CO_2)$ . This is based on the studies performed on other hydrogenated amino acid crystals [9 – 11] as well as in deuterated one [12]. Such an assignment is important because under extreme conditions torsional vibrations of  $CO_2$ can present particular behavior and furnish insights about the hydrogen bonds, as those presented by *L*-alanine-d7 under pressure variation [13].

Fig. 3 presents the Raman spectra of LAHW for nine different scattering geometries in the spectral region  $200 - 650 \text{ cm}^{-1}$ . In this region it is found the bands associated to torsional vibration of  $NH_3^+$ ,  $\tau(NH_3^+)$ . In general, the band associated to this vibration has a low intensity, at least as



FIG. 2: Raman spectra of *L*-Arginine hydrochloride monohydrate crystal in several scattering geometries in the  $25 - 225 \text{ cm}^{-1}$  spectral region.

suggested by former works on *L*-alanine [12] and *L*-leucine [14]. Other important vibration is associated to the rocking of  $NH_2$ ,  $r(NH_2)$ , which ab initio calculations show to be between 445 and 465  $cm^{-1}$ . Bending vibration of NH are also found in this spectral region, at about 530  $cm^{-1}$ , while the bending of N-C-N,  $\delta(NCN)$ , was identified as the peak at  $\sim 551 \ cm^{-1}$ , a band with small linewidth as can be observed in Fig. 3 for some scattering geometries.

Fig. 4 presents the Raman spectra of LAHW in the spectral region  $650 - 1750 \ cm^{-1}$  for nine different scattering geometries. This region presents a series of bands and the ab initio calculation was fundamental to assign in a most precise manner. The peak observed at about  $678 \ cm^{-1}$  was associated to a C2 - C3 stretching vibration, v(C2 - C3). Here we need to state the following observation: both, C2 and C3 carbon atoms are hold with other atoms ( $NH_3$  for C2 and  $O_2$  for C3) in such a way that we do not have a C - C "pure" vibration, as occurs with diamond; this explains why for diamond the v(C - C) vibration is observed at 1332  $\ cm^{-1}$ . A wagging vibration related to the N - C - N unit was associated to the band observed at 751  $\ cm^{-1}$ . The bending of  $CO_2^-$  unit,



FIG. 3: Raman spectra of *L*-Arginine hydrochloride monohydrate crystal in several scattering geometries in the  $200-650 \text{ } \text{cm}^{-1}$  spectral region.

 $\delta(CO_2^-)$ , was observed at 843  $cm^{-1}$ , as pointed out by our calculations. The bands observed at about 930 and 940  $cm^{-1}$  were associated to the wagging vibration of  $NH_2$ ,  $w(NH_2)$ . The stretching vibration of C2 - N1, v(C2 - N1), was associated to the band observed at ~ 1052  $cm^{-1}$ , and an out-of-plane  $NH_3^+$  rocking vibration,  $r(NH_3^+)$ , was assigned as the band observed at about 1090  $cm^{-1}$ . An out-of-plane  $NH_3^+$  vibration was assigned as the band observed at ~ 1250  $cm^{-1}$ .

Many bands are observed between 1300 and 1400  $cm^{-1}$ , among them the twist vibration of  $CH_2$ ,  $t(CH_2)$ . It is worth to note that the 6-311G basis of calculation suggests that the C2 - C3 stretching vibration should be observed at 1388  $cm^{-1}$ . However, because the reason given previously (carbon atoms bonded to other atoms) we believe that the suggestion is not correct. The band observed at 1406  $cm^{-1}$ was associated with the bending vibration of  $CO_2^-$ ,  $\delta(CO_2^-)$ , and a rocking vibration of the same unit,  $r(CO_2^-)$ , was associated to the band observed at  $\sim 1472 cm^{-1}$ . Finally, in this spectral region we also observed bending vibrations of the H - N - C unit, between 1600 and 1650  $cm^{-1}$ .



FIG. 4: Raman spectra of *L*-Arginine hydrochloride monohydrate crystal in several scattering geometries in the  $650 - 1750 \text{ cm}^{-1}$  spectral region.

Fig. 5 presents the Raman spectra of LAHW crystal in the spectral region  $2800 - 3600 \ cm^{-1}$  for nine different scattering geometries. The profile between 2800 and 3000  $\ cm^{-1}$  is very rich, originated mainly from stretching vibrations of *CH*. However, our calculations were not able to assign separately all bands in this region. Finally, the bands with high linewidths centered at ~  $3200 \ cm^{-1}$  and  $3400 \ cm^{-1}$  can be associated to the stretching vibration of water, v( $H_2O$ ), as it is expected from crystals with structural water [15].

With this assignment we can both (i) through light on the phase transition undergone by LAHW at low-temperatures and (ii) to give insights on future works on vibrational properties of the crystal under high pressure conditions. Related to the first point we remember that Raman spectroscopy study have suggested the occurrence of a phase transition undergone by LAHW crystal at about 110 K [8]. This phase transition was inferred mainly by the change of band profiles in the low-wavenumber region of the spectra. Obviously, as stated previously, these modifications are associated with lattice mode vibrations. Interesting enough is that we have also observed modifications in bands observed at  $1088 - 1100 \ cm^{-1}$ . The assignment of the present study confirms the identification of such bands as rocking vibrations of NH<sub>3</sub> units. In other words, our study reinforces the interpretation that the low temperature phase transition undergone by LAHW can be understood mainly as conformational change of the L-arginine molecules, mainly consequence of modifications on hydrogen bonds  $N - H \dots O$ , with the oxygen atom belong to a water molecule or to the COO- group



FIG. 5: Raman spectra of *L*-Arginine hydrochloride monohydrate crystal in several scattering geometries in the  $2800 - 3600 \text{ cm}^{-1}$  spectral region.

of another arginine molecule. This is in good accordance with experimental facts already presented by *L*-cisteine [16] and *L*-serine [17] crystals when submitted to high pressure conditions – the change of molecular conformation, due to variations in the dimensions of hydrogen bonds, modifies the symmetry of the unit cell of the crystal.

### Conclusions

In this paper we presented the polarized Raman spectra of LAHW in the spectral region between 25 and  $3600 \text{ cm}^{-1}$ . Ab initio calculations were used to assign most of the bands observed in all nine scattering geometries. With such identification we gave insights about a previously observed phase transition undergone by LAHW at low temperatures. Finally, the work will be useful to understand eventual modifications in the Raman spectra of LAHW, for example, when submitted to high pressure conditions in future works.

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