# Characterization of PbTe Epitaxial Layers Grown on BaF<sub>2</sub>/CaF<sub>2</sub>/Si Structures

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> > Received July 21, 1995

In this paper we present a characterization of PbTe epilayers grown by hot wall epitaxy on silicon using II-a fluorides as intermediate layers. The PbTe layers were characterized electrically by temperature dependent Hall effect measurements. A detailed study of the strain in CaF<sub>2</sub> layers grown on Si(111) substrates was performed by high resolution x-ray diffraction analysis. To simulate the real operation conditions of the lead chalcogenides infrared devices, the CaF<sub>2</sub>/Si structure was submitted to thermal cycles from 300 to 77K and a  $\omega$  scan was measured after each cycle.

#### I. Introduction

The interest in the epitaxy of lead chalcogenides on silicon has recently increased due to the possibility of fabricating IV-VI semiconductor compounds infrared devices integrated monolithically to the Si substrate. Since the epitaxial growth of narrow gap lead chalcogenides directly on Si presents problems due to the large differences in lattice constants and thermal expansion coefficients, researches have been carried out on the possibility to interpose II-a fluoride buffer layers between the two semiconductor materials<sup>[1,2,3]</sup>.

In the IV-VI compounds epitaxy on Si substrates, CaF<sub>2</sub> is normally chosen as the first buffer layer because of the small lattice mismatch relative to Si (0.5% at room temperature) followed by a BaF<sub>2</sub> layer which has a close match of lattice constant and thermal expansion coefficient to the lead chalcogenides ones. Although the CaF<sub>2</sub> lattice constant (5.46Å) is very close to Si (5.43Å), the large difference in thermal expansion coefficients between CaF<sub>2</sub> (19.2×10<sup>-6</sup> K<sup>-1</sup>) and Si (2.6 × 10<sup>-6</sup> K<sup>-1</sup>) imposes some limits in the epitaxial growth<sup>[4,5]</sup>. In order to obtain good monocrystalline lead chalcogenides epilayers on the  $BaF_2/CaF_2/Si$  structure, the strain relief mechanisms of the thermal mismatch in the  $CaF_2/Si$  system must be well understood.

In this work we report on the electrical and structural characterization of PbTe layers grown by hot wall epitaxy technique on  $BaF_2/(CaF_2/Si$  structure, as well as a detailed analysis of the strain situation in the  $CaF_2/Si$  system by high resolution x-ray diffraction measurements. To simulate the real condition of the operation of the infrared detectors, the  $CaF_2/Si$  system was submitted to thermal cycles from room to liquid nitrogen temperature and the x-ray rocking curves were measured after each thermal cycle.

#### **II.** Experimental procedure

The films were grown in a home made hot wall epitaxy (HWE) system<sup>[3]</sup> consisted of a high vacuum chamber with two HWE reactors for the lead telluride compounds, each one with an additional Te compensation source, and two graphite effusion cells for the fluorides. The background pressure during growth was  $10^{-7}$  mbar. The Si substrate is placed inside a furnace that fits on top of a rotator disc, that allows the growth of several layers without vacuum breaking. Due to the preferential (111) growth of the fluorides, Si(111) substrates were used. Details about the growth procedure were published elsewhere<sup>[3]</sup>. To measure the layer thickness perfilometry and ellipsometry methods were used for the fluoride layers and interferometry for the PbTe ones.

#### **III.** Electrical properties

To characterize the PbTe layer electrically Hall effect and resistivity measurements were made from room temperature down to 13K in an automatic temperature dependent Hall effect system using the Van der Pauw geometry.

The electrical properties of the PbTe layers grown on top of the BaF<sub>2</sub>/CaF<sub>2</sub>/Si structures were compared with PbTe layers grown directly on  $BaF_2(111)$  substrate. Fig. 1 shows the Hall mobility as a function of temperature for a typical p-type PbTe layer  $(1.5 \times 10^{17} \text{cm}^{-3})$  grown on BaF<sub>2</sub>/CaF<sub>2</sub>/Si structure and for a PbTe layer with the same carrier type and concentration grown directly on a BaF<sub>2</sub> substrate. For the layer  $PbTe/BaF_2$ , the mobility is limited at high temperatures by acoustic phonons ( $\mu \sim T^{-5/2}$ ) while at low temperatures some defects (or impurities) limits the mobility at  $3 \times 10^5 \text{cm}^2/\text{V.s.}$  In case of PbTe/BaF<sub>2</sub>/CaF<sub>2</sub>/Si the amount of defects start to limit the mobility already at approximately 150K causing a deviation from the phonon limited curve  $(\mu \sim T^{-5/2})$  and a saturation in the mobility at low temperatures at a value 10 times lower. These defects on the PbTe layer of the PbTe/BaF<sub>2</sub>/CaF<sub>2</sub>/Si structure are caused by dislocations created by the differences in lattice constants and thermal expansion coefficients between the semiconductor materials.

By adjusting very well the thicknesses of the fluoride buffer layers, specially by controlling the  $CaF_2$ layer thickness not to exceed 10 nm and optimizing the growth parameters it was possible to obtain a n-type PbTe layer on top of the  $BaF_2/CaF_2/Si$  structure with mobility values of  $5 \times 10^4$  cm<sup>2</sup>/V.s at 20K.



Figure 1. Hall mobility of PbTe layer as a function of temperature grown on  $BaF_2$  substrate and on  $BaF_2/CaF_2/Si$  structure.

As the quality of the  $CaF_2$  first buffer layer showed to be the most important in the whole structure, we decided to make a detailed study on the  $CaF_2/Si$  structure.

## IV. High resolution X-ray diffraction measurements

In order to analyze the CaF<sub>2</sub>/Si(111) structurally, x-ray diffraction techniques were used. In standard xray diffraction  $\Theta/2\Theta$  scans, the CaF<sub>2</sub>/Si(111) showed only the symmetrical Bragg diffraction peaks. It was not possible to distinguish the CaF<sub>2</sub> peaks under the Si ones using this technique. High resolution x-ray diffraction (HRXD) measurements were then performed using a two-crystal diffractometer with a GaAs crystal monochromator for the CuK $\alpha$  radiation with a final resolution of 17 arcsec.

Fig. 2 shows a  $\omega$  scan of a (15 nm) CaF<sub>2</sub>/Si structure where the CaF<sub>2</sub>(333) Bragg peak appears as a shoulder near the Si(333) peak. For comparison the (333) peak of a Si substrate is also plotted in the figure. This position of the CaF<sub>2</sub>(333) Bragg peak cannot be explained with a simple pseudomorphic growth model since in this case a compressive strain is expected due to the larger CaF<sub>2</sub> lattice constant relative to Si. The tensile strain in the plane of the CaF<sub>2</sub> layer can only be understood if the large difference in thermal expansion coefficients between  $CaF_2$  and Si is taken into  $account^{[6,7]}$ . This difference increases the lattice misfit between  $CaF_2$  and Si from 0.5% at room temperature to 2.2% at growth temperature (730°C).



Figure 2.  $\omega$  scan of the (333) Bragg reflection of the CaF\_2/Si structure and of a Si substrate.

In order to calculate the angular separation  $\Delta \omega$ between the CaF<sub>2</sub>(333) and the Si(333) symmetrical Bragg peak, the elastic theory combined with Bragg diffraction law was used. The relation of the perpendicular ( $\epsilon_{\perp}$ ) to the parallel strain ( $\epsilon_{\parallel}$ ) taken from the elastic theory for a cubic crystal in the (111) direction assuming tetragonal distortion and using bulk elastic constants  $C_{ij}$  for the CaF<sub>2</sub> layer is:

$$\epsilon_{\perp} = -2 \left[ \frac{(C_{11} + 2C_{12} - 2C_{44})}{(C_{11} + 2C_{22} + 4C_{44})} \right] \epsilon_{\parallel} = -1.01 \epsilon_{\parallel}$$

In case of a pseudomorphic growth (compressive strain with  $\epsilon_{\parallel} = -0.5\%$ ) the calculated angular separation is  $\Delta \omega = -0.56^{\circ}$  or in case of a completely relaxed CaF<sub>2</sub> layer on Si at room temperature ( $\epsilon_{\parallel} = 0$ ),  $\Delta \omega = -0.28^{\circ}$ .

In the proposed model, considering that the  $CaF_2$ layer is completely relaxed at growth temperature where interfacial defects such as misfit dislocations accommodate the 2.2% lattice misfit, a tensile strain in the layer is built up during cooling down to room temperature due to the thermal mismatch. Assuming complete pinning of these defects (no atomic rearrangement) during cooling, the maximum tensile strain in the layer is  $\epsilon_{\parallel} = 1.7\%$ . In this case we calculate  $\Delta \omega = 0.81^{\circ}$ . If the CaF<sub>2</sub> layer is assumed to be totally relaxed at growth temperature and R represents the degree of relaxation in % of the thermal strain during cooling (R=0 for maximum thermal strain  $\epsilon_{\parallel} = 1.7\%$  and R=100 for  $\epsilon_{\parallel} = 0$ ), a relaxation degree R=75% during cooling down process is necessary to fit the position of the measured CaF<sub>2</sub>(333) peak. This result agrees with the tetragonal distortion determined previously from RBS ion channeling measurements in CaF<sub>2</sub>/Si system<sup>[6,7]</sup>.



Figure 3.  $\omega$  scan of the (333) Bragg diffraction peak of the CaF<sub>2</sub>/Si structure after thermal cycles from 300 to 77K: (a) as grown sample, (b) after 15 sec immersed in LN<sub>2</sub> and (c) after +1 min immersed in LN<sub>2</sub>.

To simulate the real operating conditions of infrared detectors, a 15 nm thick  $CaF_2$  layer on Si was submitted to various thermal cycles from 300 to 77K. Fig. 3 shows the HRXD spectrum of the sample measured after each thermal cycle. To distinguish between the two (333)

diffraction peaks in this structure, the spectra were fitted with a double Gaussian. As shown in Fig. 3b, after the first thermal cycle with 15 seconds immersion in  $LN_2$  the CaF<sub>2</sub>(333) Bragg peak shifts towards the Si(333) peak position causing an increase in intensity of the total peak  $\omega(333) = \operatorname{CaF}_2(333) + \operatorname{Si}(333)$ . However, after the second thermal cycle with +1.0 min immersed in LN<sub>2</sub> (Fig. 3c) the  $\omega(333)$  structure broadens and decreases its intensity. For subsequent thermal cycles with larger periods of time, the form of this  $\omega(333)$  structure remains almost the same as the one plotted in Fig. 3c. The broadening of the total x-ray rocking curve  $\omega(333)$ is related to an increase in the defects density at the CaF<sub>2</sub>/Si interface. This increase in defect density and the modification of the strain in the  $CaF_2/Si$  interface may cause some effects in the lead chalcogenides layers grown on top of the BaF<sub>2</sub>/CaF<sub>2</sub>/Si structure when cooled to the liquid nitrogen temperature.

#### V. Conclusion

As demonstrated by HRXD and Hall effect measurements, the control on the  $CaF_2$  layer strain appears to be very important to optimize the Hall mobility of PbTe layers grown on Si with fluoride buffer layers.

The large difference in thermal expansion coeffi-

cients between  $CaF_2$  and Si is responsible for the inplane tensile strain observed in the  $CaF_2$  layer grown on Si. The shift and broadening of the  $CaF_2(333)$  Bragg peak observed when submitting the  $CaF_2/Si$  structure to thermal cycles from 300 to 77K indicate that some degrading effect on infrared devices grown on top of the  $BaF_2/CaF_2/Si$  structure may appear.

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